

Hazy3

a brief introduction
to Cloudy 95

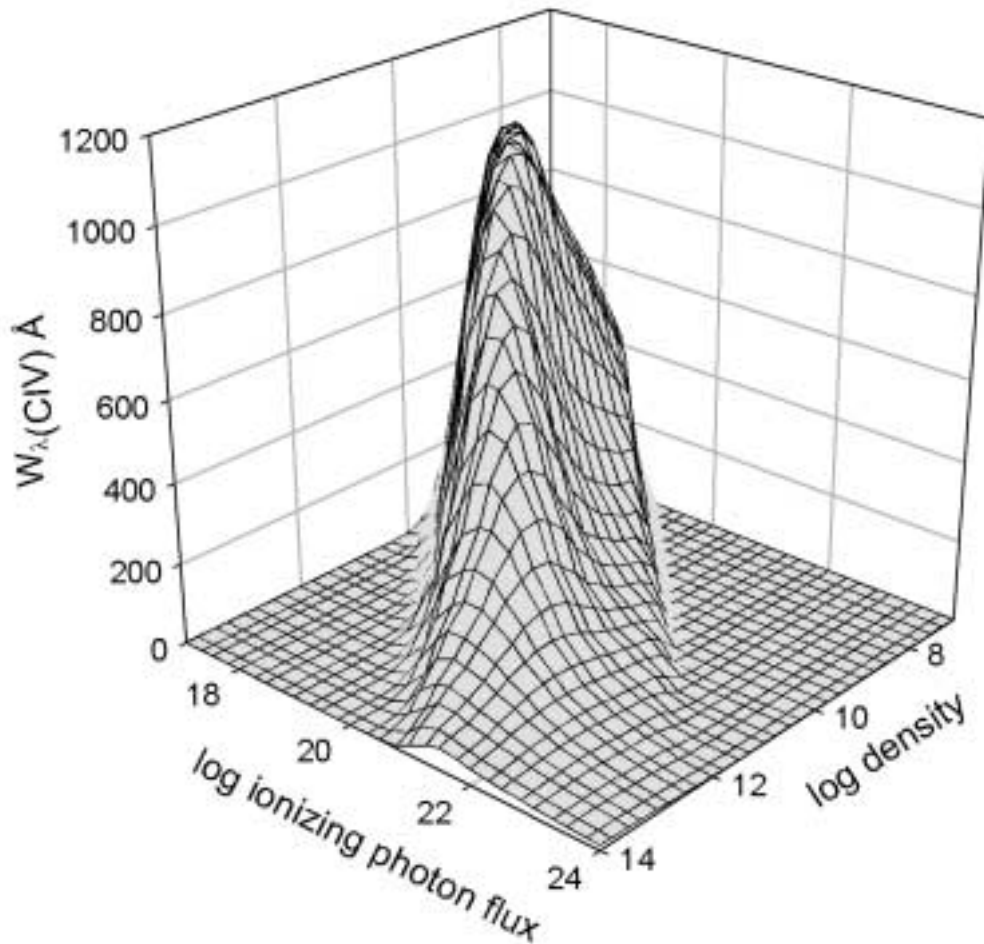
results, environment

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<http://www.pa.uky.edu/~gary/cloudy>



Use of this program is not restricted provided each use is acknowledged upon publication. The bibliographic reference to this version of Cloudy is Ferland, G.J., 2001, *Hazy, a Brief Introduction to Cloudy*, University of Kentucky Department of Physics and Astronomy Internal Report.

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Cloudy is an evolving code. Updates are made on a roughly quarterly basis, while major revisions occur roughly every three years. You should confirm that you have the most recent version of the code by checking the web site <http://www.pa.uky.edu/~gary/cloudy> or asking to be placed on the Cloudy mailing list.

CLOUDY 95

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1 MACHINE ENVIRONMENT

1.1 Overview

This section describes the machine environment needed to run Cloudy. The code is designed to run on IEEE 32-bit computers, and has been tested on a Sparc, SGI, Alpha, HP NCX, and PC. Calculations on all machines produce essentially identical results.

The `readme.htm` file that is included with the source contains the most recent information about compiling and setting up the code.

1.2 Floating Point Environment

The floating-point environment should be set to ignore floating-point underflow but crash on any other floating-point error. Floating-point underflow is an unavoidable consequence of the attenuation of radiation as a beam of light is extinguished by an absorbing medium; underflow error checking should be disabled.

Floating point overflow or division by zero *must never* occur, nor should library function domain errors (i.e., the log of a negative number). I would appreciate hearing about these errors. I can't fix it if I don't know it is broken. My email address is `gary@cloud9.pa.uky.edu`. Please send the input file and version number.

1.3 Cloudy via ftp or the Web

The home page for Cloudy is `http://www.pa.uky.edu/~gary/cloudy`. This is the preferred method of obtaining the code. This is the best place to retrieve the source since it is the only place certain to have the current version.

Cloudy is an evolving code. Bugs are fixed as soon as they are discovered, and major revisions occur from time to time. It is important to check the web page several times per year to see what changes and improvements have occurred. I also maintain a mailing list to announce changes to the code. To be placed on this list send a request to `gary@pa.uky.edu`.

I would like to keep track of what the code is being used for, if only to try to make sure that it is within its range of validity. I would appreciate receiving preprints of any work done using Cloudy.

The bibliographic reference to this version of Cloudy, which should be cited in any paper that uses it, is given on the inside cover of this document (Ferland, G.J., 2001, *Hazy, a brief introduction to Cloudy*, 95.01d). The citation should include the subversion number, i.e., version 95.01d.

1.4 Cloudy in Lexington

1.4.1 Versions of the Code

Many versions of Cloudy live on the Clouds domain in Lexington. In general there will two versions of the code — a fully tested (older) version, and the current development version.

Software development goes through several phases. The initial version of a code is often referred to as an “alpha” version, the nearly de-bugged version as a “beta” version, and the fully tested version as the “gold” code. At any one time the last gold version of the code will live in a directory with a name indicating the version number. The version of Cloudy is indicated by two parts - an integer indicating the version number, and a decimal part (perhaps with a letter following it) indicating the sub-version. Sub-versions to the gold code indicate fixes of bugs or stability problems, or major changes in the atomic data. The last gold version of Cloudy was 95. An example of the current gold version number might be something like 95.09d.

1.4.2 Running a single model

Cloudy is often used to read in the parameters for a single model, and compute the result. The easiest way to do this is to create a small file that contains the input stream for that model. As a typical case consider a simple planetary nebula:

```
hden 4
radius 17
black body 100,000K, luminosity 38
```

Assume this is saved as the file **pn.in**. Note that Cloudy stops reading the input stream when it reaches a blank line, or the end of file. No special end of input sentinel is needed.

I created a shell script with the name **run** which is in my “bin” directory, which I include on my path. The shell script **run** consists of the following:

```
echo reading input file $1.in
case $# in
0) echo there must be an input file ;;
1) /homeb/uwc0/gary/cloudy/c.sun4<$1.in >$1.out
   echo created output file $1.out ;;
2) /homeb/uwc0/gary/cloudy/c.sun4 < $1.in >$2.out
   echo created output file $2.out ;;
esac
echo $p
exit
```

run normally works by reading in a file called **file.in**, and creating a file called **file.out**.

If **run** is executed with no input parameters it will complain that at least one argument is needed and then stop. If there is one parameter it is treated as the name of the input and output files. So in the above example, typing

```
run pn
```

would read the input stream in **pn.in** and create an output file of results called **pn.out**. When two parameters occur the first is the name of the input stream and the second is the name of the output stream. The example

run `pn test`

would read the file `pn.in` and create the file `test.out`.

2 CLOUDY AS A SUBROUTINE

2.1 Overview

It is possible to use Cloudy as a subroutine of other, much larger, codes. When used this way a series of subroutine calls, described next, are used to initialize the code, specify the initial conditions, drive the code, and finally examine the predictions.

A common strategy is to call the code to compute line intensities for a large matrix of parameters. The results of one such calculation is shown in Figure 1. Such grids can be computed in a few dozen hours on modern workstations, and offer far greater insight to physical effects of changing model parameters, than does a single model.

This document gives an overview of all the routines that are intended to be “public” (needed to be accessed by programs that will call Cloudy). The definitions for all public routines are contained in the header file *cd drive.h*, which gives the best and current description of all these routines. That file is the definitive reference source for all of the material in this chapter.

2.1.1 Creating a new main program

In C there must be exactly one main program, called *main*. In the distribution this

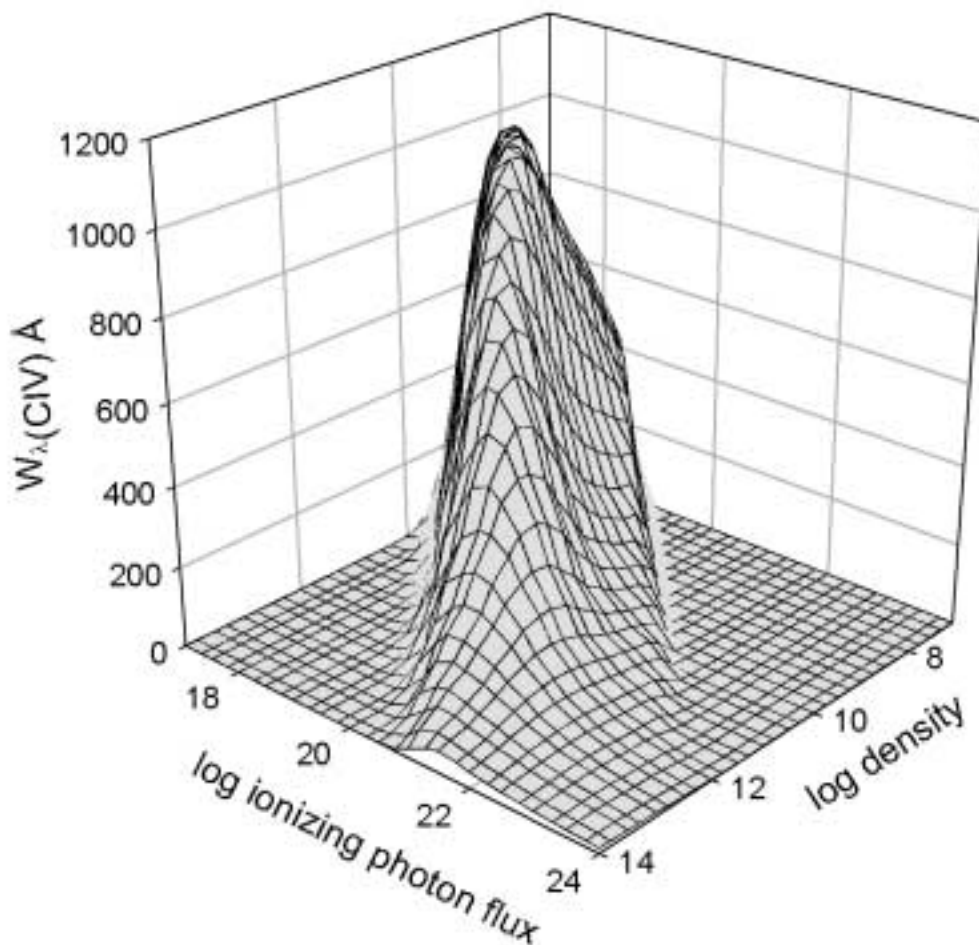


Figure 1 The results of a large grid of model calculations are shown. The x-y plane show the logs of the hydrogen density and stellar temperature. The z axis is the indicated line ratio, for solar abundances.

routine is within the file *maincl.c*. The goal is to replace the existing Cloudy main program with one that you write. The file *maincl.c* that is included in the distribution must be deleted so that the program you write will be loaded instead. The remaining routines are then compiled with a command like the following:

```
gcc -c *.c
```

which will create a large number of object files. Often the new main program will be linked with these object files with a command something like

```
gcc newmain.c *.o -lm
```

The following subsections outline how to write code for this new main program.

2.2 Writing a New Main Program

2.2.1 The *cddrive.h* header file

The file *cddrive.h* contains definitions of all public routines, the routines that a user would call to drive Cloudy. That file is the definitive reference for the material contained in this section. Comments within the file explain all routines and their parameters. Often I copy from the directory containing the majority of the Cloudy source to other directories where special main routines are developed. It is very important that all versions of *cddrive.h* be updated with the source is updated since the file changes with time.

2.2.2 A note on return conditions

Some of the routines return a value to indicate success or failure. I try to follow the C and Unix conventions to indicate success with zero and trouble with a non-zero return. This rule is not always followed, however, and *cddrive.h* should be consulted to make sure the return conditions are understood.

2.2.3 Initializing the code

Many variables must be initialized at the beginning of the calculation. Calling routine *cdInit* does this.

```
cdInit();
```

Routine *cdInit* must be called every time a new calculation is to be performed, *before* calling any of the following subroutines, but after the results of previous calculations have been read. (The results of the previous calculation are lost when *cdInit* is called.)

2.2.4 *cdMPI* - telling the exit handler about MPI

If MPI is used for parallel processing then, when the calculation is complete, the code must clean up by calling *MPI_Finalize* at exit. The exit handler is defined in *cdInit* and is called *exithandler*. Routine *cdMPI* must be called after *cdInit* and before *cdDrive* for *MPI_Finalize* to be called correctly.

2.2.5 *cdRead* - entering Commands

Command lines are entered by successive calls to routine *cdRead*. The argument of *cdRead* is the line image, a null-terminated string. The commands consist of a

series of 80 character free-format command lines. They must obey all the rules outlined in Part I.

In the examples below some commands are directly entered as strings and others created by writing variables through ***sprintf*** (a standard C io function).

```
char chLine[132];/* this will hold the command lines we will generate*/

/* this example send the string straight to cdRead */
cdRead("title Gene's series of constant pressure models" );
hden = 5.4;

/* this example writes a variable to a string then send the string to cdRead */
sprintf( chLine , "hden %5.2f ", hden );
cdRead(chLine );

/* this example sends a string with double quotes, and so must "escape" them */
cdRead("set path\"d:\\projects\\cloudy\\ccloudy\\data\" " );

sprintf( chLine , "coronal %5.2f ", temp );
cdRead(chLine );

cdRead("stop zone 1 " );
```

cdRead returns the number of commands that can still be entered before exceeding the size of the storage arrays. So this routine is an exception to the general rule that a zero return condition indicates success – here it indicates a problem -- no further commands can be entered. The return value was ignored in the examples above.

It is not now possible to read in more than 1000 command lines because of limits to the size of the character arrays used to store them. This limit is stored as the variable ***NKRD***. If more than 1000 lines are read in by calling ***cdRead*** then ***cdRead*** will stop after explaining why. It will be necessary to increase ***NKRD*** if more than 1000 command lines are needed.

2.2.6 ***cdOutp*** - sending output to a file

Cloudy normally writes its standard output on the system's ***stdout***. This can be changed to another stream by calling routine ***cdOutp***, which has a file handle to an open file as its single argument. By combining this redirection with the C ***fopen*** statement it is possible to have the standard output sent into any file.

```
#include "cddrive.h"
/* this defines a standard C file handle */
FILE *ioData ;

/* open the file output.txt for writing */
ioData = fopen("output.txt","w");

/* ioData is equal to NULL if we failed to open the file */
if( ioData==NULL )
{
    exit(1);
}

/* send output to this file*/
cdOutp( ioData ) ;
- - - code goes here
/* at end of calculation close the file */
fclose(ioData);
```

2.2.7 *cdTalk* - produce output??

Cloudy normally speaks what's on its mind. However, it does have a quiet mode in which nothing at all is printed. This quiet mode is set by the logical argument to subroutine ***cdTalk***.

```
#include "cddrive.h"
/*set no output at all*/
cdTalk( FALSE )
/*have the code produce the normal printout*/
cdTalk( TRUE )
```

The default is for Cloudy to produce output, and ***cdTalk*** does not have to be called if this is desired. However, it does need to be called with the logical variable ***FALSE*** if the quiet mode is desired. (TRUE and FALSE are defined in cddrive.h).

2.2.8 *cdDrive* - calling the Code

The calculation is performed when routine ***cdDrive*** is called. ***cdDrive*** returns an int indicating whether the calculation was successful. The value 0 indicates a successful calculation. The following shows an example of its use.

```
int lgOK;
if( cdDrive() )
{
    printf("problems!\n");
    exit(1);
}
```

If problems occurred and the results cannot be trusted then the return value is non-zero. This will only be set if the calculation suffered a complete meltdown. Routine ***cdNwcns*** (page 411) can be called to find out about less serious problems.

2.2.9 *cdNoExec* - checking without Computing

If routine ***cdNoExec*** is called after ***cdInit*** but before ***cdDrive*** then only the initial parts of a calculation will be performed when routine ***cdDrive*** is called.

```
CdInit();

/*read in commands */
cdRead( . . . );

/*tell it not to execute */
cdNoExec();

/*call the code */
lgBad = cdDrive();
```

When ***cdDrive*** is called after ***cdNoExec*** the code will generate the incident continuum, set the initial density, and the chemical composition. It will then stop just before the initial search for the physical conditions in the first zone. All of the initial printout, summarizing properties of the composition and continuum, will be generated. This provides a quick way to check that a large grid of models will be specified correctly, without actually fully computing the grid.

2.3 Checking Predictions

2.3.1 *cdLine* - emission lines intensities

The predicted line intensities or luminosities for all lines with non-zero values are stored within a set of structures which also contain the line identifiers, a four character label and integer wavelength. These are normally printed at the end of the calculation. You obtain the line intensity by calling subroutine ***cdLine***. The label and wavelength of the line are specified in the call, and the routine returns the relative intensity and the log of the absolute intensity or luminosity.

```
Include "cddefines.h"
    double relint , absint;
*
if( cdLine( "totl" , 4861 , &relint , &absint ) <= 0 )
    printf("did not find this line\n");
```

The first variable in the call is the line label, the four-character null-terminated string (upper or lower case) as used by the code to identify the line. The second integer variable gives the wavelength of the line. Both of these must exactly match the label and wavelength used by Cloudy to identify the line (see the chapter “Lines” for a full description). The third variable (***relint*** in the above example) is a pointer to the double precision relative intensity of the line (relative to the normalization line, usually H β , but set with the ***normalize*** command). The log of the intensity (erg cm⁻² s⁻¹) or luminosity (erg s⁻¹) of the line is returned as a double precision variable (***absint*** in the above example). If the intensity of the line is zero or the line was not found then this variable will be set to -37.

If ***cdLine*** finds the line it returns the index of the line within the stack of emission lines. It returns the negative of the total number of lines in the stack if the line is not found. This may occur if the line wavelength or label were mistyped. This is an exception to the normal, C-like, function return convention in which a normal return is zero and an abnormal return non-zero. A positive value indicates a successful return.

2.3.2 *cdGetLineList* - special arrays of emission lines

The routine ***cdGetLineList*** provides a way to access a block of lines in an automatic manner. This provides a way to access sets of predictions, to be stored or printed.

cdInit must be called first to initialize needed variables. Then ***cdGetLineList*** is called as described here. Finally, the actual grid of calculations begins with another call to ***cdInit***. Finally the intensities of a set of lines are extracted by calls to ***cdLine***. So the first call to ***cdInit*** is generally not followed by execution of the code when these routines are used.

When routine ***cdGetLineList*** is called a list of lines, from either one of two files included in the standard data distribution, or a user-created special file, will be entered into a pair of vectors. This list can then be used to call ***cdLine*** to obtain intensities of these lines. The argument to routine ***cdGetLineList*** is the name of the file containing the lines. If a null string is passed (“”) then *BLRLineList.dat* is used. That file is included in the standard data distribution and contains a description of the required file format. A second file, *NLRLineList.dat*, is also included in the

standard distribution and contains lines appropriate for a low-density gas. The file location is assumed to be on the path as set with the `path` command, or in `path.c`.

When the routine has been called the code creates a pair of vectors, `cdGetchLabel[n]` and `cdGetnWL[n]`. The first is an array of four-character strings that are emission line labels. The second vector is a long integer array of wavelengths. This string and integer are the label and wavelength used to identify the lines. The function returns a long int, the number of lines in the list. If problems occurred then a negative integer is returned.

The following shows an example of getting the lines from *BLRLineList.dat*, executing the code, then obtaining the predicted intensities by calling ***cdLine***.

```
/* initialize the code */
double pred[LargeEnough];
cdInit();
/* get list of lines from standard data file */
if( (nLines=cdGetLineList("")) < 0 )
{
    /* this is the error exit - could not obtain the lines */
    exit(1);
}
- - - - -
/* now set up the actual call to the code */
cdInit();
/* missing commands here, then call the code */
cdDrive();
/* now save the emission lines that were in the file into the array pred */
for( n=0; n<nLines; ++n )
{
    lgOK = cdLine( cdGetchLabel[n], cdGetnWL[n] , &relative , &absolute );
    if( lgOK<= 0 )
    {
        fprintf(stderr,"did not find %4s%5li\n",cdGetchLabel[n],cdGetnWL[n]);
        fprintf(ioDATA,"ndid not find line.\n");
    }
    pred[n] = relative;
}
}
```

2.3.3 *cdEmis* - emissivity of lines

A second routine, ***cdEmis*** functions much the same as ***cdLine*** (page 408 above) but returns the local emissivity ($\text{erg cm}^{-3} \text{ s}^{-1}$ for unit filling factor) of the line for the last computed zone. The return value is the index of the line in with the line stack if it was found, and the negative of the number of lines in the stack if the line could not be found.

2.3.4 *cdColm* - the computed column densities

The predicted column densities can be accessed by calling the routine ***cdColm***.

```
/* want C+2 */
if(cdColm("carb", 3 , &column))
{
    printf(" could not find C+2 - check in put\n");
}
else
{
    printf("The predicted C+2 column density is %e\n", column );
}
}
```

The routine returns 0 if it found the species, and 1 if it could not. It returns the predicted column density (linear, cm^{-2}) as the third argument. The first argument

chLabel is a four-character identifier that must agree with the first four characters (upper or lower case) of the name used to indicate the element in the printout. The integer variable **ion** is the spectroscopic designation of the level of ionization, i.e., 1 indicates C⁰, 3 indicates C⁺², etc.

2.3.5 *cdIonFrac* - the computed ionization fractions

The predicted ionization fractions, averaged over radius or volume, can be accessed by calling the subroutine **cdIonFrac**.

```
*
    if( cdIonFrac( "carb" , 2 , &frac , "radius",FALSE) )
    {
        exit(1);
    }
    printf("The predicted ionization fraction over radius is is%g\n", frac);
```

The routine returns the predicted ionization fraction A_{ion}/A_{tot} . **chLabel** is a four-character identifier that must agree with the first four characters (upper or lower case) used to indicate the element in the printout. The integer variable **ion** is the spectroscopic designation of the level of ionization, i.e., 1 indicates C⁰ or the atom, 3 indicates the second ion C⁺², etc. **chWeight** is a six-character variable (plus end of string sentinel) which must be either "radius" or "volume" (either upper or lower case). This string determines whether the ionization fraction returned is weighted with respect to radius or volume. The last variable determine whether (TRUE) or not (FALSE) the ionization fraction is also weighted with respect to the electron density. The function returns zero if the ion was found, and non-zero if an error occurred.

2.3.6 *cdTemp* - the computed mean temperature

Routine **cdTemp** returns the computed mean temperature. The first parameter is a four character null terminated string giving the first four letters (upper or lower case) of the name of the element as spelled by the code. The second parameter is the ionization stage, with 1 for the atom, 2 the first ion, etc. The third parameter will be the computed mean temperature. The last parameter is a six character null terminated string, either "radius" or "volume" that says whether the temperature should be weighted with respect to radius or volume. The routine returns 0 if it finds the species, and 1 if it could not find the species.

```
*
    if( cdTemp( "carb" , 2 , &temp , "radius") )
    {
        exit(1);
    }
    printf("The mean temperature is%g\n", temp);
```

The label "21cm" for the element name has a special meaning. The temperature will be the mean of n_H/T , the harmonic mean temperature weighted with respect to the atomic hydrogen density, averaged over radius. This is the temperature measured by combined 21 cm -- L α observations.

2.3.7 *cdGetCooling, cdGetHeating* – last zone's cooling or heating

The total cooling and heating rates (erg cm⁻² s⁻¹) for the last computed zone is obtained by called the void functions **cdGetCooling** and **cdGetHeating**. In equilibrium the cooling and heating rates are equal.

2.3.8 *cdGetTe* - reading the temperature

The electron temperature for the last computed zone is obtained by calling the function ***cdGetTe***. The function has no arguments and its return value is the predicted temperature.

2.3.9 *cdGetPres* – pressure

The pressure for the last computed zone is obtained by calling routine ***cdGetPres***. This routine has three arguments, pointers to the total (gas plus radiation) pressure, the gas pressure, and the radiation pressure. All are double precision variables.

2.3.10 *cdTimescales* – several timescales.

This routine has three arguments, pointers to doubles that will return the thermal timescale, the hydrogen recombination timescale, and the H₂ formation timescale.

2.4 Other information

cdDate(string) The date when the current version of the code was released will be placed as a null terminated string into the string passed as an argument. The calling program must allocate enough room in the string.

cdVersion(string) The code's version number will be placed as a null-terminated string into the string passed as an argument. The version number is a string rather than a float since it usually ends with a letter of the alphabet. The calling program must allocate enough room in string.

double cdExecTime(void) This returns the time that has elapsed since the previous call to ***cdInit***.

2.5 Printing the comments

After the calculation is complete, but before the emission lines are printed, the code generates a series of statements that indicate warnings, cautions, comments, and surprises. These should be examined to confirm that the calculation is likely to be all right. A series of public routines allows the driving code to determine whether these comments were generated, what type they were, and to print them into an arbitrary open file.

2.5.1 *Were comments generated?*

A series of routines are provided to check on problems in the calculation.

Routine ***cdNwcns*** will return the number of warnings, cautions, surprises, notes, and temperature and pressure failures:

```
cdNwcns( &lgAbort , &nw , &nc , &nn , &ns , &nte , &npe , &none , &neden )
```

where the first variable is a flag indicating whether the calculation aborted, ***nw*** is the number of warnings generated (if this number is non-zero, then the calculation has serious problems), ***nc*** is the number of cautions generated (these are less severe than warnings, but are still a cause of concern), and ***nn*** and ***ns*** are the number of notes and surprises. The next two arguments are the number of temperature and pressure failures, and both should be zero. The last two are the number of ionization and

electron density failures, again hopefully zero. All of these integer values will add up to zero in a successful calculation.

If either of the first two variables are non-zero then the code returned with an indication of serious failure. An abort is far more serious than a warning since it indicates catastrophic meltdown. I would appreciate learning about these.

2.5.2 Printing the comments.

The comments that normally appear after the last zone may be printed on any file by calling the series of subroutines described here. In all cases the routines take as an argument a file handle, which must point to a file open for writing.

```
/* output the comments into a file, but first open it */
FILE *ioOUT;
If( (IoOUT = fopen( "comments.txt", "w" ) ) == NULL )
{
    printf("error creating comments.txt file\n");
    exit(1);
}
/*print the reason the calculation stopped, and geometry*/
cdReasonGeo( ioOUT )
/*print the warnings*/
cdWarnings(ioOUT)
/*next print the cautions*/
cdCautions(ioOUT)
/*now print any surprising results*/
cdSurprises(ioOUT)
/*now print the notes
cdNotes(ioOUT)
```

cdReasonGeo (FILE *io). It is very important to understand why the calculation stopped. The first two lines after the last zone results give the reason the calculation stopped, and the type of geometry. This information will be printed on the file whose handle is the argument.

cdWarnings(FILE *io) All warnings will be printed on the output file.

cdCautions (FILE *io) All cautions will be printed on the output file.

cdSurprises(FILE *io) All surprises (denoted by a "!") are printed on the output file.

cdNotes(FILE *io) The notes concerning the calculation will be printed on the output file.

2.5.3 cdErrors(FILE *io) - printing a summary of any problems

Routine ***cdErrors(FILE *io)*** will generate a summary of any problems that happened during a calculation. The argument is a file pointer to the output file where the results will be placed. If problems occurred in the calculation, such as temperature or pressure failures, warnings, or cautions, these will be printed following the title for the calculation.

2.5.4 cdPrintCommands(FILE *io) – print the command stack

The single argument is a file handle. The entire series of input commands will be written into the file. The commands are preceded and followed by lines that begin with "c =====" to easily identify where these start.

2.5.5 cdSPEC – get continuous spectra

This routine provides an interface between Cloudy and Keith Arnaud's X-Ray spectral analysis program XSPEC. It is called after **cdDrive** has computed a model. Depending on which option is used, it will return the incident continuum, the attenuated incident continuum, the reflected continuum, the diffuse continuous continuum, outward direction diffuse continuous emission, reflected lines, outward lines. All are $4\pi \nu J_\nu$ and have units of $\text{erg cm}^{-2} \text{s}^{-1}$.

All lines and continuum emitted by the cloud assume full coverage of the continuum source. Details are given in *cdDrive.h*.

2.6 Example Call as a Subroutine

The following is an example of a very simple use of Cloudy as a subroutine.

```
/*main program that calls cloudy when used as a stand-alone program */
#include <stdio.h>
#include <stdlib.h>
#include <math.h>
#include <string.h>
#include " cddrive.h"

int main( void )
{
    int lgOK ;
    FILE *ioDATA ;

    ioDATA = fopen("d:\\projects\\cloudy\\run2\\run2a.out","w");
    if( ioDATA == NULL )
    {
        printf(" could not open xspec.out for writing.\n");
        exit(1);
    }
    /* turn off buffering so we results as they happen */
    setbuf( ioDATA , NULL );

    /* initialize the code */
    cdInit();
    cdOutp(ioDATA);
    cdRead( "hden 10.5 " );
    cdRead( "agn 6.00 -1.40 -0.50 -1.0 " );
    cdRead( "phi(h) 23 " );
    cdRead( "stop column density 21.860889 " );
    cdRead( "**constant temper 727,000 " );
    cdRead( "background z=0 " );
    cdRead( "failures 3 no map " );
    /* actually call the code */
    lgOK = cdDrive();
    fclose( ioDATA );
    eexit(0);
}
```

2.7 Storing Grids of Calculations¹

This subsection describes a strategy for creating and analyzing a large grid of models. In this strategy all results are saved to a file, which is then read in later.

The **punch results last** command is described in Part I. The line intensities and column densities computed by the code will be written into a file when this command is entered. Any post-processing software can read this file, but reading this file is facilitated by routine **cdGett**, described next.

¹ This method is seldom used today and will be removed in a future version of the code.

The output produced by the **punch results last** command can be read with routine **cdGett**. **cdGett** has one argument, a pointer to the file containing the results of the **punch results** command. Its return value indicates whether (TRUE) or not (FALSE) the end of file was encountered when the model results were read. If the logical variable is TRUE then the end of file was encountered and no results were read. If it has not been encountered then there may be more models further down the input stream. The next model in the grid will be accessed by the next call to **cdGett**.

Once **cdGett** has been called the routines **cdLine**, **cdIonf**, and **cdColm** routines can be used to obtain emission line intensities, ionization fractions, or column densities from the stored file. A structure, **getpar**, is associated with **cdGett**. This contains information concerning the initial parameters of the current model, as specified in the input stream.

This method is seldom used today. The method described next is simpler.

2.8 Computing Grids of Calculations

Today I usually use the code to compute results, extract information on the fly, and then save desired quantities. The following examples illustrate producing a series of models with increasing stellar temperature, in which the stellar temperature and the stellar temperature and [O III]/H β ratio are written to a file.

In the following only one call to **cdLine** is made to get results for a single line. In practice all desired lines are extracted at this stage and stored in a format where it can be read by other software. A call to **cdGetLineList** (page 408 above) provides an easy way to obtain large numbers of lines whose labels are stored in a file.

```
*
/* very simple main program to call cloudy as a stand-alone program */
#include <stdio.h>
#include <stdlib.h>
#include <math.h>
#include <string.h>
/* this must point to cddrive.h */
#include "d:\projects\cloudy\current\source\cddrive.h"

/*int main( int argc, char *argv[] )*/
int main( void )
{
    /* this will hold images of the command lines */
    char chCard[200];
    double TStar , rel , absol;
    FILE *ioDATA ;

    /* open file for writing some results */
    if( (ioDATA = fopen("1DGrid.out","w")) == NULL )
    {
        printf(" could not open 1DGrid.out for writing.\n");
        exit(1);
    }

    TStar = 3e4;
    while( TStar < 5e4 )
    {
        /* initialize the code */
        cdInit();
        /* redirect output to the file we opened above */
        cdOutp(ioDATA);
        /* but also say we want to output by passing 0, for FALSE */
        cdTalk( 0 );
        /* write variables into strings and send string as input file */
        cdRead( "hden 5 " );
    }
}
```

```

cdRead( "ionization parameter -2 ");
cdRead( "stop zone 1 ");
/* this is example of writing a variable into the string then passing
 * the string to cloudy */
sprintf( chCard , "blackbody, T= %f" , TStar );
cdRead( chCard );
/* actually call the code */
if( cdDrive() )
{
    printf("Beware: Cloudy returned error condition, so exit with 1.\n");
}
/* now increment TStar by 5000K */
TStar += 5000.;
/* get intensity of [OIII] relative to Hbeta - remember cdLine is different
 * from most cd routines since it returns element within line stack, 0 for failure */
if( !cdLine( "o 3" , 5007 , &rel , &absol ) <=0 )
{
    printf("could not find 5007\n");
    exit(1);
}
/* now print stellar temperature and 5007/Hbeta ratio */
fprintf(ioDATA , "%.0f %.2f\n", TStar , rel );
}
/* things were ok, exit with 0 */
exit(0);
}

```


3 OUTPUT

3.1 Overview

This section defines the output produced by Cloudy. Each section begins with a sample of the output described, and then goes on to describe the meaning of the printout in greater detail. The output actually shown is from the Meudon (Pequignot 1986) and Lexington (Ferland et al 1995) meetings Planetary Nebula test case.

3.2 Header Information

Several lines of output echo the input commands and outline some properties of the initial continuum.

3.2.1 Initial Information

```

Cloudy 95.00

*****01Feb01*****
*
* title "New" Paris meeting Planetary nebula
* c recompute "standard" PN model of the Pequignot Meudon Conference
* init file='c84.ini'
* elements read
* sphere
* black body, T=150,000K radius = 10
* hden = 3.4771213
* radius = 17
* abund -1 C-3.523 N-4.0-3.222 ne-3.824 na--10 mg-4.523 al--10
* continue si-4.523 s-4.824 ar-10 ca--10 fe-10 ni--10
* plot continuum range .1
* punch overview last 70
* c parispn.in
* c Sun
*
*****

Chemical Composition
H :1.00E+00 He:1.00E-01 Li:2.04E-09 Be:2.60E-11 B :7.60E-10 C :3.00E-04 N :1.00E-04 O :6.00E-04 F :3.02E-08
Ne:1.50E-04 Na:1.00E-10 Mg:3.00E-05 Al:1.00E-10 Si:3.00E-05 P :3.73E-07 S :1.50E-05 Cl:1.88E-07 Ar:1.00E-10
K :1.35E-07 Ca:1.00E-10 Sc:1.22E-09 Ti:8.60E-08 V :1.05E-08 Cr:4.84E-07 Mn:3.42E-07 Fe:1.00E-10 Co:2.24E-09
Ni:1.00E-10 Cu:1.87E-08 Zn:4.52E-08

```

This begins with the version number of Cloudy. Major revisions, which have noticeable effects on the emission-line spectrum or which reflect significant improvements in the physics, are denoted by integer increases in the version number, while minor changes increment the revision number by 0.01. In a static version of the code, small changes (usually minor bug fixes) are denoted by letters (i.e., .02a). The following line gives the date this version was created.

All of the input command lines, with the exception of those starting with a #, %, or *, are echoed before the calculation begins, and are saved to be reprinted after the calculation is completed.

The input information is followed by the chemical composition of the gas. The numbers are the number densities of the elements, relative to a hydrogen abundance of unity. Only the active elements are included (those turned off with the **elements off** command are not printed).

3.2.2 Properties of the Continuum

```

1568Cell Peak3.83E+00 Lo 1.00E-05=0.9108cm Hi-Con:7.63E+01 Ryd E(hi):7.35E+06Ryd E(hi): 100.01 MeV
L(nu>1ryd): 37.5396 Average nu:2.935E+00 L(X-ray): 31.4007 L(BalC): 36.1444 Q(Balmer C): 46.9767
Q(1.0-1.8): 47.1613 Q(1.8-4.0): 47.4535 Q(4.0-20): 47.0522 Q(20--): 40.7275 Ion pht flx:4.312E+12
L(gam ray): 0.0000 Q(gam ray): 0.0000 L(Infred): 34.4845 Alf(ox): 0.0000 Total lumin: 37.5571
log L/Lsun: 3.9743 Abs bol mg: -5.1858 Abs V mag: 2.4170 Bol cor: -7.6028 nuFnu(Bbet): 34.5868
U(1.0---):4.794E-02 U(4.0---):9.977E-03 T(En-Den):3.354E+01 T(Comp):1.436E+05 nuJnu(912A):2.908E+01
Occ(FarIR):2.354E-10 Occ(H n=6):8.468E-14 Occ(lRyd):1.342E-15 Occ(4R):3.707E-17 Occ (Nu-hi):0.000E+00
Tbr(FarIR):3.757E-10 Tbr(H n=6):3.703E-10 Tbr(lRyd):2.120E-10 Tbr(4R):2.351E-11 Tbr (Nu-hi):0.000E+00

```

This section gives a synopsis of the incident continuum, evaluated at the illuminated face of the cloud. The first line gives the number of numerical frequency

cells in the continuum, followed by the energy (in Ryd) of the hydrogen-ionizing continuum² with the largest flux density per unit energy interval (f_v). Next are the energies of the low and high-energy limit of the continuum, both in Ryd and cm or MeV. The last two numbers are the energies of the high-energy limit of the present version of the code, in Ryd and keV.

The second line gives the log of the energy ($\text{erg s}^{-1} \text{ cm}^{-2}$ or erg s^{-1} , depending on whether a flux or luminosity was specified) in the hydrogen ionizing continuum ($1 \text{ Ryd} \leq h\nu < 100 \text{ MeV}$), and the average energy of the hydrogen ionizing continuum, in Ryd, weighted by photon number;

$$\langle h\nu \rangle = \frac{\int_{1 \text{ Ryd}}^{\infty} 4\pi J_\nu d\nu}{\int_{1 \text{ Ryd}}^{\infty} 4\pi J_\nu / h\nu d\nu} . \quad (400)$$

The log of the energy in the X-ray continuum ($20.6 \text{ Ryd} \leq h\nu \leq 7676 \text{ Ryd}$), the log of the energy ($\text{erg s}^{-1} \text{ cm}^{-2}$ or erg s^{-1}), and the number of photons ($\text{cm}^{-2} \text{ s}^{-1}$ or s^{-1}) in the Balmer continuum (0.25 Ryd to 1.0 Ryd) is then printed.

The third line gives the log of the number of photons ($\text{cm}^{-2} \text{ s}^{-1}$ or s^{-1}) in four frequency bins ($1.0 \text{ Ryd} \leq h\nu < 1.807 \text{ Ryd}$, $1.807 \text{ Ryd} \leq h\nu < 4.0 \text{ Ryd}$, $4.0 \text{ Ryd} \leq h\nu < 20.6 \text{ Ryd}$, and $20.6 \text{ Ryd} \leq h\nu < 7676 \text{ Ryd}$). The last number “Ion pht flx” is the flux of hydrogen ionizing photons;

$$\Phi(H) = \frac{Q(H)}{4\pi r^2} \text{ cm}^{-2} \text{ s}^{-1} . \quad (401)$$

In this equation $Q(H)$ is the total number of hydrogen-ionizing photons emitted by the central object (s^{-1}), and r is the separation between the center of the central object and the illuminated face of the cloud. Unlike the majority of the quantities printed in the header, $\Phi(H)$ (per unit area) is always printed, never $Q(H)$ (into $4\pi \text{ sr}$).

The fourth line of the header gives some information about the low and high energy portions of the incident continuum. The first number is the log of the luminosity in the gamma-ray (100 keV \sim to \sim 100 MeV) continuum. The second number on the line is the log of the number of photons over this energy range. The third number is the log of the luminosity in the continuum between 0.25 Ryd and the lowest energy considered, presently an energy of $1.001 \times 10^{-5} \text{ Ryd}$. All of these entries are either per unit area, or radiated into $4\pi \text{ sr}$, depending on how the continuum was specified.

The next entry, “Alf(ox)”, is the spectral index α_{ox} , defined as in Zamorani et al. (1981), except for the difference in sign convention. This is the spectral index which would describe the continuum between 2 keV (147 Ryd) and 2500\AA (0.3645 Ryd) if the continuum could be described as a single power-law, that is,

² The printed number was incorrect in versions 80 through 88.01, but had no other effects on computed results.

$$\frac{f_{\nu}(2 \text{ keV})}{f_{\nu}(2500 \text{ \AA})} = \left(\frac{\nu_{2 \text{ keV}}}{\nu_{2500 \text{ \AA}}} \right)^{\alpha} = 403.3^{\alpha} . \quad (402)$$

The definition of α_{ox} used here is slightly different from that of Zamorani et al. since implicit negative signs are *never* used by Cloudy. Typical AGN have $\alpha_{ox} \sim -1.4$. If no X-rays are present then $\alpha_{ox} = 0$.

The last number on the line is the log of the total energy in the continuum between 1.001×10^{-8} Ryd and 100 MeV. It is given as either $\text{erg cm}^{-2} \text{ s}^{-1}$ or erg s^{-1} , depending on how the continuum was defined. If the continuum is specified per unit area, then this number is 4π times the integrated intensity of the incident continuum. If it is specified as the total luminosity radiated into $4\pi \text{ sr}$, then the quantity is the luminosity.

The next line is optional, depending on whether the continuum is specified as the total luminosity or photon number radiated into $4\pi \text{ sr}$, or as an incident surface flux. If the continuum is specified in absolute terms, i.e., the luminosity or photon number radiated into $4\pi \text{ sr}$, then this optional line is generated. The first quantity is the log of the total luminosity in the continuum, in solar units. The absolute bolometric magnitude, absolute V magnitude, and the bolometric correction, are then given, followed by the log of the continuum specific luminosity (νF_{ν}) at H β (the units of νF_{ν} (H β) are erg s^{-1}).

The next line begins with two ionization parameters. The first is the dimensionless ratio of ionizing photon to hydrogen densities, defined as

$$U \equiv \frac{\Phi(H)}{n_H c} \quad (403)$$

where n_H is the total hydrogen density. The second number is defined in a similar way, but the numerator is the number of photons with energies greater than 4 Ryd (i.e., helium-ionizing). The third number is the equivalent black body temperature corresponding to the energy density u at the illuminated face of the cloud, from the incident continuum and Stefan's radiation density constant a ; $T_u \equiv (L / 4\pi r^2 a c)^{1/4}$, and the next quantity is the Compton temperature of the incident radiation field³. The last number on the line is $4\pi \nu J_{\nu}(912 \text{ \AA})$, the flux at 912 Å ($\text{erg cm}^{-2} \text{ s}^{-1}$). In this equation J_{ν} is the mean intensity of the incident continuum as defined by Mihalas (1978).

The next two lines give some of the incident continuum photon occupation numbers $\eta(\nu)$, defined as

³For a blackbody radiation field $T_{Compton}$ is roughly 4% lower than the blackbody color temperature T_{color} when the energy density temperature T_u is $\ll T_{color}$. Only when $T_u \equiv T_{color}$ does induced Compton heating cause $T_{Compton} \equiv T_{color}$. If $T_u > T_{color}$ then $T_{Compton} > T_{color}$ because of induced Compton heating. All of the relevant physics is included in the Compton temperature printed here.

$$\eta(\nu) \equiv J_\nu(\nu) \left(\frac{2h\nu^3}{c^2} \right)^{-1}, \quad (404)$$

and the incident continuum brightness temperature $T_b(\nu)$, (K), defined as

$$T_b(\nu) \equiv J_\nu(\nu) \left(\frac{2k\nu^2}{c^2} \right)^{-1}, \quad (405)$$

for five energies. These energies correspond to the lowest frequency considered (presently an energy of 1.001×10^{-5} Ryd); the ionization potential of the $n=6$ level of hydrogen ($1/36$ Ryd); an energy of one Rydberg; four Rydbergs, and the high energy limit of the incident continuum (this depends on the continuum shape; the energy is given by the fifth number on the first line of the continuum output).

3.3 Zone Results

```
#### 1 Te:1.877E+04 Hden:3.000E+03 Ne:3.612E+03 R:1.001E+17 R-R0:1.251E+14 dR:2.502E+14 NTR: 2 Htot:7.223E-17 T912: 9.990E+07##
Hydrogen 2.117E-04 9.998E-01 H+0/Hden: 1.000E+00 7.161E-12 H- H2 4.586E-19 4.263E-13 H2+ HeH+ 1.980E-14 H col den 7.505E+17
H 2SP 3-6 2.034E-11 1.923E-17 5.607E-18 8.113E-18 1.204E-17 1.715E-17 Texc(La); 3.800E+03 T(contn): 3.348E+01 T(diffs): 4.329E+00
Helium 6.480E-07 6.538E-03 9.935E-01 He I 2S3 1.113E-08 Comp H, C 6.680E-23 8.732E-24 Fill Fact 1.000E+00 Gam 1/tot 9.651E-01
He singlet 6.368E-07 1.483E-13 4.152E-21 6.702E-21 1.141E-20 1.767E-20 He triplt 1.113E-08 2.297E-18 1.920E-20 5.899E-20 1.365E-20
HeII 6.538E-03 8.905E-12 2.598E-18 3.851E-18 6.052E-18 9.068E-18 1.843E-17 2.477E-17 3.015E-17 PRAD/GAS; 1.384E-02
Pressure NgasTgas; 1.298E+08 P(total): 1.817E-08 P( gas ): 1.792E-08 P(Radtn): 2.480E-10 Rad accel 5.561E-06 Force Mul 1.720E+03
Lithium 1.458E-05 2.598E-02 5.548E-01 4.192E-01 Beryllium 3.617E-06 8.883E-03 9.512E-01 3.987E-02 1.817E-05
Boron 3.651E-08 2.990E-04 4.829E-02 9.513E-01 6.680E-05 0.000E+00
Carbon 8.791E-10 9.441E-06 3.465E-03 7.871E-02 9.178E-01 0.000E+00 0.000E+00 H2O+/Otot 0.000E+00 OH+/Otot1 0.000E+00
Nitrogen 3.321E-10 5.934E-06 5.302E-03 1.590E-01 4.493E-01 3.864E-01 0.000E+00 0.000E+00 O2/Otot1: 0.000E+00 O2+/Otot: 0.000E+00
Oxygen 1.292E-10 2.793E-06 3.943E-03 1.959E-01 6.207E-01 1.732E-01 6.220E-03 0.000E+00 0.000E+00 Hex(tot): 0.000E+00 A:-12.580
Fluorine 0.000E+00 1.161E-06 3.523E-03 1.387E-01 7.343E-01 1.233E-01 1.532E-04 0.000E+00 0.000E+00 0.000E+00
Neon 0.000E+00 1.880E-06 4.552E-03 1.833E-01 7.138E-01 9.818E-02 1.582E-04 0.000E+00 0.000E+00 0.000E+00
Sodium 0.000E+00 9.429E-06 6.092E-03 2.091E-01 5.649E-01 2.170E-01 2.914E-03 1.103E-06 0.000E+00 0.000E+00 0.000E+00
Magnesium 0 3.071E-08 1.107E-05 1.429E-02 2.867E-01 4.640E-01 2.331E-01 1.892E-03 0.000E+00 0.000E+00 0.000E+00
Aluminum 0 4.196E-07 1.751E-04 2.929E-03 4.417E-01 4.339E-01 1.202E-01 1.063E-03 0.000E+00 0.000E+00 0.000E+00
Silicon 0 3.897E-08 7.822E-05 5.003E-03 4.523E-02 8.506E-01 9.869E-02 4.456E-04 0.000E+00 0.000E+00 0.000E+00
Phosphorus 0 2.860E-09 1.391E-05 2.403E-03 5.739E-02 2.383E-01 7.003E-01 1.570E-03 0.000E+00 0.000E+00 0.000E+00
Sulphur 0 1.579E-10 2.734E-06 8.478E-04 2.012E-02 1.396E-01 4.764E-01 3.630E-01 5.985E-06 0.000E+00 0.000E+00 0.000E+00
Chlorine 0 1.150E-10 2.004E-06 4.807E-04 1.203E-02 8.264E-02 5.192E-01 3.535E-01 3.216E-02 0.000E+00 0.000E+00 0.000E+00
Argon 0 0.000E+00 8.222E-08 1.560E-04 3.905E-03 2.841E-02 3.441E-01 5.810E-01 4.203E-02 4.291E-04 0.000E+00 0.000E+00
Potassium 0 0.000E+00 7.488E-08 2.094E-04 5.439E-03 4.070E-02 3.727E-01 4.913E-01 8.887E-02 7.459E-04 0.000E+00 0.000E+00
Calcium 0 0.000E+00 7.726E-08 2.146E-04 8.651E-03 3.813E-02 3.786E-01 4.789E-01 9.362E-02 1.864E-03 0.000E+00 0.000E+00
Scandium 0 4.538E-10 2.288E-07 1.997E-04 2.205E-02 1.020E-01 4.865E-01 3.674E-01 2.163E-02 1.787E-04 0.000E+00 0.000E+00
Titanium 0 0.000E+00 2.834E-08 3.943E-05 6.968E-03 1.732E-01 5.239E-01 2.769E-01 1.888E-02 9.026E-05 0.000E+00 0.000E+00
Vanadium 0 0.000E+00 1.782E-09 4.294E-06 9.849E-04 3.606E-02 7.237E-01 2.292E-01 1.004E-02 4.626E-05 0.000E+00 0.000E+00
Chromium 0 0.000E+00 2.975E-10 7.639E-07 2.002E-04 9.541E-03 9.232E-01 6.817E-01 1.620E-02 4.507E-05 0.000E+00 0.000E+00
Manganese 0 0.000E+00 2.018E-10 3.969E-07 1.074E-04 5.537E-03 2.073E-01 6.636E-01 1.233E-01 1.921E-04 0.000E+00 0.000E+00
Iron 0 0.000E+00 4.846E-10 1.340E-07 4.709E-05 2.649E-03 1.167E-01 6.324E-01 2.416E-01 6.693E-03 0.000E+00 0.000E+00
Cobalt 0 0.000E+00 0.000E+00 1.246E-08 4.157E-06 3.247E-04 4.726E-02 5.588E-01 3.825E-01 1.111E-02 7.686E-06 0.000E+00 0.000E+00
Nickel 0 0.000E+00 0.000E+00 1.191E-08 3.872E-06 2.691E-04 5.398E-02 5.317E-01 3.988E-01 1.528E-02 1.410E-05 0.000E+00 0.000E+00
Copper 0 0.000E+00 0.000E+00 9.451E-09 2.963E-06 2.062E-04 3.519E-02 5.963E-01 3.541E-01 1.420E-02 1.296E-05 0.000E+00 0.000E+00
Zinc 0 0.000E+00 0.000E+00 1.369E-08 3.726E-06 2.355E-04 3.976E-02 5.138E-01 4.362E-01 9.973E-03 1.011E-05 0.000E+00 0.000E+00
```

The results of calculations for the first and last zones are always printed. Results for intermediate zones can be printed if desired (see the **print every** command). The following is a line-by-line description of the output produced for each printed zone.

3.3.1 Line 1

The line begins with a series of # characters, to make it easy to locate with an editor. The zone number is the first number, followed by the electron temperature of the zone (“Te”). A lower case u will appear before the “Te” if the temperature solution is possibly thermally unstable (i.e., the derivative of the net cooling with respect to temperature is negative. See the section in Part III on thermal stability problems). The total hydrogen (“Hden”) and electron (“Ne”) densities (cm^{-3}) follow. The next number (“R”) is the distance to the center of the zone, from the center of the central object. The depth, the distance between the illuminated face of the cloud and the center of the zone, (“R-R0”, or $r-r_0$), and the thickness of the zone (“dR”, or δr), (all are in cm), follow. The inner edge of the zone is $(r-r_0) - \delta r/2$ from the illuminated face of the cloud. The line ends with a number indicating how many ionization iterations were needed for this zone to converge (NTR), followed by the total

heating⁴ (“Htot”; photoelectric and otherwise, $\text{erg cm}^{-3} \text{s}^{-1}$), and the optical depth between the *illuminated* face of the cloud and the *outer* edge of the zone at the Lyman limit (T912; the number is the *total absorption* optical depth at 912Å, and *not* the hydrogen Lyman limit optical depth).

3.3.2 [Optional] wind parameters

A line describing the velocity and acceleration of the zone is printed if the **wind** option is used. The numbers are the wind velocity at the outer edge of the current zone (km s^{-1}), inward gravitational acceleration (cm s^{-2}), total outward radiative acceleration, and the fraction of this acceleration caused by the incident continuum, line driving, and the gradient of the radiation pressure.

3.3.3 [Optional] radiation pressure

If the ratio of line radiation to gas pressure, $P(\text{radiation})/P(\text{gas})$, is greater than 5%, then a line describing the source of the radiation pressure is generated. The line begins with the label **P(Lines)** and continues with the fraction of the total radiation pressure produced that an emission line, the spectroscopic designation of the line, followed by its wavelength in Ångstroms. Up to twenty lines can be printed, although in most cases only $\text{L}\alpha$ and a few others dominate.

3.3.4 Line 1 - Hydrogen I

The line begins with the abundance of neutral and ionized hydrogen relative to all atomic-ionic hydrogen (i.e., the ratios $H^0/(H^0+H^+)$ and $H^+/(H^0+H^+)$ where H^0 is the population in all bound levels of hydrogen. If **print departure coefficients** has been specified then departure coefficients are also printed on the following line. Neutral hydrogen H^0 is defined to be the total population of atomic hydrogen in all explicitly computed bound levels. Next comes $H+o/H_{\text{den}}$, the ratio of the density of hydrogen in atomic or ionic form (this is indicated by the label “H+o”) to the total hydrogen density in all forms (including molecular).

The following five numbers are abundances of the negative hydrogen ion and several molecules (H^- , H_2 , H_2^+ , and HeH^+) relative to the total hydrogen abundance. The total hydrogen density is usually referred to by the label **hden**, and is the sum $H^0 + H^+ + \text{H}^- + 2\text{H}_2 + 2\text{H}_2^+ + 3\text{H}_3^+$. Note that, with this definition of the hydrogen density a fully molecular gas will have $n(\text{H}_2)/n(\text{H})=0.5$. These molecular abundances are also expressed as departure coefficients if this option is set with the **print departure coefficients** command. The last number on the line is the total hydrogen column density (cm^{-2}).

3.3.5 Line 2 - Hydrogen II

The first two numbers are the populations of the $\text{H}^0 2s$ and $2p$ levels relative to the ionized hydrogen density. The next four numbers are populations of levels 3 to 6, again relative to the ionized hydrogen density. All of these populations usually are relative to the ionized hydrogen density, but can also be printed as LTE departure

⁴Cloudy defines heating as the energy input by the freed photoelectron, or $h\nu - \text{IP}$, where IP is the ionization potential of the atom or ion, and $h\nu$ is the energy of the photon. See Osterbrock (1989) for more details.

coefficients if the `print departure coefficients` command is given. The excitation temperature T_{exc} of $\text{Ly}\alpha$, defined as

$$\frac{n(2p)/g(2p)}{n(1s)/g(1s)} = \exp(-h\nu/kT_{exc}) \quad (406)$$

is given. This is followed by the temperature corresponding to the energy density of the attenuated incident continuum (“T(contn)”), and the diffuse continua (“T(diffs)”). This includes all trapped lines and diffuse continuous emission.

3.3.6 Line 3 - Helium

The first three numbers are the total populations of the three ionization stages of helium, relative to the total helium abundance. The population of atomic helium is the sum of the total population in the triplets and singlets, including the population of all explicitly computed levels of each. These populations can also be expressed as departure coefficients if this option is set with the `print departure coefficients` command. The population of $\text{He } 2^3\text{S}$, relative to the total helium abundance, follows. The Compton heating and cooling rates (both $\text{erg cm}^{-3} \text{ s}^{-1}$) are next, followed by the gas filling factor. The last number is the fraction of the total hydrogen ionizations that are caused by photoionization from the ground state.

3.3.7 Line 4 - Atomic Helium

The first group are the level populations of the populations of the $n=1$ to 6 levels of the He^0 singlets. Level two is actually resolved into $2s$ and $2p$, but the total population of 2 is printed. The next group consists of populations of the $2s$, $2p$, and $n=3s, p, d$ levels of the He^0 triplets. Both sets of populations are relative to the total helium abundance. Departure coefficients are also printed if requested.

3.3.8 Line 5 - Ionized Helium

The populations of the $2s$, $2p$, and $n=3$ to 6 levels are indicated. There are relative to He^{++} ; departure coefficients are also printed if requested. The ratio of radiation pressure to gas pressure follows.

3.3.9 Optional Grains

If grains are present, then lines giving some properties of the grain populations are printed. Each line gives the results of calculations for a specific type of grain. Normally, a type of graphite and silicate are included when grains are present. There will be one line of output for each grain species. Each line begins with the name of the grain, and an asterisk appears if the species is quantum heated. The remainder of the lines gives the equilibrium temperature of the grain, the potential in volts, the charge, the drift velocity, followed by the gas heating ($\text{erg cm}^{-3} \text{ s}^{-1}$) due to grain photoionization, and the dimensionless fraction of the total gas heating due to grain photoionization. For quantum heated grains the temperature is the average weighted by T^4 .

3.3.10 Pressure

Some information concerning the pressure is printed. The gas equation of state includes thermal gas pressure, the radiation pressure due to trapped line emission, and the radiation pressure due to absorption of the incident continuum. The first

number is the gas pressure $n_{\text{gas}} T_{\text{gas}}$ (with units $\text{cm}^{-3} \text{K}$), followed by the total pressure (dynes cm^{-2}), and followed by the gas pressure ($n_{\text{gas}} kT_{\text{gas}}$) in dynes cm^{-2} . The radiation pressure follows. The second to last number is the radiative acceleration (cm s^{-2}) at the inner edge of this zone. The radiative acceleration is computed with all continuous scattering and absorption opacities included. The last number is a force multiplier, defined as in Tarter and McKee (1973), and is the ratio of total opacity to electron scattering opacity.

3.3.11 Molecules

A line giving relative abundances of some molecules is printed if there is a significant molecular fraction. All molecular abundances are relative to either the total carbon or total oxygen abundance (this is indicated in the label for each). In order, the molecules are CH, CH⁺, CO, CO⁺, H₂O, and OH.

3.3.12 Li, Be, B

Abundances of each stage of ionization relative to the total gas phase abundance of the element are printed across two lines.

3.3.13 Carbon

The abundances of the seven stages of ionization of carbon relative to the total carbon abundance begin the line. The relative abundance of H₂O⁺ and OH⁺ (relative to the total oxygen abundance) follows.

3.3.14 Nitrogen

The relative populations of the eight ionization stages of nitrogen are printed first. The relative abundance of O₂ and O₂⁺ (relative to the total oxygen abundance) follows.

3.3.15 Oxygen

The oxygen ionization stages are followed by the extra heat added at this zone (erg $\text{cm}^{-3} \text{s}^{-1}$); due to cosmic rays, turbulence, etc, and the log of the effective hydrogen recombination coefficient ($\text{cm}^3 \text{s}^{-1}$).

3.3.16 Fluorine, Neon

The fluorine and neon relative ionization balances are printed across the line.

3.3.17 Remaining Elements

There are too many ionization stages to print across the line. Although all stages with non-trivial abundances are computed, only the highest twelve stages of ionization are printed. The first number is an integer indicating how many stages are “off the page to the left”. If the number is 2, then the first printed stage of ionization is twice ionized, i.e., Fe⁺².

3.4 Calculation Stopped Because ...

```
Calculation stopped because lowest Te reached.      Iteration 1 of 1
The geometry is spherical.
!Non-collisional excitation of [OIII] 4363 reached 2.08% of the total.
!AGE: Cloud age was not set. I cannot check whether the time-steady assumption is ok.
Derivative of net cooling negative and so possibly thermally unstable in 4 zones.
Photoionization of He 2TriS reached 17.1% of the total rate out, 10.6% of that was Lyα.
Grains were not present but might survive in this environment (energy density temperature was 3.35E+01K)
The ratio of radiation to gas pressure reached 1.65E+01. Caused by Lyman alpha.
Line radiation pressure capped by thermalization length.
```

A series of messages appear after the printout of the last zone.

The first will say why the calculation stopped. In a valid calculation the model will stop because one of the specified stopping criteria specified was met. If no other criteria are specified, then the calculation usually stops when the default lowest temperature of 4000 K is reached. If the code stops because of an unintended reason (i.e., internal errors, or the default limit to the number of zones) then a warning is printed saying that the calculation may have halted prematurely.

Only one stopping criterion message will be printed. The possible messages, and their interpretations, are:

3.4.1 ... because of radiation pressure

The default density law is for a constant density. If constant pressure is specified instead (with the **constant pressure** command), then Cloudy will try to keep the total pressure, particle and radiation, constant. The radiation pressure is small at the boundaries of the cloud, so the cloud will be unstable if the ratio of radiation to total pressure exceeds 0.5. The calculation stops, and this message is generated, if this occurs after the first iteration.

3.4.2 ... because lowest EDEN reached.

The calculation can be forced to stop when the electron density (**eden**) falls below a certain value, as set by the **stop eden** command. This can be used to stop the calculation at an ionization front. The default lowest electron density is negative, so this stopping criterion applies only when the command is entered.

3.4.3 ... because low electron fraction.

The calculation can be forced to stop when the ratio of electron to hydrogen densities falls below a certain value, as set by the **stop efrac** command. This can be used to stop the calculation at an ionization front when the hydrogen density there is not known (for instance, in a constant pressure model). The default lowest electron density is negative, so this stopping criterion applies only when the command is entered.

3.4.4 ... because wind veloc too small

The code can perform a wind calculation which includes the outward force due to radiation pressure and the inward force of gravity. The solution is only valid well above the sonic point. This message is printed if the gas is decelerated to below the sonic point.

3.4.5 ... because code returned BUSTED

The calculation stopped because something bad happened. The results are suspect. I would appreciate learning about this - please send the input script and version number.

3.4.6 ... because DRAD small - set DRMIN

The Strömgren radius of the H⁺ zone is estimated at the start of the calculation, and the smallest allowed zone thickness is then set as a very small fraction of this. The calculation will stop if the zone thickness falls below this smallest thickness. This can occur because of any of several logical errors within Cloudy (adaptive logic is used to continuously adjust the zone thickness), although it can rarely occur for physical reasons as well. The smallest thickness can be reset to any number with the

`set drmin` command, but it should not be necessary to do this. I would appreciate learning about this - please send the input script and version number.

3.4.7 ... because DR small rel to thick.

The depth into the cloud is stored as the double precision variable *depth* and the zone thickness is stored as the double precision variable *drad*. If the zone size becomes too small relative to the depth ($drad/depth < 10^{-14}$) then the depth variable will underflow such that $depth + drad = depth$. The calculation will stop in this case and give the above reason if this problem prevents the density from being properly evaluated. This is a fundamental numerical problem with no clear solution.

3.4.8 ... because carbon fully molecular.

For mixtures where oxygen is more abundant than carbon the atomic carbon abundance can become vanishingly small when carbon monoxide forms. The matrix inversion routine may have trouble determining the carbon balance under these conditions. As a precaution the current version of the code will stop if the ratio of carbon monoxide to total gas phase carbon exceeds 0.80, the value of the code variable *colimt*. This limit can be reset with the `set colimt` command.

3.4.9 ... because negative mole abundan.

The matrix inversion routine can predict negative abundances of some of the heavy element molecules when the gas becomes predominantly molecular. Cloudy is not now designed to handle this situation, but should be well protected against this happening. I would appreciate learning about this occurring- please send the input script and version number.

3.4.10 ... because optical depth reached.

The default value of the largest allowed continuous optical depth is unphysically large, and can be reset with the `stop optical depth` command. The command specifies both the optical depth, and the energy at which it is to be evaluated. All absorption opacity sources included in the calculation contribute to the computed optical depths. If the calculation stops because the largest continuum optical depth is reached, then this line is printed. This line is also printed if the `stop effective column density` command is used to stop the calculation, since this command is actually a form of the `stop optical depth` command.

3.4.11 ... because outer radius reached.

The default outer radius is unphysically large, but can be changed with the `radius` or `stop thickness` commands. If the calculation stops because the outer radius set by one of these commands is reached, then this line is printed.

3.4.12 ... because column dens reached.

The default values of the largest allowed neutral, ionized, and total hydrogen column densities are unphysically large. They can be reset with the commands `stop column density`, `stop neutral column density`, or `stop ionized column density`. This message will be printed if one of these criteria stops the calculation.

3.4.13 ... because lowest T_e reached.

The default value of the lowest temperature allowed is 4000 K. This is reasonable when only optical emission lines are of interest. The limit can be changed with the `stop temperature` command. This message is printed if the calculation stops because the lowest temperature is reached.

3.4.14 ... because highest T_e reached.

The default value of the highest temperature allowed is 10^{10} K. The limit can be changed with the `stop temperature exceeds` command. This message is printed if the calculation stops because the highest allowed temperature is exceeded.

3.4.15 ... because NZONE reached.

The default condition is for up to 600 zones to be computed. This can be reset with the `stop zone` command. This message is printed if the calculation stops because the limiting number of zones is reached. A warning will be printed at the end of the calculation if it stops because it hits the default limit to the number of zones allowed, presently 600, since this was probably not intended.

The default limit to the number of zones can be increased, while retaining the check that the default limit is not hit, by using the `set nend` command.

3.4.16 ... because line ratio reached.

It is possible to set a limit to the largest value of an emission-line intensity ratio with the `stop line` command. This message is printed if the calculation stops because the largest value of the ratio is reached.

3.4.17 ... because internal error - DRAD.

An internal logical error caused this message to be printed. Send the command lines, and the version number of Cloudy to me. My email address is gary@cloud9.pa.uky.edu.

3.4.18 ... because initial conditions out of bounds.

The temperature of the first zone was not within the temperature bounds of the code. This is probably due to the incident continuum not being set properly.

3.4.19 ... because reason not specified.

This is another internal error I would appreciate learning about.

3.5 Geometry

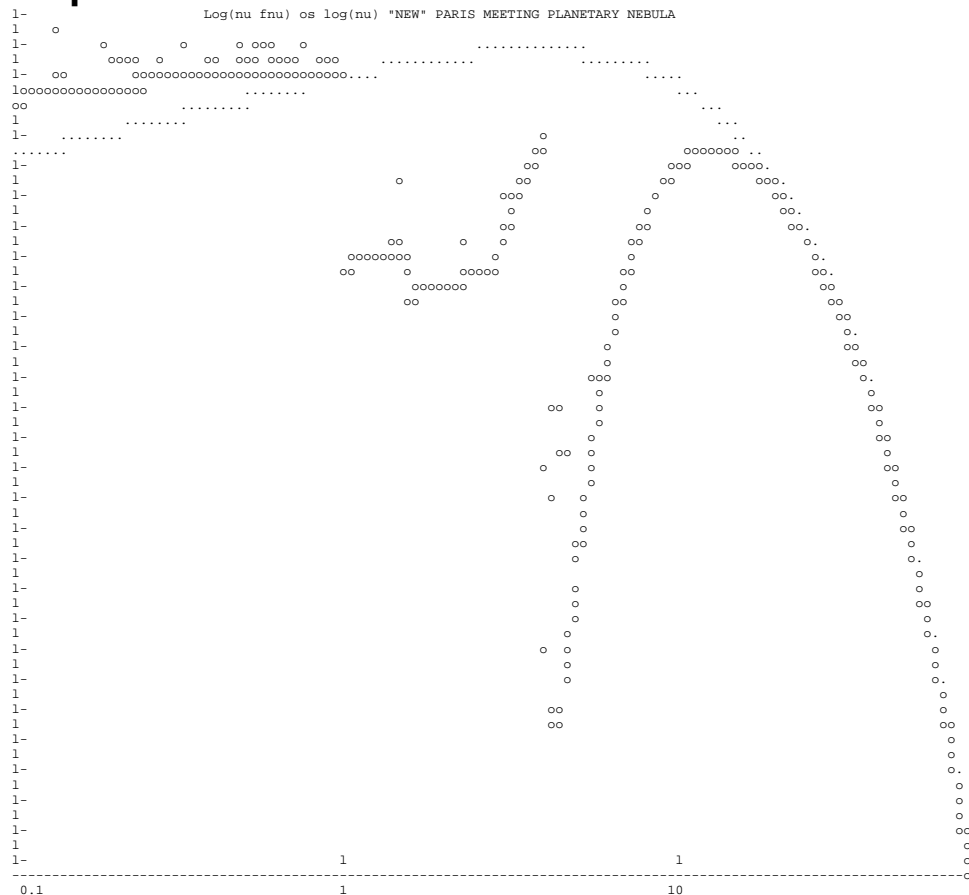
After saying why the calculation stopped, Cloudy will say whether the geometry is plane parallel ($\Delta r/r_o < 0.1$), a thick shell ($\Delta r/r_o < 3$), or spherical ($\Delta r/r_o \geq 3$), where r_o is the inner radius and Δr is the thickness of the cloud.

3.6 Warnings, Cautions, Surprises, and Notes

The next, optional, messages fall into four categories: warnings, which begin with W-; cautions, which begin with C-; surprising results, which begin with an explanation mark (!), and notes.

Cloudy checks that its range of validity was not exceeded in the calculation. Warnings are issued to indicate that the program has not treated an important process correctly. For instance, warnings occur if the temperature was high enough for the electrons to be relativistic, if the global heating - cooling balance is off by more than 20%, or if the code stopped for an unintended reason. I would like to hear about warnings, my e-mail address is gary@cloud9.pa.uky.edu. Cautions are less severe, and indicate that Cloudy is on thin ice. Examples are when the optical depths in excited states of hydrogen change during the last iteration. Surprises begin with “!” and indicate that, while the physical process has been treated correctly, the result is surprising. An example is when induced Compton heating is more than 5 percent of the total Compton heating. Notes indicate interesting features about the model, such as maser effects in lines or continua, or if the fine structure lines are optically thick. The messages are usually self-explanatory.

3.7 Optional Plot



If any of the optional plots are requested with the `plot xxx` command then they will appear next. The quantities plotted are described in the section of HAZY where the plot command is defined.

3.8 Final Printout

3.8.1 Emission-Line Spectrum

```

***** Cloudy 94.01 <*****
* title "New" Paris meeting Planetary nebula
* c recompute "standard" PN model of the Pequignot Meudon Conference
* init file='c84.ini'
* sphere
* black body, T=150,000K radius = 10
* hden = 3.4771213
* radius = 17
* abund -1 C-3.523 N-4. 0-3.222 ne-3.824 na=-10 mg-4.523 al=-10
* plot continuum range .1
* punch overview last 70
* c parispn.in
* c Sun
***** Log(U): -1.32 <*****

Emission Line Spectrum. Constant Density Model. Closed geometry. Iteration 1 of 1.
Luminosity (erg/s) emitted by shell with full coverage.

general properties.....
TOTL 4861 -16.690 1.0000 O 4 1397A -19.618 0.0012 N 3 265A -18.070 0.0417 Al 4 115A -18.802 0.0077
TOTL 1216 -14.856 68.2311 O 4 1407A -19.246 0.0028 N 3 277A -18.136 0.0358 Al 4 157A -18.512 0.0151
TOTL 1216 -14.856 68.2311 O 4 1405A -18.737 0.0090 N 3 312A -18.098 0.0390 Al 4 110A -19.172 0.0033
Inci 0 2.563***** O 4 1401A -18.585 0.0127 N 3 273A -18.314 0.0238 Al 4 108A -19.435 0.0018
TotH 0 -15.017 47.0916 O 4 789A -15.572 13.1211 N 3 267A -18.431 0.0181 Al 4 106A -19.673 0.0010
TotC 0 -15.338 22.4813 rec 789 -16.972 0.5224 N 3 249A -18.474 0.0164 Al 4 123A -19.632 0.0011
BFHl 0 -15.382 20.2978 O 5 630A -15.569 13.1903 N 3 271A -18.468 0.0167 Al 5 108A -19.318 0.0024
BFHe 0 -15.359 21.3951 rec 630 -19.299 0.0025 N 3 306A -18.476 0.0164 Al 5 104A -19.572 0.0013
TotM 0 -15.975 5.1816 TOTL 1218 -18.965 0.0053 N 3 269A -18.595 0.0124 Al 5 104A -19.590 0.0013
HFFc 0 -16.544 1.3992 O 5 1218A -19.336 0.0023 N 3 301A -18.587 0.0127 Al 5 131A -19.586 0.0013
H FF 0 -16.544 1.3992 O 5 1211 -19.206 0.0030 N 3 239A -19.009 0.0048 Al 5 281A -18.948 0.0055
ComH 0 -19.609 0.0012 TOTL 1035 -17.103 0.3864 N 3 234A -19.116 0.0037 Al 6 311A -19.451 0.0017
CT H 0 -17.355 0.2159 O 6 1032A -17.275 0.2595 N 3 287A -19.051 0.0043 Si 2 1180A -17.915 0.0595
Ca B 1216 -15.292 24.9963 O 6 1037A -17.586 0.1269 N 3 231A -19.307 0.0024 Si 2 993A -18.667 0.0105
Q(H) 4861 0.154***** O 6 150A -18.002 0.0487 N 3 279A -19.287 0.0025 Si 2 894A -19.291 0.0025
Q(H) 1216 1.688***** Ne 3 156ml -17.234 0.2857 N 3 230A -19.501 0.0015 Si 3 565A -18.445 0.0176
H-like iso-seq..... Ne 3 360ml -18.377 0.0206 N 3 275A -19.454 0.0017 Si 3 465A -18.550 0.0138
CaBo 4861 -16.830 0.7239 Ne 3 3869 -17.319 0.2346 N 3 272A -19.564 0.0013 Si 3 436A -19.034 0.0045
CaBo 1640 -15.973 5.2094 Ne 3 3968 -17.840 0.0707 N 3 270A -19.622 0.0012 Si 3 407A -19.294 0.0025

```

The final printout begins by reprinting the input commands. The box surrounding it gives both the version number of Cloudy (at the top) and the log of the ionization parameter (the ratio of ionizing photon to hydrogen densities) at the bottom.

The set of printed emission lines also includes other predicted quantities. Some continua, and various indications of contributors to lines and continua, are mixed in what follows. The previous section of the document describes how to convert these into some observed quantities. Not all are printed by default – the **print** commands described in Part I and also in section 5.3.2 starting on page 446 tell how to get more or fewer predictions.

The line following the box summarizes some properties of the model and output. The first part of the line indicates whether the energy in the emission lines is given as the luminosity radiated by a spherical shell covering Ω sr (erg s⁻¹; $\Omega / 4\pi$ is the covering factor) or the intensity produced by a unit area of gas (erg s⁻¹ cm⁻²). Which of the two choices is printed is determined by whether the luminosity of the continuum was specified as the luminosity radiated by the central object into 4π sr or the intensity ($4\pi J$) of the incident continuum (erg cm⁻² s⁻¹) at the illuminated face of the cloud. If the model is spherical and the incident continuum specified per unit area, then the emergent emission-line spectrum will be per unit area in units of the inner radius r_0 (that is, the total line luminosity radiated by a shell covering 4π sr will be the listed intensity $4\pi J$ times $4\pi r_0^2$). The second part of this line indicates the density structure of the model (i.e., wind, constant density, constant pressure, constant gas pressure, power-law density distribution, etc). The next section tells whether the geometry was open or closed. The last part indicates which iteration this is.

The organization and meaning of the different of lines in the printout is discussed on page 440 below.

A list of emission lines with negative intensities may follow the main block of lines. These are lines, which heat rather than cool the gas (heating is negative cooling). This is not a problem, but occurs if the line de-excitation rate exceeds the line excitation rate. The most common reason for this to occur is if the line is radiatively excited but collisionally de-excited.

3.8.2 Thermal balance

```
Cooling:  O 3 5007:0.245  O 3 4959:0.082
Heating:  BPH1  0:0.720  BFHe  0:0.233
```

```
IONIZE PARMET:  U(1-) -1.3193  U(4-) -2.0010  U(sp): -2.51  Q(ion):  43.458  L(ion):  33.712  Q(low): 49.69  P(low)  37.341
ENERGY BUDGET:  Heat:  37.222  Coolg:  37.222  Error:  0.2%  Rec Lin:  37.064  WorkF:  37.437  F-F  H 21.885  RadBetaMax:1.65E+01
```

Cooling: This line indicates the fraction of the total cooling (defined here as in Osterbrock 1989; that is, the energy of the freed photoelectron) carried by the indicated emission lines. The designation of the line is given as in the emission-line spectrum, and this is followed by the ratio of the energy in the line to the total cooling. This is an important indication of the fundamental power-losses governing conditions in the model. The labels used are the same as those in the line array.

Heating: This line indicates the fraction of the total heating produced by various processes. The labels used are the same as those in the line array.

IONIZE PARMET The line begins with the log of the H “U(1-)” and He+ “U(4-)” ionization parameters defined in the header. The third number “U(sp)” is the log of a spherical ionization parameter often used in spherical geometries, such as H II regions or planetary nebulae. It is defined as

$$U_{sph} = \frac{Q(H)}{4\pi R_s^2 n_H c} \quad (407)$$

where R_s is the Strömgren radius, defined as the point where the hydrogen neutral fraction falls to $H^0/H_{tot} = 0.5$. If no ionization front is present, then U_{sph} is evaluated at the outer edge of the computed structure. The next two numbers are the log of the number of hydrogen ionizing photons ($h\nu \geq 1$ Ryd) exiting the nebula “Q(ion)”, and the log of the energy in this ionizing continuum “L(ion)”. The last two numbers are the equivalent quantities, for non-ionizing ($h\nu < 1$ Ryd) radiation. These are either per unit area or by a shell covering 4π sr. These have been corrected for the r^{-2} dilution if per unit area, and so are directly comparable with the numbers given at the start of the calculation.

ENERGY BUDGET This line gives an indication of the energy budget of the nebula. The first number “Heat” is the log of the total heating (in ergs s⁻¹, but again either into 4π sr or cm⁻²). The second number “Coolg” is the log of the total cooling, in the same units. Cooling, as defined in Osterbrock (1989), is the total energy in collisionally excited lines and part of the recombination energy, but *does not* include recombination lines. The percentage error in the heating-cooling match “Error” follows. The next number “Rec Lin” is the log of the total luminosity in recombination lines. The number indicated by “WorkF” is an indication of the work function (that is, the log of the energy needed to remove bound electrons from the atom or ion) of the cloud. The work function and the total cooling do not add up to the total energy absorbed from the incident continuum because some recombination lines of helium and heavy elements contribute to both. The next number “F-F H” is

the log of the amount of energy deposited by free-free heating, and the last number “RadBetaMax” is the largest value of the ratio of radiation to gas pressures which occurred in the calculation.

3.8.3 Column densities, et al.

```
Column density  H12:9.024E+20  H  II:8.704E+20  HI:3.204E+19  H-:  1.445E+12  H2:  9.379E+11  H2+:2.190E+11  He  H+:4.187E+12
                CH:0.000E+00  CH+:0.000E+00  OH:0.000E+00  OH+:  0.000E+00  O2:  0.000E+00  C2:0.000E+00  CO:0.000E+00
                CO+:0.000E+00  H2O:0.000E+00H2O+:0.000E+00  O2+:  0.000E+00  C2+:  0.000E+00  H3+:0.000E+00  H3O+:0.000E+00
                CH2+:0.000E+00  CH2:0.000E+00  CH3:0.000E+00MoH2O  0.000E+00
                Col(Heff):  6.286E+20  snd travl time  2.34E+11 sec  NeN+dl: 2.98E+24  Te-low:3.81E+03  Te-hi:1.88E+04
He/Ha:9.82E-02  =  0.98*true  Lthin:1.00E+30  lter/zn:2.016  Hlu/zn:6.56E+00
<a>:0.00E+00  erdeFe0.0E+00  Tcompt8.11E+07  Tthr6.76E+09  <Tden>: 1.33E+04  <dens>:7.11E-21  <Mol>:6.47E-01
Mean Jeans  1(cm)7.67E+19  M(sun)8.46E+05  smallest:  len(cm):4.10E+19  M(sun):1.29E+05  Alf(ox-tran):  0.0000
```

Column density This line lists the column densities (cm^{-2}) of some ions and molecules. The first number “H12” is the total hydrogen column density (both H^0 and H^+). The following two numbers are the column densities in H^+ and H^0 only. The last four numbers are column densities in four ion - molecules (H^- , H_2 , H_2^+ , and HeH^+).

The next series of three lines give column densities in various molecules.

Col (Heff) The effective column density “Col(Heff)”, as defined in the section on the **stop effective column density** command, is printed. This is followed by “snd travl time”, the sound travel time across the nebula in seconds. Constant pressure is only valid if the cloud is static for times considerably longer than this. The third number “NeN+dl” is the emission measure, the integral over radius of the product $n_e n_p f(r) dr$, where $f(r)$ is the filling factor. The last two numbers are the lowest “Te-low” and highest “Te-hi” electron temperatures found in the computed structure.

He/Ha This line gives some quantities deduced from the predicted emission-line spectrum. The first (He/Ha) number is the apparent helium abundance He/H , measured from the emission-line intensities using techniques similar to those described in Osterbrock (1989);

$$\left(\frac{\text{He}}{\text{H}} \right)_{\text{apparent}} = \frac{0.739 \times I(5876) + 0.078 \times I(4686)}{I(H\beta)} . \quad (408)$$

The intensity of both $\text{H}\beta$ and $\text{HeI } \lambda 5876$ are the total predicted intensities, and includes contributions from collisional excitation and radiative transfer effects. The intensity of $\text{He II } \lambda 4686$ is taken from Case B results, which are better than those of the model atom at low densities. The second number (i.e., $1.07 \times \text{true}$), is the ratio of this deduced abundance to the true abundance. This provides a simple way to check whether ionization correction factors, or other effects, would upset the measurement of the helium abundance of the model nebula. This is followed by the longest wavelength in centimeters “Lthin” at which the nebula is optically thin. Generally the largest FIR opacity source is brems, and the number will be 10^{30} if the nebula is optically thin across the IR. The last two quantities are related to the average number of iterations needed to converge each zone.

<a> The mean radiative acceleration (cm s^{-2}) is printed if the geometry is a wind model and zero otherwise. This is followed by some time scales. The first “erdeFe” is the time scale, in seconds, to photoerode Fe (Boyd and Ferland 1987; this number is 0s if the γ -ray flux is zero). The next two are the Compton equilibrium timescale “Tcompt”, and the thermal cooling timescale “Tthr”. Both are in seconds. The

density (gm cm^{-3}) weighted mean temperature “<Tden>”, radius weighted mean density “<dens>” (gm cm^{-3}), and mean molecular weight “<Mol>” follow.

Mean Jeans This line gives the mean Jeans length “l(cm)” (cm) and Jeans mass “M(sun)” (in solar units), followed by the smallest Jeans length “smallest len(cm)” and the smallest Jeans mass “M(sun)” which occurred in the calculation. The last quantity “Alf(ox-tran)” is the spectral index α_{ox} , defined as in the header, but for the transmitted continuum (attenuated incident continuum plus emitted continuum produced by the cloud).

H and He atoms , This line gives the number of levels of the model hydrogen atom, the “topoff” level, above which the remainder of the recombination coefficient is added, the type of topping off used for this calculation, and the number of levels used for the helium singlets and ion. The last number on the line is the execution time in seconds.

3.8.4 Averaged Quantities

Averaged Quantities										
	Te	Te(Ne)	Te(NeNp)	Te(NeHe+)	Te(NeHe2+)	Te(NeO+)	Te(NeO2+)	NH	Ne(O2+)	Ne(Np)
Radius:	1.334E+04	1.358E+04	1.363E+04	1.109E+04	1.542E+04	1.115E+04	1.139E+04	3.000E+03	3.339E+03	3.429E+03
Volume:	1.179E+04	1.207E+04	1.214E+04	1.100E+04	1.446E+04	1.115E+04	1.117E+04	3.000E+03	3.305E+03	3.294E+03

Peimbert T(OIIIr)1.15E+04 T(Bac)1.21E+04 T(Hth)1.14E+04 t2(Hstro) 2.72E-02 T(O3-BAC)1.15E+04 t2(O3-BC) 1.81E-03 t2(O3str) 7.36E-03

This begins with several temperature and density averages, over either radius or volume. The volume averages are only printed if the **sphere** command is entered. The quantity which is printed is indicated at the top of each column. The averaged quantity is the first part of the label, and the weighting used is indicated by the quantity in parenthesis. For instance, **Te(NeO2+)** is the electron temperature averaged with respect to the product of the electron and O^{2+} densities.

Peimbert This series of quantities deal with temperature fluctuations (t^2 , Peimbert 1967) . The code attempts to analyze the predicted emission line and continuum spectrum using the same steps that Manuel outlined in this paper. The code does not attempt to correct the predicted emission line intensities for collisional suppression or reddening, so this line is only printed if the density is below the density set with the **set tsqden** command - the default is 10^7 cm^{-3} . This code does not attempt to deredden the spectrum: a caution is printed if grains are present.

The nature of temperature fluctuations is, in my option, the biggest open question in nebular astrophysics. Theory (Cloudy too) predicts that they should be very small, because of the steep dependence of the cooling function on the temperature, while some observations indicate a very large value of t^2 (see Liu et al. 1995, and Kingdon and Ferland 1995 for a discussion). If something is missing from our current understanding of the energy source of photoionized nebulae then the entire nebular abundance scale (for both the Milky Way and the extragalactic nebulae) is in error by as much as 0.5 dex.

Two fundamentally different t^2 s enter here - the “structural” t^2 and the observational t^2 . The structural value comes from the computed ionization and thermal structure of the nebula, while the observational value comes from an analysis of the predicted emission line spectrum following the methods outlined in Peimbert’s 1967 paper.

The structural t^2 for the H^+ ion is defined as

$$t^2(H^+) = \left\langle \left[\frac{T(r) - \langle T \rangle}{\langle T \rangle} \right]^2 \right\rangle = \frac{\int [T(r) - \langle T \rangle]^2 n_e n_{H^+} f(r) dV}{\langle T \rangle^2 \int n_e n_{H^+} f(r) dV} \quad (409)$$

where $\langle T \rangle$ is the density-volume weighted mean temperature

$$\langle T \rangle = \frac{\int T(r) n_e n_{H^+} f(r) dV}{\int n_e n_{H^+} f(r) dV}. \quad (410)$$

This quantity is given in the averaged quantities block as the column “Te(NeNp)”.

The observational t^2 - related quantities are the following: “T(OIIIr)” is the electron temperature indicated by the predicted [OIII] 5007/4363 ratio in the low density limit. This number is meaningless for densities near or above the critical density of these lines. “T(Bac)” is the hydrogen temperature resulting from the predicted Balmer jump and H β . “T(Hth)” is the same but for optically thin Balmer continuum and case B H β emission. “t2(Hstrc)” is the structural HII t^2 . The entries “T(O3-BAC)” and t2(O3-BC)” are the mean temperature and t^2 resulting from the standard analysis of the [OIII] and HI spectra (Peimbert 1967). Finally “t2(O3str)” is the structural t^2 over the O $^{2+}$ zone. Only the structural t^2 s are meaningful for high densities. This section was developed in association with Jim Kingdon, and Kingdon and Ferland (1995) provide more details.

3.8.5 Grains

```
Average grain Properties:
Gra-Ori Sil-Ori
<Tdust>: 1.540E+02 1.218E+02
<Vel D>: 3.635E+04 2.891E+04
<Pot D>: 1.337E+00 1.419E+00
```

The next lines give some information concerning grains if these were included in the calculation. These lines give the mean temperature, drift velocity, and potential, for all of the grain populations included in the calculation. An asterisk will appear to the right of the name of any species with quantum heating included. In this case the mean temperature is weighted by T^4 .

```
Contin Optical Depths: COMP: 6.80E-04 H-: 5.59E-05 R(1300): 2.14E-04 H2+ 1.50E-06 HeTri:6.09E-04
Pfa:3.40E-04 Pa:3.40E-04 Ba:3.65E-04 Hb:3.60E-04 La:1.29E-01 lr:4.995E+07 1.8:1.19E+07 4.:1.738E+06
Line Optical Depths: 10830: 1.23E+02 3889: 5.24E+00 5876: 3.29E-06 7065: 1.82E-06 2.06m 2.54E-03
```

3.8.6 Continuum optical depths

Contin Optical Depths The first two lines give the continuum optical depths at various energies. These are the total optical depths, including the correction for stimulated emission, and will be negative if maser action occurs. These include grain opacity if grains are present. The labels, and their interpretation, are as follows. COMP is Thomson scattering. H- is the negative hydrogen ion at maximum cross section. R(1300) is Rayleigh scattering at 1300Å, H $_2^+$ is the molecular hydrogen ion. HeTri is the helium triplet at threshold. The next line gives total continuous optical depths at the energies of various hydrogen and helium ionization edges and lines. These are the Pfund α , Paschen α , Balmer α and β , L α , and the ionization edges of hydrogen, atomic helium, and the helium ion.

Heavy element line optical depths are printed also if the **print line optical depths** command is entered.

```
Old hydro optical depths: 1 9.99E+07 2 1.00E-20 3 1.00E-20 4 1.00E-20 5 1.00E-20 6 1.00E-20 7 1.00E-20
```

Old H Lines:	2-1	9.96E+19	3-2	3.34E-02	4-3	3.33E-04	5-4	1.66E-05	6-5	2.56E-06	7-6	2.56E-07	8-7	2.56E-08
New hydro optical depths:	1	4.99E+07	3-2	1.46E-05	3	2.82E-05	4	6.80E-06	5	1.25E-12	6	3.05E-12	7	6.49E-12
New H Lines:	2-1	2.02E+06	3-2	9.37E-03	4-3	1.98E-09	5-4	6.50E-10	6-5	5.01E-09	7-6	2.94E-08	8-7	1.81E-07
Old He I optical depths:	1	2.38E+07	2	1.00E-20	3	1.00E-20	4	1.00E-20	5	1.00E-20	6	1.00E-20	7	1.00E-20
Old He II Lines:	2-1	9.96E+19	3-2	1.00E-20	4-3	1.00E-20	5-4	1.00E-20	6-5	1.00E-20	7-6	1.00E-20	8-7	1.00E-20
New He I optical depths:	1	1.19E+07	2	1.46E-05	3	2.83E-05	4	6.06E-06	5	1.41E-12	6	3.33E-12	7	6.93E-12
New He II Lines:	2-1	6.15E+04	3-2	1.12E-03	4-3	3.50E-11	5-4	2.59E-11	6-5	1.25E-10	7-6	5.98E-09	8-7	2.36E-08
Old He I optical depths:	1	9.96E+19	2	1.00E-20	3	1.00E-20	4	1.00E-20	5	1.00E-20	6	1.00E-20	7	1.00E-20
Old He II Lines:	2-1	9.96E+19	3-2	1.00E-20	4-3	1.00E-20	5-4	1.00E-20	6-5	1.00E-20	7-6	1.00E-20	8-7	1.00E-20
New HE I optical depths:	1	1.70E+06	2	4.88E-07	3	1.89E-04	4	1.39E-05	5	2.16E-05	6	2.75E-05	7	2.72E-05
New HE II Lines:	2-1	1.13E+06	3-2	9.14E-05	4-3	1.36E-11	5-4	4.73E-11	6-5	2.17E-10	7-6	5.52E-09	8-7	2.11E-08

Hydrogen and helium optical depths in continua and $n \rightarrow n-1\alpha$ transitions follow. The first two lines are the optical depths assumed at the start of the present iteration, and the second pair of lines gives the newly computed total optical depths. Negative optical depths indicate maser action. For each of the pairs of lines, the first line is the optical depth at thresholds of the first seven levels of hydrogen. The second line gives the optical depths in the first seven of the $n \rightarrow n-1\alpha$ transitions of hydrogen or helium.

3.8.7 Line optical depths

3.8.8 Mean ionization

[illegible]

The two large blocks of output give the mean ionization, averaged over volume, and over radius. The numbers printed are the log of the mean ionization fraction in the various stages. The volume average ionization fraction for ion i of element a is given by

$$\left\langle \frac{n_a^i}{n_a} \right\rangle_{vol} = \frac{\int n_a^i f(r) dV}{\int n_a f(r) dV}. \quad (411)$$

and the radius average by

$$\left\langle \frac{n_a^i}{n_a} \right\rangle_{rad} = \frac{\int n_a^i f(r) dr}{\int n_a f(r) dr}. \quad (412)$$

3.8.9 Continuum

This is only printed if the `print continuum` command is included. Then the following tables, all related to the transmitted continuum, will be printed.

X-Ray Continuum. The next line gives the photon fluxes ($\text{cm}^{-2} \text{s}^{-1}$) in various X-ray bands, if the continuum extends to X-ray energies. The units of the energy bands are keV. The numbers are the numbers of photons exiting the cloud, integrated over the energy bands. This is the net continuum, that is, the incident continuum, less attenuation, with diffuse re-emission from the cloud added on. This is only printed if the `print continuum` command is entered.

Normalized Continuum. This block is a set of ordered pairs giving the emergent Balmer continuum, relative to the continuum which entered the cloud. The first number of each pair is the frequency in Rydbergs. The second is the ratio of the emergent continuum to the incident continuum (i.e., that which went into the cloud). In the absence of optical depth or diffuse emission effects, this block will be equal to 1.000 throughout. This is only printed if the `print continuum` command is entered.

Emergent Continuum. This block gives ordered pairs of energy (in Rydbergs) and the emergent continuum. It is expressed as photon fluxes ($\text{phot Ryd}^{-1} \text{cm}^{-2}$) corrected for r^2 dilution, so as to be directly comparable with the continuum which went into the cloud. This is only printed if the `print continuum` command is entered.

4 OBSERVED QUANTITIES

4.1 Overview

This section describes how to convert the quantities actually used or predicted by Cloudy into commonly observed ones.

4.2 Incident and Diffuse Continua

The emission line printout gives the intensity of the incident continuum (λF_λ or νF_ν) at 4860 and 1215 Å. These appear with the label **inci** followed by the wavelength. The entire incident continuum can be obtained with the output of the **punch continuum** command.

The diffuse continuum, that emitted by the cloud, is not normally included in the line output. The **print diffuse continuum** command will add the total emitted continuum to the emission line list. These are in units λF_λ or νF_ν at the indicated wavelengths and have the label **nfnu**. The entry with the label **nTnu** is the sum of the reflected plus attenuated incident continuum. The inward total emission and the reflected incident continua will be printed if this command appears together with the **print line inward** command. Two contributors to the inward emission are predicted. That labeled **InwT** is the total inwardly emitted continuum, and includes both diffuse emission and the back-scattered incident continuum. The component labeled **InwC** is the back-scattered incident continuum alone.

4.3 Line Equivalent Widths

The equivalent width of an emission or absorption line is defined as the number of Ångstroms of the continuum that is equivalent to the energy in the line. It can be defined as

$$W_\lambda = \int \frac{F_\lambda^c - F_\lambda^l}{F_\lambda^c} d\lambda \approx -\lambda \frac{F_{line}}{\lambda F_\lambda^c} \quad (413)$$

where the fluxes are in the interpolated continuum (F_λ^c) and the integrated line (F_{line}). By this convention the equivalent width of an emission line is negative.

The code predicts the integrated fluxes of all lines. It also predicts the product λF_λ^c for the incident continuum at a few wavelengths. These are given the label **inci** and the wavelength where it is evaluated follows. The entry **inci 4860** is the intensity of the incident continuum at a wavelength near H β . The units of this incident continuum are either erg cm⁻² s⁻¹ or erg s⁻¹ depending on whether the incident continuum was specified as a flux or luminosity. The fluxes of lines and these continuum points can be read from the output, or obtained by software calling the **cdLine** routine. The continuum flux at any wavelength can be obtained with the **punch continuum** command. If the line intensity is given by F_{line} and the continuum intensity λF_λ^c , then the equivalent width of a line relative to the continuum where λF_λ^c is specified will be given by the last term in equation 413.

A covering factor will complicate this slightly. (Covering factors are defined in the section **Definitions** in Part I of this document.) If luminosities are predicted then partial coverage of the source is taken into account with the **covering factor** command, and the luminosities are correct for this coverage. The ratio of line to continuum given in equation 413 will represent what is observed. If fluxes are specified instead then the line flux is given per unit area of cloud, no matter what covering factor is specified. In this second case the ratio in equation 413 must be scaled by the covering factor.

4.4 Emission Line Asymmetries

The inward fraction of the total emission of each line is always predicted by the code, but not normally printed out. Many lines are significantly inwardly beamed, and this can lead to emission line asymmetries if the envelope is expanding. The inward part of the lines will be printed if the **print line inward** command is entered. The effects of this line beaming is very geometry dependent.

4.5 Line to Continuum Contrast

The code has several **punch** commands that will produce ancillary files containing the predicted line and continuum spectra.

There is an ambiguity in how strong the lines should appear to be relative to the continuum in a plot where the lines are not resolved. This is described in Part I of this document where the **punch continuum** and **set PunchLWidth** commands are introduced.

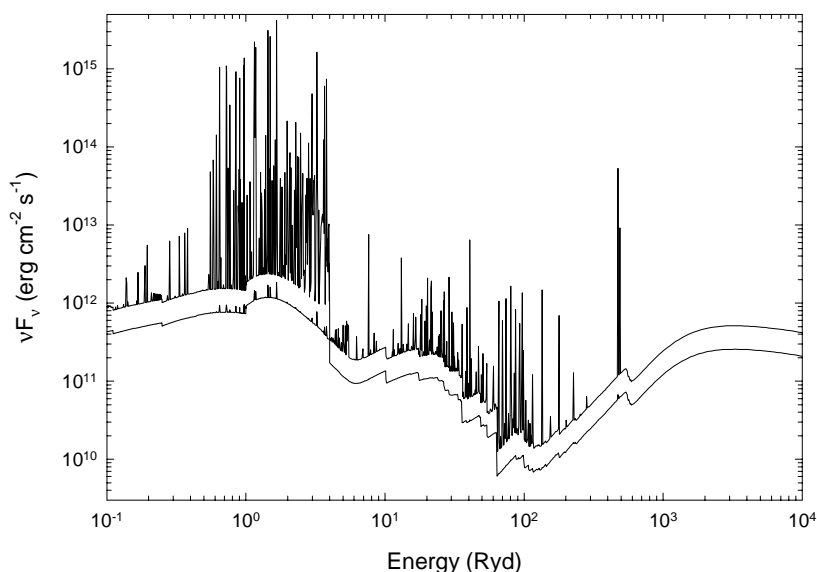


Figure 2 This is the continuum predicted by the input file reflector.in. The lower curve has been divided by two and shows the total spectrum produced by setting the line width to the speed of light. The upper curve shows the same thing, but with the line width set to 100 km/sec. reflector

Figure 2 shows the continuum predicted with the **reflector.in** test case. The lower curve shows the total continuum that would be predicted if the **PunchLWidth** variable is set equal to the speed of light⁵. Here lines are added to the continuum such that the difference between νF_ν at the line peak and νF_ν for the underlying diffuse continuum is equal to the line flux. As a result the resulting line to continuum contrast is very small. The

⁵ This was the default for version 90.00 through version 90.03. In C90.04 the default was changed to 1000 km/s. Before version 90 the line to continuum contrast depended on the cell width at the particular energy.

upper curve shows the same model but with the line contrast enhanced by entering the command `set PunchLWidth 100 km/sec`. The entire spectrum is shifted by a factor of two to make the two appear separated. The default line width is 1000 km s⁻¹.

The only effect of the `set PunchLWidth` command is to change the contrast in the punch output. The computed results and line intensities in other output are not affected. If the width is set to the speed of light then the intensities in the punch output will be correct but the line to continuum contrast too small. If the width is set to a small value the contrast is increased but the total intensity in the punch output will be greater than the actual emission. (Energy *will not* appear to have been conserved in this punch output).

4.6 Surface Brightness

Cloudy will normally predict $4\pi I$, the intensity radiated into 4π sr by a unit area of cloud, erg cm⁻² s⁻¹. Observations of resolved sources often measure the surface brightness, with units erg arcsec⁻² s⁻¹. Be careful! – some workers may report surface brightness with units erg arcsec⁻² s⁻¹ sr⁻¹. Remove the sr⁻¹ before continuing by multiplying by 4π .

To obtain the surface brightness we must divide the intensity $4\pi I$ predicted by Cloudy by the number of square seconds of arc in 4π sr. One radian is $180/\pi = 57.29578$ deg, so 1 sr is $(180/\pi)^2 = 3282.806$ deg², and there are 5.3464×10^{11} square arc seconds in 4π sr. The surface brightness (per square second of arc) is the intensity $4\pi I$ (per square centimeter) multiplied by the inverse of this, or 1.8704×10^{-12} cm² arcsec⁻².

Note that this is only correct for a line that is emitted isotropically, because the code predicts $4\pi I$ while an observer measures I along a particular direction. This discussion is only correct if $I = J$.

4.7 Flux to luminosity

The luminosity is the flux of a line multiplied by the total area of the shell. For full coverage this is $4\pi r^2$ where r is the radius of the shell. If the shell only partially covers the continuum source then this should be multiplied by the covering factor.

4.8 Relative hydrogen line intensities

Very accurate ratios of Balmer or Paschen lines of hydrogen can be used to determine reddening. Ferguson and Ferland (1997) describe Cloudy's hydrogen atom. It gives good results for levels below 10 in the code's default state, which uses a 15 level atom. The number of levels can be increased to ≤ 50 using the `hydrogen levels` command, and this gives better results at the expense of more compute time. The larger atom should give results accurate to better than 5% for lines arising from below principal quantum number 10, and 10% accuracy for lines with upper levels between 10 and 15. The accuracy decreases for upper levels higher than 15 although the total recombination efficiency of the atom is computed to high precision.

For pure recombination lines you can easily get better predicted relative intensity than those predicted by Cloudy. The code is limited by the size of the model

hydrogen atom that can be computed on the fly. The definitive calculation for hydrogen recombination is that of Hummer and Storey (1987), who used a 1000 level atom with all l -states explicitly considered (that works out to something like a million levels!). Storey and Hummer (1995) placed a program on the web that will interpolate on their tables of case B hydrogen emission, for any temperature and density they computed. The best way to obtain a very high quality hydrogen recombination spectrum is to get the mean H^+ temperature and the electron density (perhaps those predicted by Cloudy) and then use their interpolating code to provide the hydrogen spectrum for these conditions.

The Hummer and Storey (1987) calculation is for case B conditions, which assume that many processes are unimportant (see Ferguson and Ferland 1997). Neglected processes include collisional excitation from the ground or first excited states, induced processes where the incident continuum causes the atom to fluoresce, and line transfer in all non-Lyman lines. This is an excellent assumption for conventional nebulae, such as planetary nebular or HII regions. They are questionable for gas denser than 10^6 cm^{-3} or when X-Rays are present. When any of these processes are important the hydrogen spectrum is far more model dependent and Cloudy's results are more realistic than the case B results.

4.9 Line Intensities in a dusty open geometry

Two sets of line intensities are printed if a dusty open geometry is computed. The second block of lines is the conventional set of intrinsic emission line intensities. When grains are present these intensities would need to be corrected for line of sight reddening to be compared with observations.

The first block of emission line intensities would be that emitted from the illuminated face of a molecular cloud. The geometry is appropriate for the Orion Nebula, a blister HII region on the surface of Orion Molecular Cloud 1 (OMC1). An idealized geometry is shown in Figure 3. The code computes the fraction of the line emission that is directed towards the illuminated face. The remainder is emitted towards the neutral gas, which is assumed to have an infinite optical depth due to grains. The local albedo of the gas-grain mixture is computed, and the fraction reflected is passed back towards the illuminated face. The resulting intensities are roughly half what would

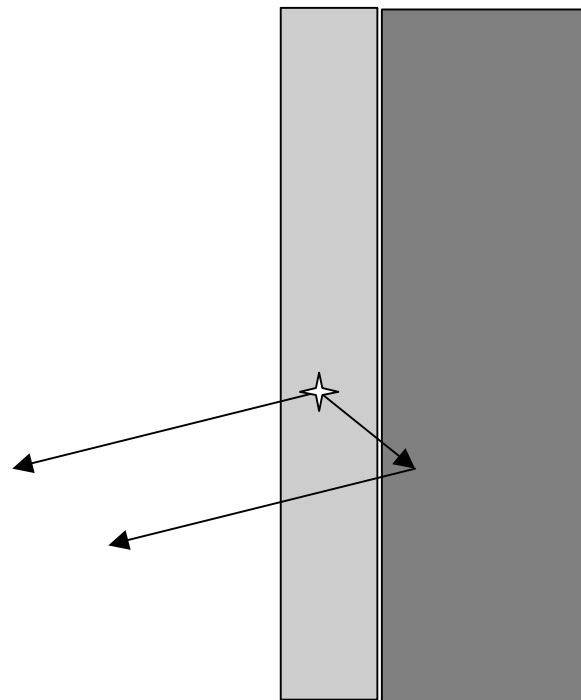


Figure 3 This shows that geometry assumed when computing the first block of lines in an open dusty geometry. The light area in the center is the H II region, which is assumed to be a layer on the surface of an infinitely optically thick molecular cloud, the dark area on the right. Light can be emitted towards, and freely escape from, the illuminated face of the cloud. A fraction of the light emitted towards the molecular cloud is reflected back towards the illuminated face.

be expected were the cloud emitting from both sides. Something like 10% of the line striking the molecular cloud will be reflected back to the observer.

So, for the illustrated blister the first block of lines gives what would be seen by an observer a large distance off to the left.

4.10 Continuum pumping contribution to line intensities

Continuum pumping or fluorescence is included for all lines. The contribution is usually not explicitly printed, but will be if the **print line pump** command is entered. Whether or not this contribution actually adds to the observed line emission depends on the geometry. Continuum pumping increases the line emission if no related continuum absorption occurs. This will be the case if the continuum source is either not observed or not covered by absorbing gas. If absorbing gas covers an observed continuum source then the situation is like the P Cygni problem, and pumping may not increase the net intensity of the line at all (the absorption component will have the same equivalent width as the associated emission). The printed line intensity includes this contribution unless the **no induced processes** command is entered. (The **no induced processes** command has many other effects and so should not be used except as a test.)

The output produced by the **punch continuum** commands does not include the pumped part of the line contribution. This is correct if the continuum source is included in the beam, but is not if only the gas is observed.

5 THE EMISSION LINES

5.1 Overview

The following sections outline the emission lines predicted by Cloudy. Before version 90 of the code all lines were listed in sub-section 5.3.1, immediately following this section. The code is being modified to bring all lines into a common line class, as the code moves to C++ and objects.

5.2 The main emission line printout

The main emission line printout was briefly described on page 428 above. This section goes into more detail.

Output organization. The printed list is sorted into four large groups of columns, with each large column sub-divided into four smaller sub-columns. The first sub-column is either the spectroscopic designation of the ion producing the line or an indication of how the line is formed. The second sub-column is the line wavelength, with a 0 to indicate a continuum. The third sub-column is the log of the power in the line, in the units given in the header (4π sr or cm^{-2}). The last sub-column is the intensity of the line relative to the reference line, usually $\text{H}\beta$, unless this is reset with the `normalize` command.

Line intensities when grains are present. The computed emission-line spectrum follows. Emission lines are divided into two large groups. The first includes the effects of grain scattering and absorption, and is indicated by the header “Emergent Line Intensities”. This first group is only printed if grains are present and the geometry is open (i.e., `sphere` not set). The intensities are the *total* intensities observed from the illuminated face, including both absorption and scattering by grains. This is discussed on page 438 above.

The second, larger, group of lines, called “Intrinsic line intensities”, is always printed. This usually gives the intrinsic intensity of the lines, and does not include the reddening effects of internal grains due to the photon's passage out of the nebula (unlike the first group). This second group usually gives the total intrinsic intensity of the lines. Although reddening effects of internal (or external) dust are not taken into account, photon destruction by background opacity sources during the transfer process is. This distinction is only important for forbidden lines, which have no local destruction since they are optically thin, but can be absorbed along their way out. This predicted spectrum should be compared with the reddening-corrected observed spectrum.

Line wavelengths. In the current version of the code the line wavelength is an integer. This is to allow a precise match with a wavelength specified during a search. An integer format has obvious limitations in specifying fractions. To overcome this, the format of the wavelength changes with wavelength, in an attempt to have several significant figures.

Many wavelengths are too small or large to be expressed in Ångstroms. Many FIR lines would overflow the output format if their wavelengths were expressed in Ångstroms. The first letter following the line wavelength in the printout indicates

the units of the line. An “A” indicates Ångstroms and ”m” microns. This letter may be followed by an integer indicating how many places to the left the decimal point should be moved. For instance, the $P\alpha$ line of hydrogen would have a wavelength of 187m2, indicating 1.87 μ . The [C II] $\lambda 157$ micron line would be C 2 157 and the [O III] 88 micron line would be O 3 883m1. The logic deciding this format is given in routine *iWavLen*.

Visible wavelengths are usually given in air. Continua are usually indicated by a wavelength of zero.

5.2.1 Blocks of lines....

Within each column lines are organized by common origin with a comment beginning the section. As an example, the first commented block of lines begins with “general properties.....”. The following subsections give overviews of the lines.

5.2.2 General properties....

This mainly summarizes heating and cooling agents for the model.

TOTL 4861 and **TOTL 1216**, are the total intensities of $H\beta$ and $L\alpha$, as predicted by the multi-level H atom. These intensities are the results of calculations that include all collisional, radiative, and optical depth effects.

Inci - The total energy in the incident continuum.

TotH and **TotC** give the total heating and cooling. These will be nearly equal in equilibrium.

BFH1 and **BFHx** are the heating due to photoionization of ground state and excited state hydrogen respectively.

He1i, **3He1**, heating due to ground state He and the triplets.

BFHe and **TotM** are the heating due to helium and metal photoionization.

Pair – heating due to pair production.

ComH, **ComC**, – Compton heating, cooling.

CTH CT C – charge transfer heating and cooling.

extH extC “extra” heating or cooling added to model.

e-e+ 511 The positron line.

Expn, expansion, or adiabatic, cooling

HFB, H radiative recombination cooling

HFBc, **HFBc**, hydrogen net free-bound cooling and heating

Iind, cooling due to induced recombination of hydrogen

3He2, cooling due to induced recombination of fully ionized helium

Cycn, cyclotron cooling

5.2.3 Continua....

These give intensities of various continua. These are either the total integrated continuum or the product νF_ν at certain energies.

Bac 3646 residual flux at head of Balmer continuum, νF_ν

cout 3646 cref 3646, outward, reflected continuum at peak of Balmer Jump

thin 3646, residual flux at head of Balmer continuum, optically thin limit

Inci 4860, **Inci 1215**, incident continua near $H\beta$ and $L\alpha$

Ba C 0, integrated Balmer continuum

PA C 0, integrated Paschen continuum

HeFF 0, He brems emission

HeFB 0, He recombination cooling

MeFB 0, heavy element recombination cooling

MeFF 0, metal brems emission

ToFF 0, total brems emission

FF x, part of H brems, in x-ray beyond 0.5KeV

eeff, electron - electron brems

nFnu 122m, **nInu 122m**, **InwT 122m**, **InwC 122m**, a large list of continua at selected wavelengths will be printed if the print diffuse continuum command is entered. The first is the total continuum at the wavelength, given as νF_ν . **nInu** is the transmitted and reflected incident continuum. **InwT** is the total reflected continuum. **InwC** is the reflected incident continuum.

5.2.4 Molecules....

H2 I 2 is the intensity of the H_2 lines near $2\mu m$.

H2dC, is the cooling due to collisional dissociation of H_2 .

H2dH, heating by H_2 dissociation by Lyman continuum

H2vH, heating by coll deexcit of vib-excited H_2

H2vC, cooling by coll deexcit of vib-excited H_2

H2 v, line emission by vib-excited H_2

H-FB and **H-FF** are the free-bound and free-free continua of the H^- ion.

H-CT 6563, H-alpha produce by H^- mutual neutralization

H- H 0, H^- heating

H-Hc 0, H^- heating

H2+ and **HEH+** are the cooling due to formation of H_2^+ and HeH^+ .

Codh, carbon monoxide photodissociation heating

CO C, old cooling due to collisions of vibrational rotational levels (used pre-c95)

CO 12, $C^{12}O^{16}$ cooling

CO 13, $C^{13}O^{16}$ cooling

12CO 2588m, **Inwd 2588m**, **Coll 2588m**, **Pump 2588m**, **Heat 2588m**, et al. Next follows intensities and contributors to the ^{12}CO and ^{13}CO lines included in the calculation.

5.2.5 Grains....

Information in this block concerns emission, absorption, heating, and cooling by any grains included in the calculation.

GrGH, gas heating by grain photoionization

GrTH, gas heating by thermionic emissions of grains

GrGC, gas cooling by collisions with grains

GraT, total grain heating by all sources, lines, collisions, incident continuum

GraI, grain heating by incident continuum

GraL 1216, grain heating due to destruction of Ly alpha

GraC, grain heating due to collisions with gas

GraD, grain heating due to diffuse fields, may also have grain emission

5.2.6 H-like iso-seq...

This block includes all hydrogen-like isoelectronic species.

HFFc 0, net free-free cooling, nearly cancels with cooling in lte

HFFh 0, net free-free heating, nearly cancels with cooling in lte

H FF 0, H brems (free-free) cooling

FF H 0, total free-free heating

Clin 912, total collisional cooling due to all hydrogen lines

Hlin 912, total collisional heating due to all hydrogen lines

Cool 1216, collisionally excited La cooling

Heat 1216, collisionally de-excited La heating

Crst 960, cooling due to n>2 Lyman lines

Hrst 960, heating due to n>2 Lyman lines

Crst 4861, cooling due to n>3 Balmer lines

Hrst 4861, heating due to n>3 Balmer lines

Crst 0, cooling due to higher Paschen lines

Hrst 0, heating due to higher Paschen lines

LA X 1216, la contribution from suprathermal secondaries from ground

Ind2 1216, "Ly alpha" produced by induced two photon

Pump 4861, H-beta produced by continuum pumping in optically thin ld limit

CIION 0, net col ionz-3 body heat collision ionization cooling of hydrogen

3bHt 0, heating due to 3-body recombination

Strk 1216, Stark broadening component of line

Dest 1216, part of line destroyed by background opacities

Fe 2 1216, part of La absorbed by Fe II

Q(H) 4861 is the intensity of H β predicted from the total number of ionizing photons, Q(H), assuming that each hydrogen-ionizing photon produces one hydrogen atom recombination.

Q(H) 1216 indicates the L α intensity produced if each hydrogen ionizing photon results in one L α photon in the high density limit (i.e., no two-photon emission).

CaBo 4861 These are the "old" case B predictions, as printed in versions 90 and before of the code.

Ca B 6563A The entries starting with Ca B are the Case B intensities computed from the actual model ionization and temperature structure, but assuming that H β emits with its case B emissivity.

Next the predicted intensities of all lines of the hydrogenic iso-electronic sequence are given. The lines have labels that identify the species and stage of ionization, such as H 1, He 2, Li 3, C 6, etc. The entries with a wavelength of zero are the total intensities of the $2s-1s$ two-photon emission.

5.2.7 He iso-sequence....

Atoms and ions of the helium-like iso-electronic sequence are treated as multi-level atoms. Each species and stage of ionization are specified by labels like He 1, Li 2, C 5, etc. A wavelength of zero indicates the two-photon continuum.

5.2.8 level 1 lines....

In the current version of the code, the lines printed under this title include both the lines that have been moved to the new common *EmLine* class, but also older lines that are still scalar quantities. This part of the code is still in a state of flux, and this is reflected in the current documentation. The remaining part of this subsection outlines the methods used for most of the heavy element atoms. The method for producing a list of transferred lines, those that have been moved to the *EmLine* class, is described in the section beginning on page 445 below. The old-style scalar lines are described in the section beginning on page 447 below, although this is not totally up to date.

These lines have accurate collision strengths and wavelengths. Many are two-level atoms, but some are the result of multi-level atoms. The following is a summary of the general approach.

Li-sequence. Examples include C IV $\lambda 1549$, O VI $\lambda 1034$, and Mg II $\lambda 2798$. A three level atom, with full treatment of optical depths and collisional excitation, is used. The “TOTL” intensity is the sum of both lines in the doublet, and is followed by the individual intensities of each member.

Be-sequence. Examples include C III] $\lambda 1909$, O V] $\lambda 1215$, and Si III] 1895. A four level atom, solving for populations of the individual 3P_j states, is used. The first printed intensity is the total intensity of the multiplet (both $j=0,1$ decays), and this is followed by the intensities of individual lines. The intensity of the permitted $^1P_0 - ^1S$ transition is also calculated. Optical depth and collisional effects on both the permitted and intercombination lines are included.

B-sequence. Examples include C II and O IV. The ground term is treated as a two level atom, with optical depth and collisional effects included, when the gas is too cool to excite the UV lines. The $^4P - ^2P_0$ lines are also predicted with a full multi-level atom that resolves fine structure. The TOTL intensity printed is the total intensity of the multiplet and is followed by individual lines.

3P - ground term. Examples include such spectra as [O III] and [O I]. The infrared fine structure lines are computed with full treatment of collisional and optical depth effects. A comment is printed at the end of the model if these lines mase or become optically thick. The populations of 1D and 1S are computed with a three-level atom.

The intensity of the $^1D - ^3P$ transition is only that of the individual line (i.e. 5007), not the doublet.

$^4S^0$ - ground term. Examples include [O II] and [S II]. They are treated as a five level atom. Intensities of all individual lines, as well as co-added multiplets, are given.

5.2.9 Recombination . . .

These are a set of heavy-element recombination lines that are predicted assuming that they are optically thin. This consists of all recombination lines of C, N., and O, with coefficients taken from Nussbaumer and Storey (1984) and Pequignot, Petitjean, and Boisson (1991). For this set, the spectroscopic designation is followed by the wavelength and the log of the recombination coefficient evaluated at 10^4 K.

These are all predictions for optically thin pure recombination. These should be accurate for classical nebulae, such as planetary nebulae and HII regions. They will not be accurate for dense environments where optical depths and collisional effects come into play. There are several instances where more than one line of an ion will have the same wavelength due to the integer Ångstrom format used for wavelengths. The worst case is O V 4953, where three lines of the same multiplet have the same wavelength.

5.2.10 Level 2 lines . . .

These are resonance lines that use Opacity Project wavelengths, which are generally accurate to about 10%. These lines have g-bar collision strengths, which are not very accurate at all.

5.3 The transferred lines

The group of “transferred lines” includes all those that have been moved to the *EmLine* class, in anticipation of the codes move to C++ and objects.

In versions of HAZY for Cloudy versions 90 and before, this section included descriptions of all predicted lines, and was automatically generated by the code. Today there is no limit to the number of lines the code is capable of predicting, since many iso-electronic sequences can now have a nearly arbitrarily large number of levels. Rather than waste paper by including the iso-electronic sequences here, instructions are given for creating your own automatic list of lines.

To generate a line list, set up a calculation with the atoms set to whatever size is desired (see the **atom** command in Part I). Then execute this script with the **punch line data** command included (described in Part I). The punch output will include the line list. This will include the level 1 ,level 2, CO, and recombination lines, but not the scalar forbidden lines. These are described in a list following this subsection.

5.3.1 Punch line data output

In previous versions of this document a large list of emission lines appeared here. This list is now far too large to include here. Rather, the list can be generated by executing the code with the command **punch line data “filename.txt”** included. This will create a file that includes the full set of lines that are predicted.

Note that the lines that are output are only those that exist when the code is run. It is possible to make many of the model atoms and molecules as large or small as you like, and the actual lines that exist when the punch command is entered will be output.

This contains several groups of lines. All quantities were evaluated to 10^4 K. The description of the command in Part I of this document explains how to evaluate the quantities at other temperatures.

The ion is the first column of the table. This is in a uniform format, beginning with the two character element symbol and followed by an integer indicating the level of ionization. "C 2" is C^+ or CII. This is followed by the integer wavelength label used to identify the line in the printout. The third column, with the label "WL", is the correct wavelength of the line, with units of either microns ("m") or Angstroms ("A"). The remaining columns give the statistical weights of the lower and upper levels, the product of the statistical weight and the oscillator strength, and then the Einstein A.

The last column is the electron collision strength, generally evaluated at 10^4 K. Exceptions are lines whose collision strengths are only evaluated for temperatures far from 10^4 K, for instance, a Fe XXV transition. Usually these collision strengths are for only the indicated transition, although in some cases (the Be sequence) the value is for the entire multiplet. This is discussed further in the section on the evaluation of the cooling function in HAZY II.

5.3.2 Output produced for the transferred lines

Because the lines have a common format within their storage vectors, the output has a common format too. Generally only the total intensity of the transition, the result of the solution of a multi-level atom with all processes included, is done. The approach used to compute the level populations is described in Part II of Hazy, and includes continuum pumping, destruction by background opacities, and trapping.

The total intensity of the transition is printed in a form like "C 2 1335", with the spectroscopic identification given by the first part, as found in the first column of the table, and the wavelength as indicated by the integer in the second column of the table.

In a few cases (for instance, the C 4 $\lambda\lambda$ 1548, 1551 doublet), a total intensity is also derived. In these cases the label "Totl" will appear together with an average wavelength (1549 in this case). These lines are all explicitly shown in section 5.3.1 on page 445 above.

It is possible to break out various contributors to the lines with options on the **print line** command, described in Part I of this document and in the following. These contributors are printed following the total intensity.

print line heating An emission line will heat rather than cool the gas if it is radiatively excited but collisionally de-excited. The print out will include an entry beginning with the label "Heat" if this printout is turned on.

print line collisions The collisional contribution to the lines will be printed, with the label "Coll".

print line pump The contribution to the total line, produced by continuum pumping, is printed with the label “Pump”. What is observed? Whether or not this is a net emission process contributing to the observed line intensity depends on the geometry, whether or not continuum source is in the beam. At some velocities within the line profile this can be a net emission process, due to absorption at other velocities. If the continuum source is in the beam and gas covers it, this is not a *net* emission process, since photons are conserved.

print line inward The inwardly directed part of the total emission is printed with the label “Inwd”. This can be greater than half of the line intensity if the line is optically thick.

print line optical depths At the end of the calculation the optical depths for all optically thick lines will be printed. This is not done by default since it is quite long.

5.4 Forbidden Lines

These are a series of entries that contain most of the optical forbidden lines, some continua, and identify various contributors to the main lines. These are older lines that have not yet been moved to the *EmLine* class. This description is not totally up to date since this is a part of the code that is slowly being removed as lines go to the new style.

For this set of lines, the first column gives the four character label printed in the final array listing and the second column gives the wavelength of the line, using the conventions described above. The label in the first column is the one used to access the line using the *cdLine* routine described elsewhere.

The third column character indicates whether the entry is a heat source (indicated by h), a coolant (c), a recombination line (r), or an intensity entered for information only (i). The last column gives a brief description of the meaning of the line prediction.

Label	λ	Description
Mion	0 c	cooling due to collisional ionization of heavy elements
Li3r	19 i	these lines added to outlin in metdif - following must be false
Be4r	19 i	these lines added to outlin in metdif - following must be false
Bo5r	19 i	these lines added to outlin in metdif - following must be false
REC	1656 i	C 1 1656 recomb; n.b. coll deexcitation not in
C Ic	9850 c	C 1 9850, coll excit
C Ir	9850 i	was a big mistake
TOTL	9850 i	total intensity, all processes, C I 9850
C 1	8727 c	C 1 8727; equivalent to 4363
C 1	4621 c	1S - 3P
Phot	2326 i	photoproduction, Helfand and Treffitz
REC	1335 i	C 2 1335 recombination,
C II	3134 c	C 2 intercombination line with same upper state as 1335
C3 R	977 i	dielectronic recombination contribution to C 3 977
P386	977 r	C 3 977 pumped by continuum near 386A
TOTL	1909 i	C 3 1909 collision, both lines together
C 3	1907 i	C 3 1908 j-2 to ground
C3 R	1909 i	C 3 1909 recombination from Storey
Phot	1909 i	C 3 1909 following relax following inner shell photoionization
Rec	1175 i	dielectronic recombination contribution to C 3 1175
TOTL	1549 i	total intensity of C 4 1549, all processes
Inwd	1549 i	inward part of C 4

5 THE EMISSION LINES

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DEST 1549 i part of line destroyed by photoionization of Balmer continuum
C4 r 1549 i recombination C 4 1549 from CV
C 6r 34 i these lines added to outlin in metdif - following must be false

N 1 5200 i N 1 5200, both 5198, 5200, collisions and recombination
Coll 5200 c N 1 5200, both 5198, 5200, collisions and recombination
REC 5200 i recombination contributon to [NI] 5200
N 1 3466 c [N 1] 3466, 3 - 1 transition, whole multiplet
N 1 10400 c [N 1] 10400 3 - 2 transition, whole multiplet

N 2 6584 c N 2 6584 alone
N 2 6548 c N 2 6548 alone
REC 6584 i N 2 6584 alone, recombination contribution
N 2 5755 i N 2 5755 total, collisions plus charge transfer
Coll 5755 c N 2 5755 collisional contribution

C T 5755 c N 2 5755 charge transfer contribution
Rec 1085 i dielectronic recombination contribution to N 2 1085
N2cn 1 i continuum pumped N 2 6584
N2cn 5755 i continuum pumped N 2 5755
N3cn 4640 i continuum pumped "Bowen" N 3, optically thin excited line

N3cn 4634 i continuum pumped "Bowen" N 3, optically thin excited line
N3cn 4642 i continuum pumped "Bowen" N 3, optically thin excited line
extr 990 i total N 3 990, both electron excitation and continuum pumping
rec 990 i part of N 3 990 due to recombination
N 3p 990 r N 3 989.8, continuum pumped

TOTL 1486 i N 4] 1486, total intensity of both lines
N 4 1485 i the N 4] slow transition by itself
rec 765 i N 4 765 recombination,
TOTL 1240 i N 5 1240, total emission, collisions plus pumping
Inwd 1240 i inward part of N 5

N 7r 25 i these lines added to outlin in metdif - following must be false
Fl7r 19 i these lines added to outlin in metdif - following must be false
O 1 6300 c total Oxygen I 6300, including line optical depth
O 1 6363 c total Oxygen I 6363, including line optical depth
O 1 5577 c auroral OI

TOIc 0 c total collisional cooling due to 6-level OI atom
TOIh 0 h total collisional heating due to 6-level OI atom
6lev 8446 i be moved to call PutLine
6lev 1304 i OI 1304 from six level atom
6lev 1039 i OI 1039 from six level atom

6lev 4368 i OI 4368 from six level atom
6lev 13 i OI 1.3 micron from six level atom
6lev 11 i OI 1.1 micron from six level atom
6lev 29 i OI 2.9 micron from six level atom
6lev 46 i OI 4.6 micron from six level atom

TOTL 3727 c O II 3727, all lines of multiplet together
TOTL 7325 c O II 7325, all lines of multiplet together
IONZ 3727 i line produced by photoionization of Oo; already in TOTL
IONZ 7325 i line produced by photoionization of Oo; already in TOTL
O II 3729 i five level atom calculations; D5/2 - S3/2

O II 3726 i D3/2 - S3/2 transition
O II 2471 c both 2P 1/2 and 3/2 to ground
O II 7323 i P1/2-D5/2 and P3/2-D5/2 together
O II 7332 i P1/2-D3/2 and P3/2-D3/2 together
TOTL 1665 i total intensity of OIII] 1665, all processes

Phot 1665 i contribution to OIII 1665 due to inner shell (2s^2) ionization
Augr 1665 i contribution to OIII 1665 due to K-shell ionization
O 3 5007 c fac = c5007/(1.+1./2.887)
O 3 4959 c O III 4959 alone, collisions, tot OIII is this times 4
LOST 5007 i O III 5007 lost through excit photo

TOTL 4363 i O III 4363, sum of rec, coll, ct excitation
Coll 4363 c O III 4363,collisions from five level atom
Rec 4363 i O III 4363 recombination, coef from Burgess and Seaton
O 3 2321 c collisional excitation of 2321, 5-level atom
C EX 4363 i charge exchange, Dalgarno+Sternberg ApJ Let 257, L87.

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C EX 5592 i charge exchange rate, D+S
rec 835 i O III 834A, dielectronic recombination only
InSh 1401 i inner shell photoionization, relaxation
rec 789 i O IV 789A, dielectronic recombination only
rec 630 i O V 630A, dielectronic recombination only

TOTL 1218 i O V 1218], total intensity of both lines
O 5 1211 i the slow transition by itself
O 5 5112 i BS O V 5112, recombination
TOTL 1035 i O VI 1035, total of pumping and collisional excitation
Inwd 1035 i inward part of OVI line

O 8r 19 i recombination from fully stripped ion
Ne 3 3869 c Ne III 3869, of 3968+3869 doublet
Ne 3 3968 c Ne III 3968, of 3968+3869 doublet
Ne 3 3343 c NeIII auroral line
Ne 3 1815 c NeIII auroral line

Ne 4 2424 c Ne IV 2424, collisional excitation
Ne 4 4720 c Ne IV N=3 lines, three level atom approx
Ne 4 1602 c Ne IV N=3 lines, three level atom approx
Ne 5 3426 c Ne V 3426 of 3426, 3346 doublet
Ne 5 3346 c Ne V 3346 of 3426, 3346 doublet

Ne 5 2976 c auroral line
Ne 5 1575 c collisionally excited
Ne 5 1141 c both components of 5S-3P 1146.1, 1137.0 doublet
TOTL 895 i Ne VII 895, collisionally excited, both lines
Ne 7 890 i Ne VII 890, single line

TOTL 774 i Ne VIII 774, collisionally excited
Inwd 774 i inward part of NeVIII 774 line
NeLr 12 i these lines added to outlin in metdif - following must be false
Na 5 1365 c [NaV] 1365, sum of 1365.1+1365.8; cs only guess
Na 5 2067 c [NaV] 2067, sum of 2066.9+2068.4; cs only guess

Na 5 4017 c [NaV] 4017, sum of 4010.9+4016.7+4022.7; cs only guess
Na 6 2569 c [Na VI] 2568.9
Na 6 1357 c [Na VI] 1356.6
Na 6 2972 c [Na VI] 2971.9
Na 6 2872 c [Na VI] 2872.7

NaLr 10 i these lines added to outlin in metdif - following must be false
TOTL 2798 i Mg II 2798
Inwd 2798 i inward part of Mg II 2798
Mg 6 1806 c MG VI
TOTL 615 i Mg 10 614.9 bothof doublet, li seq 2s 2p

MgLr 7 i these lines added to outlin in metdif - following must be false
totl 2665 i total emission in Al II] 2669.7, 2660 doublet
Al 2 2660 i emission in Al II] 2669 alone
TOTL 1860 i Al III
Inwd 1860 i inward part of AlIII line

Al 6 2428 c [Al VI] 2428.4
Al 6 2601 c [Al VI] 2601.0
Al 6 1170 c [Al VI] 1169.86
Al 6 2125 c [Al VI] 2124.95
TOTL 556 i Al 11, Li seq 2s2p

AlLr 6 i these lines added to outlin in metdif - following must be false
diel 1260 i SI II 1260, rough guess of dielec contribution
diel 1909 i dielectronic recombination SiII 1909
rec 1207 i Si III 1207, dielectronic recombination only
TOTL 1888 i Si III] 1892+1883, total intensity of both lines

Si 3 1883 i Si III] 1883 by itself
PHOT 1895 i photoproduction by inner shell removal
TOTL 1397 i Si IV 1397, collisionally excited
Inwd 1397 i inward part of SiIV 1397
Si 7 2148 c SI VII, 2148, O III like, collisionally excited

Si 7 2148 c
Si 8 1446 c SI VIII 1446, OIII like, collisionally excited
Si 9 1985 c SI IX 1985, 2150, collisionally excited
Si 9 949 c collisionally excited

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Si 9 1815 c collisionally excited

Si 9 691 c both components of 5S-3P doublet
Si10 606 c SI 10 606A, actually group of 4 intercombination lines.
Si11 583 c SI XI 582.9, collisionally excited
TOTL 506 i
SiLr 6 i these lines added to outlin in metdif - following must be false

Pl5r 19 i these lines added to outlin in metdif - following must be false
S 1R 1807 i this is to check whether photoexcit of S II is ever important
S 2 6720 c S II 6731 + 6716 together
S 2 4074 c S II 4070 +4078 together
S 2 10330 c S II N=3 lines, all four lines together

S II 6731 i individual line from five level atom
S II 6716 i individual line from five level atom
S II 4070 i individual line from five level atom
S II 4078 i individual line from five level atom
S II 10323 i individual line from five level atom

S II 10289 i individual line from five level atom
S II 10373 i individual line from five level atom
S II 10339 i individual line from five level atom
S 3 9532 c [S III] 9532 alone
S 3 9069 c [S III] 9069 alone

S 3 6312 c [S III] 6312, transauroral temperature sensitive
S 3 3722 c [S III] 3722, same upper level as 6312
TOTL 1198 i S V 1198] both lines together
S 5 1188 i Be seq, weaker of the two transitions
TOTL 933 i total S VI 933+944

S 9 1715 c S IX 1715, 1987, collisionally excited
S 10 1213 c S X 1213, 1197, collisionally excited
S 11 1826 c S XI 1615, 1826, collisionally excited
S 12 520 c group of four intercombination lines
S 13 488 c S XIII 488.4, 1909 like, collisionally excited

TOTL 427 i S 14 506 li seq 2s2p
S LR 5 i these lines added to outlin in metdif - following must be false
S LR 5 i
Cl 2 8579 c Chlorine II 8581, 9127 doublet
Cl 2 9127 c Chlorine II 8581, 9127 doublet

Cl 2 9127 c
Cl 2 6164 c Chlorine II 6164 auroral line
Cl 2 3676 c Chlorine II 3679 auroral line
TOTL 5525 c Cl III 5519, 5539 doublet, both together
TOTL 3350 c Cl III 3354, 3344 doublet, both together

TOTL 8494 c Cl III 8504, 8436, 8552, 8483 multiplet, all together
Cl 3 5538 i Cl III 5538
Cl 3 5518 i Cl III 5518
Cl 3 3354 i Cl III 3354
Cl 3 3344 i Cl III 3344

Cl 3 8504 i Cl III 8504
Cl 3 8436 i Cl III 8436
Cl 3 8552 i Cl III 8552
Cl 3 8483 i Cl III 8483
Cl 4 8047 c ClIV 8047

Cl 4 7532 c ClIV 7532
Cl 4 3119 c ClIV 3119
Cl 4 5324 c ClIV 5324
Cl 4 5324 c
ClRr 4 i Cl 17 ly a recombination 3.7A from fully stripped ion

Ar 3 7135 c Argon III 7135
Ar 3 7751 c Argon III 7751
Ar 3 5192 c Argon III 5192
Ar 3 3109 c Argon III 3109
Ar 3 3005 c Argon III 3005

TOTL 4725 i Argon IV 4711 + 4740 together, 4740=90%
TOTL 2860 i [ArIV] 2868, 2854 together

```

```

TOTL 7250 i [AvIV] auroral lines, 7237, 7331, 7171, 7263
Ar 4 4740 c [Ar IV] 4740
Ar 4 4711 c [Ar IV] 4711

Ar 4 2868 c [Ar IV] 2868
Ar 4 2854 c [Ar IV] 2854
Ar 4 7263 c [Ar IV] 7263
Ar 4 7171 c [Ar IV] 7171
Ar 4 7331 c [Ar IV] 7331

Ar 4 7237 c [Ar IV] 7237
Ar 5 7005 c Argon V, 3P lines, 7005, collisionally excited
Ar 5 6435 c Argon V, 3P lines, 6435, collisionally excited
Ar 5 6435 c
Ar14 4413 c Ar XIV 4413, predicted lambda, not observed(??)

Ar15 409 c collisionally excited
ArRr 4 i these lines added to outlin in metdif - following must be false
K19r 4 i these lines added to outlin in metdif - following must be false
Ca 2 3933 c coll excit calcium k+h
Ca 2 8579 c infrared triplet

Ca 2 7306 c forbidden lines, 7291+7324 together
Phot 3933 i fraction H Ly-alpha destruction of excited levels
Phot 7306 i fraction H Ly-alpha destruction of excited levels
Ca2K 3934 i individual lines from five level atom
Ca2H 3969 i individual lines from five level atom

Ca2X 8498 i individual lines from five level atom
Ca2Y 8542 i individual lines from five level atom
Ca2Z 8662 i individual lines from five level atom
CaF1 7291 i individual lines from five level atom
CaF2 7324 i individual lines from five level atom

Rec 3933 i recombination contribution to CaII emission
Ca 5 6087 c Ca V optical and uv lines, collisional excitation, 3-level atom
Ca 5 5311 c Ca V optical and uv lines, collisional excitation, 3-level atom
Ca 5 2414 c Ca V optical and uv lines, collisional excitation, 3-level atom
Ca 5 3997 c Ca V optical and uv lines, collisional excitation, 3-level atom

Ca 7 5620 c Ca VII optical and uv lines, collisional excitation, 3-level atom
Ca 7 4941 c Ca VII optical and uv lines, collisional excitation, 3-level atom
Ca 7 2112 c Ca VII optical and uv lines, collisional excitation, 3-level atom
Ca 7 3688 c Ca VII optical and uv lines, collisional excitation, 3-level atom
CaLr 3 i these lines added to outlin in metdif - following must be false

ScLr 3 i these lines added to outlin in metdif - following must be false
Sc 2 21 c Sc II 2.08 (1-3)
Sc 2 41 c Sc II 4.1 mic (1-2)
Sc 2 42 c Sc II 4.22 (2-3)
Sc 3 3933 c Sc III 3936

Sc 6 5054 c Sc VI 5054 (1-2)
Sc 6 3592 c Sc VI 3595 (2-3)
Sc 6 2100 c Sc VI 2100 (1-3)
TiLr 3 i these lines added to outlin in metdif - following must be false
Ti 3 12 c Ti III 1.21 micron, (actually multiplet) 2-1 transition from model atom

Ti 3 9594 c Ti III 9594, 3-1 transition, (actually multiplet) from model atom
Ti 3 45 c Ti III 4.57 micron, 3-2 transition, (actually multiplet) from model atom
V Lr 3 i these lines added to outlin in metdif - following must be false
V 3 8823 c V III 8823
V 3 8507 c V III 8507

V 3 8507 c
V 4 7735 c V IV 7741 1-3
V 4 9489 c V IV 9496 2-1
V 4 42 c V IV 4.19 mic 3-2
CrLr 3 i these lines added to outlin in metdif - following must be false

Cr 3 5828 c [CrIII] multiplet blend at 5828A
Cr 4 7267 c [CrIV] 2 - 1 multiplet blend at 7272
Cr 4 6801 c [CrIV] 3 - 1 multiplet blend at 6806
Cr 5 7979 c [CrV] 2 - 1 multiplet blend at 7985
Cr 5 6577 c [CrV] 3 - 1 multiplet blend at 6582

```

5 THE EMISSION LINES

```

Cr 5      37 c  [CrV] 3 - 2 multiplet blend at 3.75 microns
MnLr      3 i  these lines added to outlin in metdif - following must be false
Fe 2 6200 i  Fe 2 the 3-2 transition of Netzer's atom
Fe 2 4300 i  Fe 2 forbidden 2-1 transition from Netzer's atom
Fe 2 2400 i  Fe 2 UV3, 3-1 transition from Netzer's atom

Fe2c      0 c  total of all UV+optical Fe 2 cooling
Fe2h      0 h
Fe 2 1100 i  1 to 6 transition of Fred's Fe 2 atom
Fe 2 1500 i  2 to 6 transition of Fred's Fe 2 atom
Fe 2 11500 i 3 to 4 transition of Fred's Fe 2 atom

Fe 2 2500 i  3 to 5 transition of Fred's Fe 2 atom
Fe 2 2300 i  4 to 6 transition of Fred's Fe 2 atom
Fe 2 8900 i  5 to 6 transition of Fred's Fe 2 atom
Fe 2      0 c  all cooling due to 16 level atom
Fe 2 166 i  Fe 2 1.664 microns 8-13

Fe 2 160 i  Fe 2 1.599 microns 7-12
Fe 2 153 i  Fe 2 1.534 microns 6-11
Fe 2 164 i  Fe 2 1.644 microns 6-10
Fe 2 128 i  Fe 2 1.279 microns 12-4
Fe 2 130 i  Fe 2 1.295 microns 11-3

Fe 2 133 i  Fe 2 1.328 microns 11-4
Fe 2 126 i  Fe 2 1.257 microns 10-1
Fe 2 132 i  Fe 2 1.321 microns 10-2
Fe 2 259 i  Fe 2 25.988 microns 2-1
Fe 2 353 i  Fe 2 35.348 microns 3-2

Fe 2 178 i  Fe 2 17.936 microns 7-6, label is 178 to be unique
Fe 2 245 i  Fe 2 24.518 microns 8-7
Fe 2 358 i  Fe 2 35.776 microns 9-8
Fe 2 181 i  Fe 2 1.810 microns 10-7
Fe 2 168 i  Fe 2 1.677 microns 11-7

Fe 2 180 i  Fe 2 1.800 microns 11-8
Fe 2 171 i  Fe 2 1.712 microns 12-8
Fe 2 179 i  Fe 2 1.798 microns 12-9
Fe 2 229 i  Fe 2 22.902 microns 11-10
Fe 2 347 i  Fe 2 34.660 microns 12-11

Fe 2 8619 i  Fe 2 8619A 14-06
Fe 2 8894 i  Fe 2 8894A 15-07
Fe 2 9229 i  Fe 2 9229A 15-08
Fe 2 9270 i  Fe 2 9270A 16-09
Fe2b      2 i  emission from lage FeII atom, integrated over band

Fe 3      0 c  sum of 3p and 3g states together
Fe 3 5270 c  Fe 3 5270, predictions from garstang et al 78
Fe 3 4658 c  Fe 3 5270, predictions from garstang et al 78
Fe 4      0 c  total cooling due to 12-level Fe 4 atom
Fe 4 3096 i  Fe 4 3096.A, 4-1 and 5-1 transitions together

Fe 4 2836 i  Fe 4 2835.7A, 6-1 transition, 4P5/2 - 6S5/2
Fe 4 2829 i  Fe 4 2829.4A, 7-1 transition, 4P3/2 - 6S5/2
Fe 4 2567 i  Fe 4 2567.6+ 2567.4. 11-1 and 12-1 transitions
Fe 4 277 i  Fe 4 2.774 microns 12-7 transition
Fe 4 271 i  Fe 4 2.714 microns 12-6 transition

Fe 4 272 i  Fe 4 2.716 microns 11-6 transition
Fe 4 281 i  Fe 4 2.806 microns 10-7 transition
Fe 4 287 i  Fe 4 2.865 microns 10-8 transition
Fe 4 284 i  Fe 4 2.836 microns 9-6 transition
Fe 5 3892 c  Fe 5 3892+3839

Fe 6      0 c  all of 2G lines together first
Fe 6 5177 c  Fe 6 5177, approximate correct
Fe 7 6087 c  [Fe 7] 6087
Fe 7 5722 c  [Fe 7] 5722
Fe 7 242 c  Fe 9 242 j=1 slower decay

Fe11 2649 c  Fe 11 2649 collisional excitation
Fe11 1467 c  Fe 11 1467 collisional excitation
Fe12 1242 c  Fe 12, 1242, 1349 together, collisional excitation
Fe12 2170 c  Fe 12, 2170, 2406 together, collisional excitation

```

```

Fe12  2568 c   Fe12 2904, 2567, 3567, 3073 together, collisional excitation

Fe14  5303 i   Fe 14 optically thin in line 344
Coll  5303 c   contribution from collisional excitation
Pump  5303 r   continuum fluorescence
      347 5303 c66 error! put this in
Fe19   592 c   Fe 19 from loulergue et al '85

Fe19  7082 c   Fe 19 from loulergue et al '85
Fe19  1118 c   Fe 19 from loulergue et al '85
Fe19  1328 c   Fe 19 from loulergue et al '85
Fe22   846 c   Fe 22 845.6A
Fe23   263 c   Fe 23 1909-like 262.6

FeKa   2 i     total intensity of K-alpha line
FeLr   2 i     recombination from fully stripped ion
TotH   2 i     total hot iron Ka; Auger "hot" iron, plus recom
AugC   2 i     Auger production of "cold" iron, less than or 17 times ionized
CoLr   1 i     these lines added to outlin in metdif - following must be false

NiLr   1 i     these lines added to outlin in metdif - following must be false
CuLr   1 i     these lines added to outlin in metdif - following must be false
ZnLr   1 i     these lines added to outlin in metdif - following must be false
Stoy   0 i     optional sum of certain emission lines, set with "print sum"

```

5.5 Atomic data sources

Codes like Cloudy can only exist because of the large body of work done by the atomic and molecular physics community. This work will only continue to be supported if it is cited in the literature whenever it is used. The following is a partial list of citations for the atomic data used within the code.

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6 Coding Conventions

Cloudy is large, complex, and the result of many hands. It is essential that clarity and integrity of purpose be sustained. This can only be achieved by having the **self-restraint** to follow a coherent set of standards. These standards are outlined in this section. All are arbitrary standards, but these are the standards Cloudy follows. It is far better to follow a single set of standards than to have total anarchy.

Cloudy is evolving towards a simple formulation of the Hungarian naming convention (Simonyi 1977). In this convention the first few characters of a variable name indicate the type and function of that variable.

6.1 Variable names and strong typing

In part the naming convention used in the code today looks back to an under appreciated advantage in the FORTRAN II and FORTRAN 66 languages - the fully implicit designation of variable types by the first letter of its name. The naming convention forced by early versions of FORTRAN (integers begin with i-n, real numbers with other characters) is still useful since the type can be determined at a glance. This is used in the following.

6.1.1 Integers

Integers begin with the characters i, j, k, l, m, or n.

Counters begin with n. Examples include *nLevel* or *nLoop*.

Loop indices are generally i, j, or k. Sometimes they are counters.

Indices within arrays begin with ip. Examples include *ipContinuum* or *ipCIV1549*.

6.1.2 Double or float variables

These begin with letters between a through h, and o through z. Examples include *PumpRate*, *DestRate*, or *CollisIoniz*.

At this time the naming convention does not distinguish between floats and doubles. Eventually the code will be totally double precision.

In some cases floating numbers naturally will have names beginning with one of the letters reserved for integers. In this case a lower case x is used as the first character. Examples include *xJumpDown*, *xMoleDen*.

6.1.3 Character strings

Character variables begin with “*ch*”. Examples are *chName* or *chReadInput*.

6.1.4 Logical variables

These begin with “*lg*”. Examples are *lgOK*, *lgDone*. These are of intrinsic type int.

6.2 Structure names

Variables with a common purpose are grouped together into structures. The electron density variable *eden* is an element of the structure *phycon* with the name *phycon.eden*. The declaration for a structure occurs within an included file with the same name ending with “.h” – the *phycon* structure is in *phycon.h*.

6.3 Braces

The format of braces consumes a staggering amount of debate but is important since the format must be followed consistently across the code for it to be instantly legible. There are three major styles of braces:

Style 1;

```
if( a>0 )
{
    b = 0.;
}
else
{
    b = 1.;
}
```

Style 2:

```
if( a>0 )
{
    {
        b = 0.;
    }
}
else
{
    {
        b = 1.;
    }
}
```

Style 3:

```
if( a>0 ) {
    b = 0.; }
else {
    b = 1.; }
```

The code uses the first style. Any one of the three could have been chosen, but the first one was chosen. We must have the **self-restraint** to follow this arbitrary choice, for the clarity of the overall code.

6.4 Changes to the code

A comment line just before the affected line indicates changes to logical flow within the code that could impact results or convergence. These have the following style:

```
/* >>chng 95 dec 20 eden had eold, was undefined here, affect electron density */
```

The flag **>>chng yy mmm dd** indicates a change. Here **yy** is the last two digits of the year, **mmm** is a 3-character abbreviation of the month, all in lower case, and **dd** is a two-digit date. It is important that this style be followed consistently so that changes can be searched from the code with a pattern matcher such as **grep**, and then sorted by date.

6.5 Atomic data references

Codes such as Cloudy only exist because of the foundation of basic atomic and molecular data. It is important to the survival of this field that the original sources of the basic data be cited, since this in turn affects their ability to generate support. The code follows the convention of preceding all uses of atomic data with a citation to the original paper in the following form:

```
/* >>refer Si+2 AS Berrington, K., AtData Nuc Data Tab 33, 195.
```

The flag “>>refer” indicates a reference. If the reference cannot fit on a single line it may continue on the following line, starting with the flag “>>refercon”. This style must be followed consistently so that a pattern search will generate a list of references used.

6.6 Sanity checks and asserts

Sanity checks and asserts are redundant tests for variable values that are totally impossible (Maguire 1993). Examples include negative collision strengths or electron temperatures. A major improvement to Cloudy version 86 and later is the inclusion of large numbers of sanity checks, while in version 94 and later C assert macros were used. These checks do not have a major impact on performance but they do slow the code down a bit. For production runs with a gold version of the code it would be reasonable to not include these checks.

The asserts can be neglected, and the code run slightly faster without this checking, by defining the macro NDEBUB on the compiler command line, as in

```
cc -DNDEBUB file.c
```

6.7 Broken and Test Code

Broken() It is sometimes necessary to physically break the code, either by writing specific code to override the correct behavior or disable a physical process. When such code is entered, it should be accompanied by a call to routine **broken**. This routine does nothing but set a flag that broken code is present. This flag generates a warning after the calculation is complete, to serve as a reminder of the presence of the broken code. This routine is not normally used.

TestCode() Trial code is identified by a call to routine **TestCode**. This routine does nothing but set a flag that test code is present. This flag generates a comment after the calculation is complete, to serve as a reminder of the presence of the test code. This routine is not normally used.

fixit() sets a flag saying that the code needs to be fixed.

6.8 Version numbers

Cloudy uses version numbers to keep track of changes to the code. The version number is stored in **version**, in the structure **date**. The variable **chDate**

contains the date of the last major revision, and the variable ***chVersion*** is a string giving the version of the code.

7 PROBLEMS

7.1 Overview

This section describes some of the errors that can cause Cloudy to stop. Floating point errors should never occur. Several other internal errors, which the code is designed to catch and then complain about, can occur. Finally, it is possible that the code will stop because of thermal stability problems.

The most important single thing to understand about any calculation is why it stopped, and whether this affects the predictions. This is discussed further in the section ***Stopping Criteria*** in Part I of this document.

7.2 Thermal stability and temperature convergence

This section describes thermal stability problems, how to identify them, and what to do about them.

7.2.1 Types of thermal maps

Three types of thermal maps, showing the heating or cooling of gas as a function of temperature, can be produced by Cloudy. Each is the answer to a different question.

Figure 4 shows the heating and cooling rates as a function of temperature for a photoionized gas in which the electron temperature was varied. This figure was produced by running the test case `map.in`, one of the standard test cases included

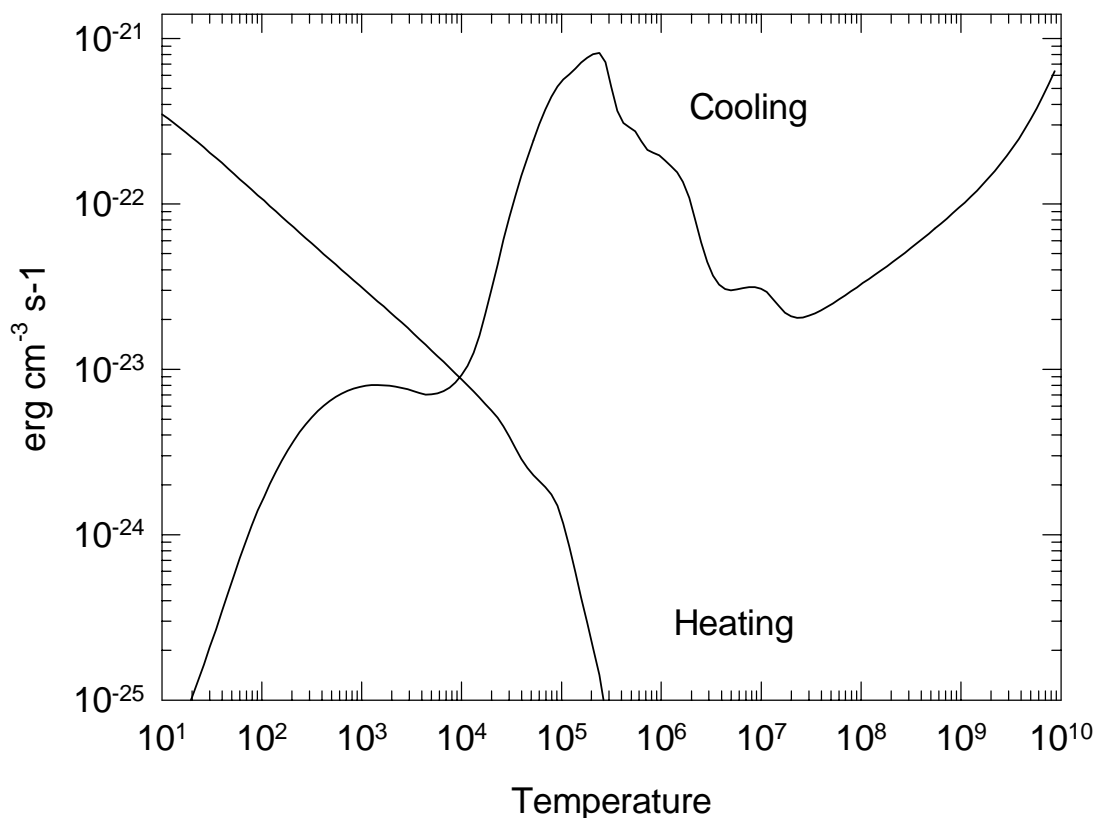


Figure 4 A typical cooling function for low density photoionized gas. The cooling and heating rates ($\text{erg cm}^{-3} \text{s}^{-1}$) are shown. cooling

in the code distribution. Both the gas density and the flux of ionizing photons was held constant, so only one temperature is meaningful. The `map.in` file uses the `punch map` command to determine heating and cooling rates at a variety of temperatures. This is exactly what the code does to determine the equilibrium temperature, so this plot can be useful to find out why the code ran into problems. This is why the command was introduced. Note that the only valid solution, and the only one with physical meaning, is the one where heating and cooling match – the others are simple bookkeeping exercises.

Gas in collisional equilibrium has a well-defined cooling rate that is only a function of temperature. The sample program `coolcurve.c` (included in the source distribution) does such a calculation, and Figure 5 shows it. Here the electron temperature is set by some physics external to the problem. Each temperature, and the entire ionization solution, is valid for each temperature, under this assumption. The unspecified heat source would have to provide a local heating rate that is equal to the calculated cooling rate for the solution to be time steady.

The third map is the type of thermal stability map shown by Krolik, McKee, and Tarter (1981) and plotted in Figure 6. The program that generated these results is given in the file `kmt.c`. Here the equilibrium temperature is determined self-consistently for gas over a wide range of densities, but for a single flux of ionizing photons (or equivalently, distance from the central object).

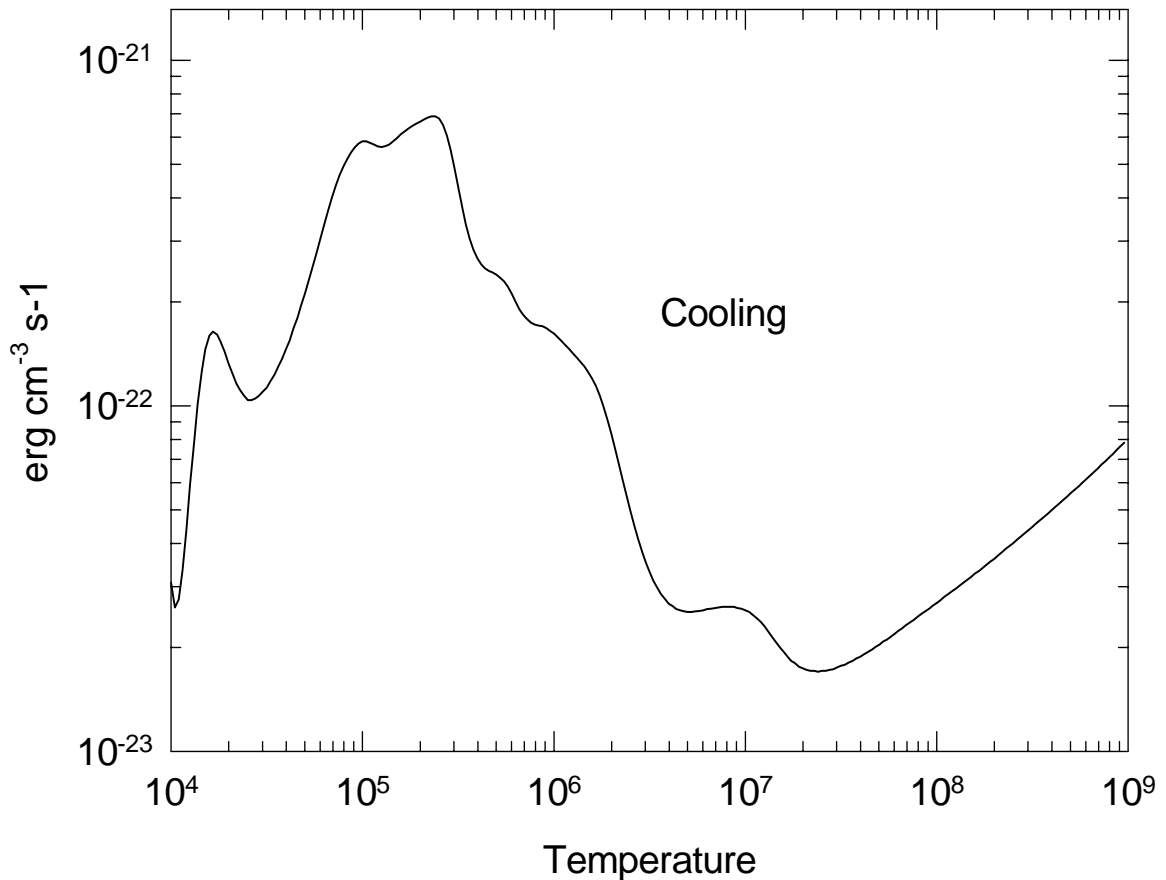


Figure 5 A typical cooling function for low density collisionally ionized gas. coolcurve

7.2.2 No Temperature Convergence

A temperature failure occurs when the heating-cooling balance is not within a certain tolerance, set by the `tolerance` command, after 20 tries. Normally Cloudy will punt after an excessive number of temperature failures occur. The limit to the number of failures is reset with the `failures` command, which sets the variable

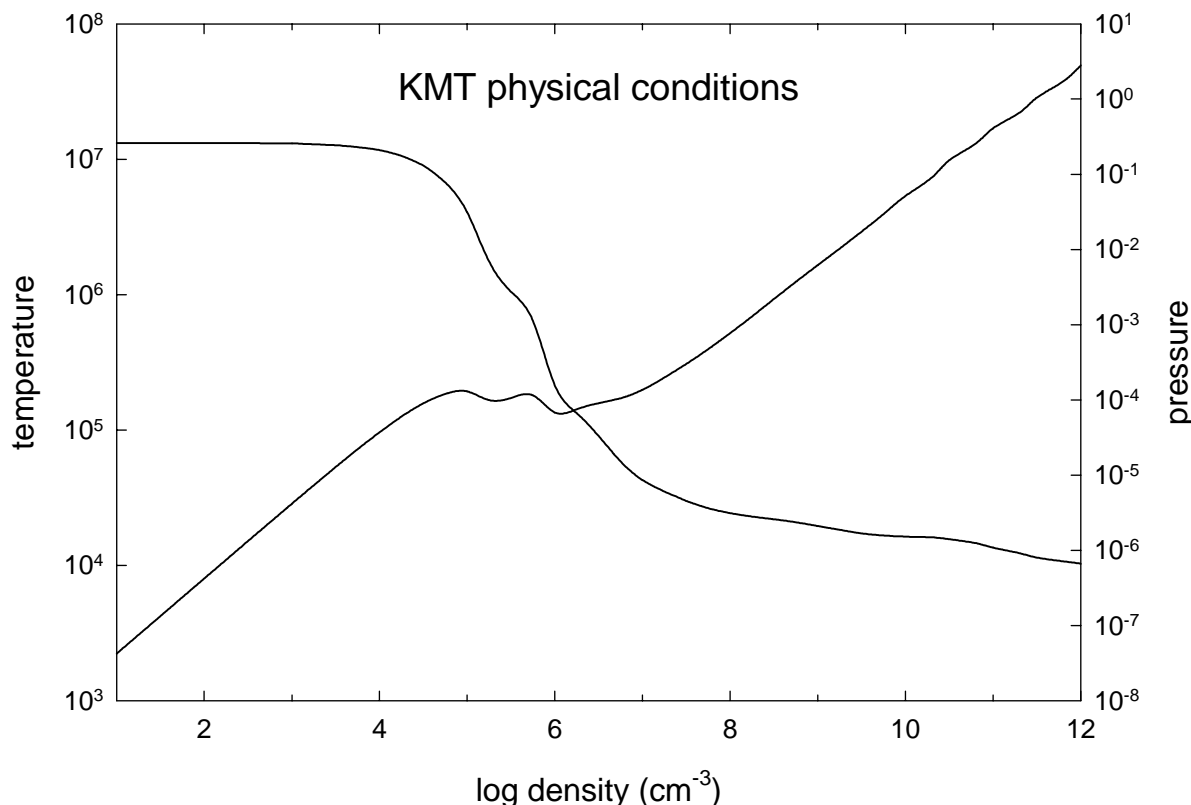


Figure 6 Equilibrium temperate as a function of density. Kmtplot

limfal. The default value is 20. (When Cloudy stops because of excessive failures it first produces a map of heating-cooling space to give an indication of where the equilibrium temperature should have been.)

Temperature failures most often occur for temperatures in the range 1 to 4×10^3 K, and 10^5 to 10^6 K. These are ranges where the cooling function permits more than one thermal solution (see, for example, Williams 1967). Figure 4 shows a typical cooling function for gas in photoionization equilibrium.

A peak is reached at a temperature near 10^3 K. This can occur when the fine-structure lines are major coolants. At lower temperatures their cooling rate goes up exponentially (as expected), until roughly 10^3 K, when their Boltzmann factors are near unity. Above this temperature their cooling rate is nearly proportional to the Coulomb focusing factor $T^{-1/2}$, and the cooling *decreases* until the temperature is high enough for optical forbidden lines to become important (at roughly 4000 K). A similar phenomenon occurs near the $\sim 10^5$ to 10^6 K peak in the cooling function.

When failures occur because more than one temperature solution is possible, the reported failures are a physical (not numerical) problem. Cloudy will try to deal with this problem by forcing the temperature to values below the peak in the cooling

function. Increasing the number of allowed failures (with the **failures** command) to prevent the code from stopping prematurely is permissible as long as the global energy balance is preserved. A warning will be issued at the end of the calculation if the heating-cooling balance is not preserved.

7.2.3 Thermal Stability

The thermal solution may be unstable when the temperature derivative of the net cooling function (cooling minus heating) is negative (Field 1965). Possibly unstable solutions are indicated by a “u” just before the equilibrium temperature in the zone printout. (The temperature derivative is for isochoric, not isobaric, conditions.) Comments are printed at the end of the calculation if possibly unstable thermal solutions are present in the calculation.

7.2.4 Thermal fronts

Just as an ionization front is a region where the level of ionization changes dramatically over a small scale, a thermal front occurs where the temperature changes dramatically over a small scale. This can be caused by a real physical change of state of the gas such as those that occur near the peaks in the cooling curve. In reality electron conduction may step in as a significant heating/cooling agent, although Cage & Seaton (19xx) argue that this is not important. This could become a problem in the current implementation since electron conduction is not included. There is no obvious solution to this type of jump since it is a physical effect. Jumps can also occur if the grid zoning becomes too coarse for changes in the physical conditions, although the code should be protected against this.

The code will generate a caution or comment if the electron temperature changes discontinuously from one zone to the next.

7.2.5 Map Output

If an excessive number of temperature failures occur (the default limit is 20) then the program stops and produces a map of the heating and cooling as a function of temperature for the last computed zone. The limit to the number of failures allowed before the code punts is reset with the **failures** command. The map is described here. The start of the output from the test case **map** is shown below.

```
90.02x map of heating vs cooling
te, heating, cooling.
Cloudy punts, Te= 9.254E+03 HTOT= 9.123E-24 CTOT= 9.118E-24 nzone= 1
COOLING array is
O 4 25 0.340 O 3 5007 0.182 O 3 88 0.075 H FB 0 0.057 S 4 10 0.048 O 3 51 0.042 S 3 9532
0.035 H ff 0 0.022 S 3 33 0.020 Ne 3 15 0.019 Hefb 0 0.015 N 3 57 0.015 Ne 3 3869 0.013 S 3 18
0.013 Ne 5 24 0.010 Ne 5 14 0.009 C 3 1910 0.008 Heff 0 0.007 Si 2 34 0.006 Fe 5 3892 0.006 O 2 3727
0.005
Line heating array follows
Te Heat-----> Cool-----> dH/dT dC/DT Ne NH HII Helium
1.0000E+01 3.4774E-22 1 1 0.636 4.6095E-26 H FB 0A 0.723 -8.19E-24 1.56E-27 9.1178E-01 1.0000E+00 -0.07 -0.40 -0.24 -1.75
1.0209E+01 3.4490E-22 1 1 0.635 4.6814E-26 H FB 0A 0.720 -7.98E-24 1.65E-27 9.1353E-01 1.0000E+00 -0.07 -0.40 -0.23 -1.73
1.0423E+01 3.4233E-22 1 1 0.635 4.7510E-26 H FB 0A 0.717 -7.74E-24 1.74E-27 9.1491E-01 1.0000E+00 -0.07 -0.41 -0.23 -1.73
```

The output begins with a listing of the strongest coolants. Then the program steps through increasing temperatures and prints the heating, cooling, and ionization of the gas. From this information it should be possible to determine the temperature where the equilibrium thermal solution should have been. Each solution is completely self-consistent (except that heating and cooling do not balance). Both the local attenuated radiation field and collisional ionization contribute to the ionization

balance at each temperature. All processes contribute to the thermal balance, including collisional ionization. The map is at constant density.

The first column gives the temperature. Columns 2 and 6 give the volume heating and cooling. Both have units $\text{erg s}^{-1} \text{cm}^{-3}$. Columns 3 and 4 constitute an indication of the main heating source. Columns 7 and 8 give the label and wavelength of the strongest coolant. Columns 5 and 9 give the fraction of the total heating or cooling due to these agents. Columns 10 and 11 give the heating and cooling derivatives. Columns 12 and 13 give the electron and hydrogen densities (cm^{-3}) and the remaining columns give the logs of the hydrogen and helium ionization fractions. The location of the probable thermal solution is indicated by a comment surrounded by dashed lines.

7.3 Floating Point Errors

The code should be compiled and linked with options enabled so that the code will crash on overflow or division by 0, but ignore underflow. *Floating point errors should never occur.* The logic within the code is designed to identify problems, and complain, but not fail. The logic is only as good as the tests they were designed to pass. It is inevitable that circumstances will occur for which the logic now in the code is not sufficient. It is possible that the code will fail when these circumstances occur. I would be grateful for reports of any such failures, since they inevitably identify shortcomings in the code, and lead to its improvement. My email address is gary@cloud9.pa.uky.edu.

7.4 Optical depth convergence problems

The code generally will not converge if it has not done so within ten or so iterations. Convergence problems most commonly occur when the specified column density or thickness is very near a prominent ionization front. In this case very small changes in the physical conditions results in large changes in the optical depths. The code will not have convergence problems if an optical depth is used as a stopping criterion instead.

7.5 Negative Populations of H, He

It is possible that the code will stop because negative level populations were predicted for atoms or ions of hydrogen and helium. This is not supposed to occur, but sometimes happens because of numerical instabilities in the matrix inversion routine. Please send me the input stream and version of Cloudy.

7.6 I can't fix it if I don't know its broken.

Machines are growing faster far more rapidly than people are getting smarter. Reliability in the face of complexity is the major challenge to the development of any large-scale computer code. There can be little doubt that Cloudy contains bugs.

The code is well tested in many simple limits, and behaves in the correct manner. Simulations of H II regions, planetary nebulae, and other simple objects, are in good agreement with predictions of other photoionization codes (Ferland et al. 1995).

Bugs can be discovered by strange behavior in situations where the code has not been well-tested. The discovery of the existence of problems is itself a major challenge. If problems arise or the code crashes then it is likely that a problem has been isolated. I would appreciate learning about such problems since they identify shortcomings which usually lead to improvements in the code (or the documentation). My email address is gary@cloud9.pa.uky.edu.

8 REVISIONS TO Cloudy

8.1 Overview

This section outlines some of the major versions of Cloudy, and gives an indication of the direction development will take in the next few years. Its development began in August of 1978, at the Institute of Astronomy, Cambridge, and has been continued at The University of Kentucky, The Ohio State University, and during extended visits to the Joint Institute for Laboratory Astrophysics, the Royal Greenwich Observatory, Cerro Tololo Interamerican Observatory, and the Canadian Institute for Theoretical Astrophysics. Figure 7 shows the evolution of the code, as indicated by its size as a function of time⁶.

8.2 Cloudy and Moore's Law

Moore's Law is due to Gordon Moore, one of the founders of Intel Corporation. He observed that modern CPU's become about twice as powerful every 18 months. This trend has held true for the past twenty years, shows no sign of failing, and seems to be associated with the control of complexity.

By this standard the growth of Cloudy has been conservative, in that it is growing slower and complex on the same timescale. As an example, the Meudon 1985 Meeting planetary nebula test (`parispn.in` in the distributed test cases) has always

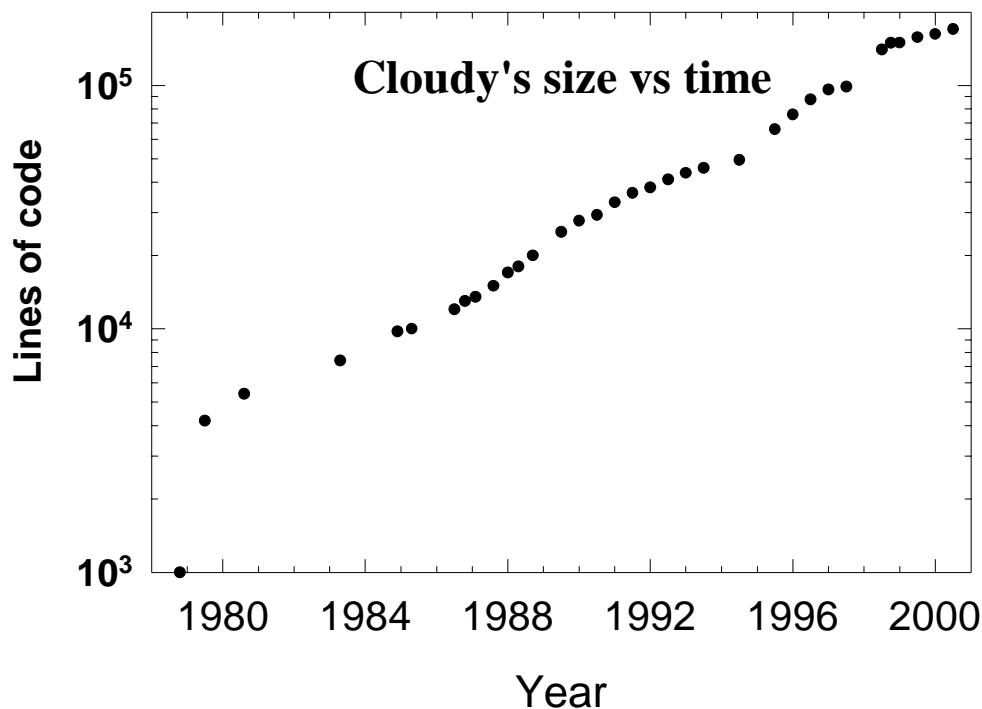


Figure 7 The size of the code, indicated by the number of lines of executable Fortran, as a function of time. size

⁶ Before mid-1995 the size was the total number of lines in the distributed source. After 1995 the size only includes the number of lines excluding block datas. The number of lines of block data

taken about one minute to compute.

8.3 Major Past Versions

- 67 August 1987. Hydrogen atom goes to LTE in limit of large electron densities. Many small bugs uncovered as result of careful comparison with Netzer's ION.
- 68 Cambridge, Fall 1987. Development work in progress.
- 69 December 1987. Hydrogen atom goes to LTE in limit of large photon densities. Ferland and Rees (1988).
- 70 September 1988. He II $\text{L}\alpha$ transfer improved. Improved form of escape probabilities with explicit damping constants. H^- and improved free-free heating. Many high excitation metal lines transferred. (Rees, Netzer, and Ferland 1989; Ferland and Persson 1989).
- 71 December 1988. Photon array rewritten, now Compton exchange problem is exact for black bodies with temperatures between 2.7 K and 10^{10} K. He II $\text{L}\alpha$ radiation pressure included.
- 72 January 1989. Static version, minor bug fixes.
- 73 1989. Major rewrite of helium treatment. Dust changed to two populations, scattering and absorption included. Default radius and thickness increased by ten orders of magnitude. He^+ goes to LTE. Development work on getting helium to go to LTE. IR power law for default AGN continuum now $\nu^{+2.5}$ below 100 micron break. No **anumm** array, all continuum one array. **table star**. **coolr** broken up. Helium to LTE for high electron density.
- 74 1990 January. Hydrogen double precision, many bug fixes. 10 tables among the continua. Kurucz (1989) atmospheres. Improved dust treatment, including photoionization and charge.
- 75 1990, JILA visit. Cosmic abundances changed to Grevesse and Anders. Major bug in constant pressure for HII regions, PNs, etc; did not affect BLR. Calculation of νf_{ν} ($\text{H}\beta$) was incorrect. Molecules at low temperatures. HDEN now $n(\text{H}^0) + n(\text{H}^+) + n(\text{H}^-) + 2n(\text{H}_2) + 2n(\text{H}_2^+)$. Improved Rayleigh scattering treatment. Dielectronic recombination for sulfur (guess). Many changes in dust; Orion paper (Baldwin et al. 1990). Subordinate lines changed to Hummer's K2 function.
- 76 1990, static version from end of JILA visit.
- 77 through Nov. 1990. **optimize** option added using Bob Carswell's code. Gaunt factor for brems input spectrum. Reflected continuum predicted. Frequency partition adjusted. X-Ray optical depth now at 0.5 keV. Read in table of points from previous calculation. **opsav** deleted, now single pointer **opsv** for all opacities. Numerical array to 100 MeV. Hummer $\text{L}\alpha$ escape destruction prob. **tautot** arrays now both in and outward directions. Bound

became very large, thanks to Dima and Katya Verner, so that the total size is roughly 2.5 times the indicated size.

Compton included for all ionization levels. Mean ionization arrays rewritten to make sense. Bug in wind velocity fixed, result now exact. C, O outward diffuse fields changed to OTS.

- 78 through May 91. Continuum escape probability formal H-only opacity. H recombination, cooling over wide range of temp. X-Ray optical depth back to 1 keV. **abundances no dust** no longer changes abundances of depleted elements. OI-L β treatment is now six-level atom. Default table AGN changed. Continuum normalization rewritten. Milne relation for diffuse fields of H, all He. Fe K α divided into hot and cold. Beams paper (Ferland, Peterson, Horne, et al. 1992) Many high ionization lines included as OTS and outward ionization sources.
- 79 Summer 91. H molecules completed, C, N, O molecules included. Continuum binning changed for Ca, Fe L-shell ionization potentials. Extensive testing.
- 80 July 91, static version .05. Version 80.06 fixed many small problems discovered by several people. New collision strength for [NeV] put in. This static version ended with 80.09, in January of 1993.
- 81 Late 1991. New collision strengths for [NeV] IR lines (10x larger). Also for [CII], [NIII], and [OIV]. Some are 2x larger. Opacity arrays totally rewritten with eye to Opacity Project data. NIII paper (Ferland 1992).
- 82 Early 1992. X-Ray opacity arrays rewritten. Improved pressure convergence. HJBAR now function. Error in cooling due to collisional ionization of H, He. Sodium and nickel added. /TAU/ array broken up. All heavy element opacities converted to table look-up. Revised collision strengths for NIII lines, improved treatment of atom. NLR abundances deleted, ISM put in their place. Summary comments now driven by subroutine. Cap on 911 OTS field. Note on [FeXI] maser (Ferland 1993). Luminosity command separated into luminosity and intensity commands.
- 83 Autumn 1992, Cambridge and CTIO visits. Hydrogen molecule network completed, Ferland, Fabian, and Johnstone (1993). Collision strengths for fine structure lines changed to Blum and Pradhan 1992, Hollenbach and McKee 89; these changed temperatures for cold ISM by factors of 2. Heavy element molecule network as in Hollenbach and McKee 1989. Code works in fully molecular limit. Kevin Volk's stars (Atlas 91, and Werner models). More accurate treatment of secondary ionization after Voit visit. [OI] lines each include escape prob. Opacities, destruction probabilities, evaluated within all loops, code far more stable, but roughly three times slower. Transferred HeI 2.06 line correctly, after Shields papers.
- 84 1993 Feb 13, Static version following CTIO visit.
- 85 1994, Revisions following Lexington meeting. Outward only now default continuum transport.
- 86 1995, All of first thirty elements are now in code. Photoionization database changed to Dima Verner's *phfit*.

- 87 1995 summer, Map now converges electron densities. Negative populations of OI and FeII atoms solved. Dima's 6k lines included in cooling and radiative acceleration. Total rewrite of *nextdr* logic. Ionization predictor corrector logic completely rewritten. Dima's *phfit* now fitted to all Opacity Project data.
- 88 1995 Fall-winter. Kirk's extensive grids of BLR models run. March 1996 visit to Tel Aviv to compare results with Hagai. Kirk carefully went over atomic data base. Iron recombination changed to Arnaud and Raymond (1992). Default iron abundance down 33%. Entire line data base revised.
- 89 1995 winter-spring. These were mainly beta versions with no major changes, but many fixes to problems.
- 90 1996 June 17, static version, extensive year of debugging. Ferland et al. (C90) paper.
- 91 1977, preparation of C version.
- 92 1998, sabbatical at CITA, revision and debugging of C version.
- 94 1999 December 24, gold version of C code.
- 95 2001 February, He isoelectronic sequence, CO multilevel molecule, pgrains distributed grains, atomioc data update.

8.4 Version 95 vs 94

8.4.1 Commands

The **atom** command is now the primary method of changing treatments of H-like, he-like, and molecular species.

8.4.2 Physics

The helium isoelectronic sequence has been broken out.

¹²CO and ¹³CO are now multi-level molecules with the full rotation spectrum predicted.

The **pgrains** treatment of grain physics has been incorporated, allowing fully resolved treatment of grains charging, collisions, and emission.

Two photon emission, and induced two photon emission, is fully treated for the complete H and He iso-electronic sequences.

8.4.3 Miscellaneous

The ISM oxygen abundance was changed from 5.01e-4 to 3.19e-4 as recommended by Meyers et al. (1987).

8.5 Version 94 versus 90

8.5.1 Commands

The **hydrogen** and **feii** commands have been combined into the **atom** command, which has many options.

8.5.2 Hydrogen

The full hydrogenic isoelectronic sequence is now treated with a single model atom and code base. The model atom can have up to 400 levels.

8.5.3 grains

Parameters used in the treatment of the old-style grains have been revised as per the Weingartner & Draine (2000) paper. The **pgrains** command makes it possible to resolve the grain size distribution function, solving for grain properties and emission as a function of their size. This is described in van Hoof et al. (2000).

8.5.4 Other changes

The code is now ANSI 1989 C, making it especially gcc and Linux friendly.

8.6 Version 90 versus 84

8.6.1 Commands

The **abundances** command now needs 29 numbers by default. A new command “**init**” allows a commonly used set of commands to be saved as a single file and used by a variety of scripts.

8.6.2 Continuum Transport

Versions 86 and before used a modified version of on-the-spot approximation (OTS) for the Lyman continua of hydrogen and helium. This method was numerically stable and gave results in excellent agreement with Van Blerkom and Hummer (1967). This has been changed to outward-only to obtain better agreement with predictions of Pat Harrington’s and Bob Rubin’s codes (Ferland et al. 1995). The OTS code is still in place and will be used if the **diffuse ots** command is entered, but outward-only is the default. The two methods result in temperatures at the illuminated face which can differ by as much as several thousand degrees, but the resulting spectra are surprisingly similar.

8.6.3 Hydrogen

The model hydrogen atom has been generalized to an arbitrary multi-level atom (Ferguson and Ferland 1996). The **hydrogen levels** command is used to specify the number of levels to be used. The collision strengths have been changed to Vriens and Smeets (1980) for levels higher than 3, and Callaway (1994) for collisions with 1,2 and 3.

Predicted infrared line intensities are now correct for all densities and temperatures greater than 10^3 K. Versions before 89 used a well l-mixed hydrogen atom, and its predictions were not correct for some infrared lines at low densities.

The routine that computes the free-free gaunt factors has been extended to include the full range the code can handle.

8.6.4 The helium ion

The helium ionization balance at low photon and particle densities, and at high particle densities, has always been exact, and this continues to be the case. There was a problem in the helium ion for high radiation densities, in versions 87 and before. The code used three pseudo levels to represent the levels between 7 and 1000, for H, He, and He⁺. This seemed to work well for the atoms for the cases of high densities, but testing has shown that it did not represent the physics of the high radiation density limit well. The problem is that the pseudo-levels had very large statistical weights, they represented line energies in the far infrared, and had A’s appropriate

for lower levels. As a result they had very large induced rates when the photon occupation numbers were large, and this affected populations of lower levels. As a result the atom became too ionized - as much as a factor of two for He⁺. The following test illustrates this problem:

```
title helium ionization in high photon density limit
print departure coef
set dr 0
stop zone 1
constant temper 4
hden 11.000
phi(h) 20.750 range 1
stop thickness 11.7
table agn
```

In versions 88 and later no pseudo levels are used for any hydrogen or helium atom or ion.

8.6.5 Heavy elements

The atomic data base, the organization of aspects of the code dealing with storing heavy element ionization, and all the associated routines, have been totally rewritten. The lightest 30 elements are now included. Photoionization data are from Verner et al. (1996), recombination data partially from Verner and Ferland (1996), and roughly 10⁴ lines of the heavy elements have been added (Verner, Verner, and Ferland 1996).

The number of resonance lines has increased by more than an order of magnitude. All resonance lines listed by Verner, Verner, and Ferland (1996) are included. As a result of these many additional lines the cooling function tends to be larger and smoother.

All lines are now fully transferred, and include pumping by the attenuated incident continuum as a general excitation mechanism. Pumping can be a significant contributor to the formation of weak high excitation lines.

The default solar mixture has been changed to Grevesse and Noel (1993). The biggest change is in the iron abundance. Previous versions had used a higher photospheric abundance. The current version is the 1993 suggested meteoritic abundance.

8.6.6 Free-free, line heating and cooling

These are counted in a different but equivalent manner. Now the *difference* between cooling and heating is used, since this is more numerically stable at high radiation densities. This difference has no physical affect on the predictions, but the printed contributors to the total heating and cooling do appear different.

8.6.7 Excited state photoionization cross sections

OP data are now used. For the excited state of Mg⁺ this is nearly ten times smaller than old screened hydrogenic values. This affects the intensity of Mg II λ 2798 in some BLR calculations.

8.6.8 The O^+ photoionization cross section

The Reilman and Manson photoionization cross sections, used before version 87, show a jump in the photoionization cross section at the $2s - 2p$ edge, and low values above that threshold extending up to the valence electron threshold.

Opacity Project cross sections are used in the current version of the code, and these do not show the $2s$ edge (the OP calculations find that the $2s$ and $2p$ electrons are highly correlated). The cross section remains large up to the valence threshold. The difference approaches a factor of two, and this affects high ionization parameter clouds since O^+ is the dominant opacity for some energies.

Verner, et al. (1996) comment on all other cases where the photoionization cross sections have changed. There are generally atoms and first ions where Opacity Project data are now available.

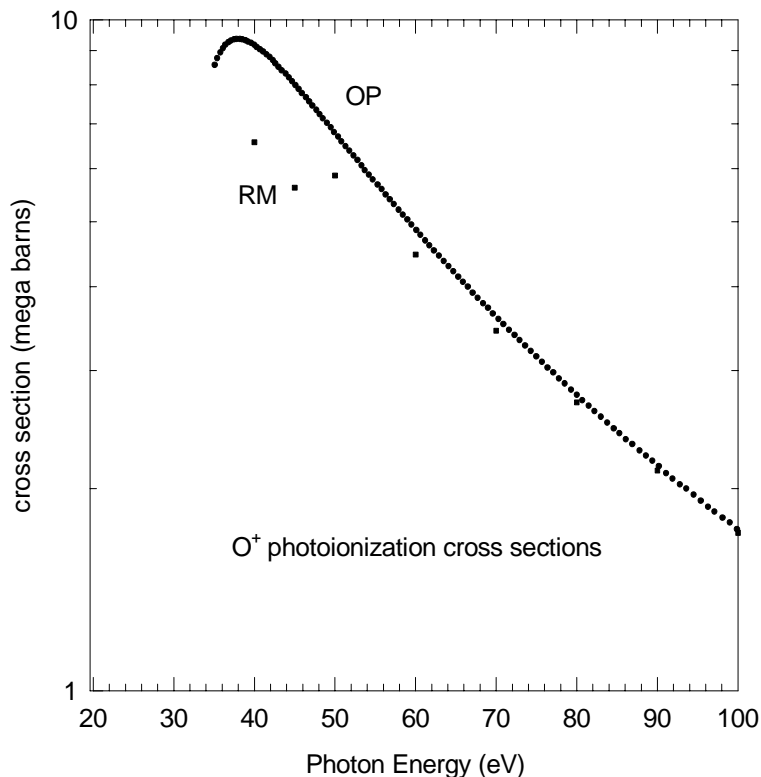


Figure 8 This figure shows the O^+ photoionization cross sections now used, compared to the Reilman and Manson values, used in version 84 and before. o2photo

8.7 Version 84 versus 80

8.7.1 Commands

In previous versions of the code both luminosity (quantity radiated into 4π sr) and intensity (a surface flux) were specified with the same command. The code decided which was intended by checking the resulting ionization parameter. This method never failed to the best of my knowledge, but, as the code grows more capable of considering ever more extreme cases, there might eventually come a time when it made the wrong decision. All luminosity-intensity commands have now been split up. For instance, the $Q(H)$ command is now two commands, $Q(H)$ (for number of photons radiated into 4π sr) and $\phi(H)$, for the surface flux. These commands are all discussed in Part I of HAZY.

8.7.2 Mg II $\lambda 2798$

The biggest difference between the two versions is in the predicted intensity of Mg II $\lambda 2798$. The intensity of this line is now a factor of two stronger in many models. The new version uses L-shell photoabsorption cross sections from Reilman and Manson (1979), and the older version used inner shell cross sections extrapolated from the table in Weisheit (1974). The cross sections differ by a factor of nearly 3, in the sense that Mg now tends to be more neutral, and Mg II stronger. As a result of the increased cooling by $\lambda 2798$ other lines formed in the same region tend to be weaker.

8.7.3 General Results

The following may affect certain specific models, but did not result in changes in any of the “standard” test cases.

The treatment of line-continuum fluorescence in the optically thin limit has been much improved, following Ferland (1992). This can affect hydrogen line emission in clouds that are optically thin in the Lyman continuum.

The treatment of molecules has been vastly improved. The code now goes to the fully molecular (H_2 and CO) limit, and reproduces the Tielens and Hollenbach (1985a, b) PDR results for Orion.

The elements Na and Ni have been added.

8.8 Known Modes for Cloudy 95

- Time dependent model turned off.
- Molecular abundances can become negative for very large CO fractions.

8.9 Making a Revision

8.9.1 The code

- Compile code with array bounds checking. Run all test cases.
- Run all test cases on Alpha, HP, Sparc, SGI, and PC.
- Confirm *fabden* stops immediately.
- Verify `aaaa.c` has correct version.
- Summarize changes in `c95rev.htm` in the web site. Use comments in `version.c`. This is the primary documentation for changes to the code.
- Create the test case files for each of the platforms. Each lives in a separate subdirectory off the hazy directory. Create a compressed tar file with the name of the platform included. Place this in the ftp site with the script `cp.ftp`.
- Tar all input files into single file with name `tests.tar`. The command is `tar -cvf tests.tar *.in *.for`. Compress and copy to ftp.

8.9.2 Printing Hazy

- **Compile the line list from `lines.c` and bring these into HAZY.** To do this, go to the `programs` subdirectory off the `hazy` main directory, and run `do.all` to create formatted output file call `lines.out` for HAZY. Edit the links in HAZY `lines.doc` to update this.
- **Create the list of transferred lines.** Compile the line data list for the transferred lines by running the input file `punchline.in` in the Cloudy main directory. This will create a punch file named `LineData.txt`. This is the file that will be linked into HAZY. Edit the links in HAZY’s file `lines.doc` to update the list of transferred lines.

- ***Update the list of subroutine names.*** Do this by running `do.all` in the `include` directory. The list of routine names in is the file `routines.txt`, and this is automatically generated by `do.all`. This was probably done to update the source, as described in the previous section. Edit links in HAZY `routines.doc` to update this.
- ***Update all comparison tables.*** These are in the section starting on page 475. Go over these to confirm that line predictions are OK.
- ***Update all test case input scripts.*** These are listed in the section beginning on page 492 of this document. Edit the links to manually update them.
- ***Confirm that all cross-linked variables are ok.*** Change labels in HAZY `headinfo.doc`.
- ***Summarize changes to the code.*** These are listed in `version.for` and should go into the past major revisions in this section of HAZY.

9 COMPARISON CALCULATIONS

9.1 Overview

This section presents comparisons between the current predictions of the code, and results from other independent calculations. The “other” calculations are from the compendium resulting from the Lexington meeting on model nebulae (Ferland et al. 1995).

The scatter among the calculations, as well as the changes that have occurred in the predictions made by Cloudy, are in some sense an indication of the stability and reliability of these types of extreme non-LTE calculations. The largest discrepancies between current predictions made by Cloudy and the other models from the Meudon meeting (which were computed in 1985) are due to changes which have occurred in the atomic data base between 1985 and the present. In general, the strongest lines are in very good agreement (as they must because of energy conservation) while weak lines (which are very sensitive to changes in the computed temperature and ionization structure) scatter by nearly a factor of two.

9.2 Cool HII Region

This is an HII region ionized by a very cool star. It is one of the Lexington Meeting test cases and is computed with the input script `coolhii.in` in the code's test suite. This is the simplest model since helium is predominantly neutral. The entry ***L(total)*** comes from the “Stoy” printed entry.

Table 1 HII Region Ionized by Cool Star

		Mean	Ferland	Harrington	Netzer	Pequignot	Rubin
L(H β)	E36	4.90	4.98	4.93	4.85	4.83	4.93
[NII]	6584+	0.87	0.91	0.82	0.97	0.82	0.84
[OII]	3727+	1.21	1.16	1.22	1.32	1.14	1.21
[NeII]	12.8 μ	0.30	0.35	0.29	0.29	0.29	0.29
[SII]	6720+	0.57	0.64	0.55	0.61	0.52	0.52
[SIII]	18.7 μ	0.31	0.27	0.36	0.17	0.37	0.37
[SIII]	34 μ	0.51	0.47	0.60	0.27	0.61	0.62
[SIII]	9532+	0.57	0.48	0.55	0.64	0.60	0.56
L(total)	E36	21.3	21.3	21.7	20.7	21.0	21.8
T(in)		6860	6952	6749	6980	6870	6747
T(H+)		6767	6740	6742	6950	6660	6742
<He+>/<H+>		0.047	0.041	0.044	0.068	0.048	0.034
R(out)	E18	8.96	8.93	8.94	9.00	8.93	9.00

Table 2 Cool HII Region vs Cloudy

		Mean	STD	90.05	94.00	95.00
L(H β)	E36	4.90	0.06	4.91	5.03	4.91
[NII]	6584	0.65	0.05	0.64	0.63	0.56
[OII]	3727+	1.21	0.07	1.08	1.04	1.08
[NeII]	12.8 μ	0.30	0.03	0.26	0.26	0.26
[SII]	6720+	0.57	0.05	0.45	0.44	0.49
[SIII]	18.7 μ	0.31	0.09	0.43	0.45	0.44
[SIII]	34 μ	0.51	0.15	0.61	0.64	0.91
[SIII]	9532	0.43	0.05	0.42	0.42	0.36
L(total)	E36	21.29	0.45	21.0	21.7	21.5
T(in)		6860	110	7261	7188	7185
T(H+)		6767	108	6600	6530	6560
<He+>/<H+>		0.047	0.01	0.048	0.043	0.047
R(out)	E18	8.96	0.04	8.81	8.92	8.85

9.3 Paris HII Region

This compares current predictions of the code with those of other participants at the Meudon meeting on photoionization calculations for the case of a simple spherical HII region. The input used to generate this model HII region is contained in the sample input file *parishii.in*.

Table 3 Paris Meeting HII Region

		Meu	Lex	Ferland	Harrington	Netzer	Pequignot	Rubin
L(H β)	E37	2.06	2.04	2.06	2.04	2.02	2.02	2.05
HeI	5876	0.116	0.111	0.109	0.119	0.101	0.116	
CII	2326+	0.17	0.17	0.19	0.17	0.16	0.14	0.18
CIII]	1909+	0.051	0.07	0.059	0.059	0.078	0.065	0.076
[NII]	122 μ		0.031	0.033			0.036	0.031
[NII]	6584+	0.73	0.80	0.88	0.74	0.87	0.78	0.73
[NIII]	57 μ	0.30	0.28	0.27	0.29	0.26	0.30	0.30
[OII]	3727+	2.01	2.20	2.19	2.14	2.3	2.11	2.26
[OIII]	51.8 μ	1.10	1.06	1.04	1.11	0.99	1.08	1.08
[OIII]	88.4 μ	1.20	1.20	1.07	1.28	1.16	1.25	1.26
[OIII]	5007+	2.03	2.09	1.93	1.96	2.29	2.17	2.10
[NeII]	12.8 μ	0.21	0.21	0.23	0.19	0.22	0.20	0.20
[NeIII]	15.5 μ	0.44	0.41	0.43	0.43	0.37	0.42	0.42
[NeIII]	3869+	0.096	0.091	0.103	0.086	0.100	0.079	0.087
[SII]	6720+	0.14	0.18	0.23	0.16	0.22	0.17	0.13
[SIII]	18.7 μ	0.55	0.53	0.48	0.56	0.5	0.55	0.58
[SIII]	34 μ	0.93	0.87	0.82	0.89	0.81	0.88	0.94
[SIII]	9532+	1.25	1.31	1.27	1.23	1.48	1.27	1.30
[SIV]	10.5 μ	0.39	0.38	0.37	0.42	0.36	0.41	0.33
Sum	Sum	11.59	11.87	11.58	11.70	12.17	11.90	12.01
L(total)	E37	24.1	24.4	24.1	24.1	24.8	24.3	24.6
T(in)		7378	7655	7815	7741	7670	7650	7399
T(H+)		7992	8052	8064	8047	8000	8060	8087
He+>/<H+>			0.76	0.71	0.77	0.76	0.75	0.83
R(out)	E18	1.45	1.46	1.46	1.46	1.47	1.46	1.46

Table 4 Paris HII Region vs Cloudy

		Mean	STD	76.03	80.06	84.15	90.05	94.00	95.00
L(H β)	E37	2.04	0.02	2.06	2.04	2.01	2.05	2.02	2.04
HeI	5876	0.11	0.01	0.12	0.12	0.12	0.11	0.117	0.119
CII]	2326+	0.17	0.02	0.24	0.17	0.17	0.15	0.145	0.146
CIII]	1909+	0.07	0.01	0.06	0.11	0.09	0.063	0.060	0.065
[NII]	122 μ	0.03	0.00			0.03	0.028	0.028	0.027
[NII]	6584	0.60	0.05	0.58	0.60	0.58	0.58	0.57	0.49
[NIII]	57 μ	0.28	0.02			0.29	0.289	0.297	0.304
[OII]	3727+	2.20	0.08	2.29	2.43	2.31	2.02	1.996	2.06
[OIII]	51.8 μ	1.06	0.05	1.15	1.09	1.11	1.22	1.24	1.23
[OIII]	88.4 μ	1.20	0.09			1.06	1.11	1.13	1.12
[OIII]	5007	1.57	0.11	1.83	1.62	1.50	1.66	1.62	1.61
[NeII]	12.8 μ	0.21	0.02	0.22	0.23	0.23	0.18	0.18	0.18
[NeIII]	15.5 μ	0.41	0.02	0.47	0.45	0.45	0.28	0.29	0.29
[NeIII]	3869	0.0689	0.008	0.098	0.090	0.082	0.063	0.061	0.063
[SII]	6720+	0.18	0.04	0.27	0.29	0.129	0.126	0.12	0.134
[SIII]	18.7 μ	0.53	0.04	0.54	0.53	0.52	0.61	0.61	0.63
[SIII]	34 μ	0.87	0.05			0.88	0.85	0.86	1.24
[SIII]	9532	0.98	0.07	1.04	1.04	1.05	0.90	0.89	0.79
[SIV]	10.5 μ	0.38	0.04	0.09	0.11	0.36	0.53	0.53	0.52
Sum	Sum	11.87	0.24	10.0	9.9	11.8	11.9	11.8	12.08
L(total)	E37	24.44	0.31	20.9	20.4	23.9	24.5	24.3	24.4
T(in)		7655	157			6547	7822	7746	7776
T(H+)		8052	32			7530	7970	7920	7967
<He+>/<H+>		0.76	0.04			0.75	0.72	0.74	0.79
R(out)	E18	1.46	0.00			1.45	1.45	1.46	1.45

9.4 Blister HII Region

This is one of the Lexington Meeting test cases, and is meant to be similar to inner regions of the Orion Nebula and is called **blister.in** in the test suite.

Table 5 Blister HII Region

		Mean	Ferland	Harrington	Netzer	Pequignot	Rubin
I(H β)		4.69	4.59	4.81	4.69	4.67	4.70
HeI	5876	0.12	0.13	0.11	0.12	0.12	
CII	2326+	0.16	0.14	0.20	0.10	0.15	0.23
CII	1335+	0.15	0.17	0.14	0.13	0.16	
CIII]	1909+	0.18	0.22	0.17	0.18	0.15	0.20
[NII]	6584+	0.81	0.58	0.94	0.74	0.90	0.87
[NIII]	57 μ	.033	.035	.033	.033	.032	.034
[OII]	7330+	0.12	0.10	0.13	0.09	0.12	0.14
[OII]	3727+	0.86	0.73	0.98	0.69	0.86	1.04
[OIII]	51.8 μ	0.29	0.31	0.29	0.28	0.28	0.28
[OIII]	5007+	4.18	4.74	3.90	4.40	3.90	3.96
[NeII]	12.8 μ	0.34	0.32	0.33	0.35	0.33	0.35
[NeIII]	15.5 μ	1.06	1.24	1.07	0.96	1.04	1.00
[NeIII]	3869+	0.34	0.48	0.32	0.35	0.26	0.29
[SIII]	18.7 μ	0.33	0.31	0.34	0.31	0.33	0.35
[SIII]	9532+	1.46	1.41	1.46	1.51	1.42	1.53
[SIV]	10.5 μ	0.51	0.54	0.52	0.51	0.53	0.43
Sum	Sum	10.78	11.32	10.80	10.63	10.46	10.71
I(total)		51.3	52.6	52.4	50.4	49.4	50.3
T(in)		7911	8206	7582	8200	8200	7366
T(H+)		8303	8324	8351	8310	8200	8328
<He+>/<H+>		0.85	0.94	0.78	0.93	0.79	0.84
ΔR	E17	2.99	2.88	3.08	2.93	2.98	3.09

Table 6 Blister HII Region vs Cloudy

		Mean	STD	90.05	94.00	95.00
I(H β)		4.69	0.08	4.71	4.64	4.91
HeI	5876	0.12	0.01	0.129	0.132	0.132
CII]	2326+	0.16	0.05	0.133	0.124	0.135
CII	1335+	0.15	0.02	0.179	0.171	0.18
CIII]	1909+	0.18	0.03	0.222	0.234	0.26
[NII]	6584	0.54	0.09	0.468	0.417	0.40
[NIII]	57 μ	0.03	0.00	0.037	0.038	0.038
[OII]	7330+	0.12	0.02	0.115	0.113	0.127
[OII]	3727+	0.86	0.15	0.737	0.713	0.794
[OIII]	51.8 μ	0.29	0.01	0.303	0.309	0.307
[OIII]	5007	3.13	0.28	3.74	3.84	3.88
[NeII]	12.8 μ	0.34	0.01	0.140	0.141	0.14
[NeIII]	15.5 μ	1.06	0.11	0.458	0.462	0.466
[NeIII]	3869	0.34	0.08	0.157	0.159	0.166
[SIII]	18.7 μ	0.255	0.015	0.301	0.304	0.349
[SIII]	9532	1.09	0.037	0.950	.955	0.87
[SIV]	10.5 μ	0.51	0.04	0.68	0.68	0.665
Sum	Sum	10.78	0.32	10.6	10.7	10.8
I(total)		51.29	1.39	50.9	50.2	50.0
T(in)		7910	406	8447	8370	8430
T(H+)		8302	59	8325	8330	8429
<He+>/<H+>		0.85	0.07	0.895	0.90	0.916
ΔR	E17	2.99	0.09	2.99	2.94	2.94

9.5 Paris Planetary Nebula

This compares current predictions of the code with those of other participants at the Meudon (1985) and Lexington (1993) meetings on photoionization calculations, for the case of ionization by a very hot black body. The input used to generate this model planetary nebula is shown in the sample input section and is called **parispn.in** in the test suite. The model results are very sensitive to the detailed transfer of HeII $L\alpha$; this line is the dominant heat source across the He⁺⁺ region of the model nebula. The parameters were chosen to be roughly similar to NGC 7027, a very well studied object.

Table 7 Paris Meeting Planetary Nebula

Line		Meudon	Lexington	Ferland	Harrington	Netzer	Pequignot
L(H β)	E35	2.60	2.68	2.63	2.68	2.73	2.68
HeII (35)	erg/s	0.87	0.88	0.83	0.88	0.94	0.85
He I	5876	0.11	0.11	0.11	0.10	0.10	0.11
He II	4686	0.33	0.33	0.32	0.33	0.35	0.32
C II]	2326+	0.38	0.33	0.33	0.43	0.27	0.30
C III]	1909+	1.70	1.77	1.82	1.66	1.72	1.87
C IV	1549+	1.64	2.33	2.44	2.05	2.66	2.18
[N II]	6584+	1.44	1.49	1.59	1.45	1.47	1.44
N III]	1749+	0.11	0.12	0.13	0.13	0.11	0.13
[NIII]	57 μ		0.13	0.12	0.13	0.13	0.13
N IV]	1487+	0.12	0.19	0.20	0.15	0.21	0.19
N V	1240+	0.09	0.17	0.18	0.12	0.23	0.15
[O I]	6300+	0.15	0.14	0.15	0.12	0.14	0.14
[O II]	3727+	2.23	2.25	2.23	2.27	2.31	2.18
[O III]	5007+	20.9	20.76	21.1	21.4	19.4	21.1
[O III]	4363	0.16	0.15	0.16	0.16	0.14	0.16
[O III]	52 μ	1.43	1.43	1.42	1.44	1.40	1.46
[O IV]	26 μ	3.62	3.67	3.52	3.98	3.32	3.86
O IV]	1403+	0.13	0.26	0.20	0.23	0.26	0.33
O V]	1218+	0.09	0.20	0.20	0.11	0.29	0.19
[Ne III]	15.5 μ	2.51	2.78	2.75	2.76	2.80	2.81
[Ne III]	3869+	2.59	2.69	3.33	2.27	2.74	2.44
Ne IV]	2423+	0.56	0.78	0.72	0.74	0.91	0.74
[Ne V]	3426+	0.73	0.67	0.74	0.60	0.73	0.61
[Ne V]	24.2 μ	1.67	0.87	0.94	0.76	0.81	0.99
Mg II	2798+	1.48	1.58	2.33	1.60	1.22	1.17
[Mg IV]	4.5 μ	0.09	0.12	0.12	0.13		0.12
[Si II]	34.8 μ	0.13	0.19	0.16	0.26	0.19	0.17
Si II]	2335+	0.11	0.16	0.15		0.16	0.16
Si III]	1892+	0.20	0.41	0.39	0.32	0.46	0.45
Si IV	1397+	0.15	0.18	0.20	0.15	0.21	0.17
[S II]	6720+	0.39	0.35	0.21	0.45	0.33	0.43
[S III]	18.7 μ	0.49	0.48	0.48	0.49	0.46	0.49
[S III]	9532+	2.09	1.96	2.04	1.89	2.05	1.87
[S IV]	10.5 μ	1.92	1.98	1.92	2.21	1.81	1.98
L(total)	E35	129	137	139	136	135	136
T(in)	E4		1.81	1.83	1.78	1.84	1.78
T(H+)	E4		1.25	1.22	1.21	1.35	1.21
<He+>/<H+>			0.72	0.74	0.74	0.71	0.71
R(out)	E17		4.06	4.04	4.04	4.07	4.07

Table 8 Paris Planetary vs Cloudy

Line		Lexington	STD	74.23	76.03	80.06	84.15	90.05	94.00	95.00
L(H β)	E35	2.68	0.04	2.57	2.66	2.52	2.34	2.62	2.55	2.56
HeII (35)	erg/s	0.88	0.05				0.83	0.83	0.866	.873
He I	5876	0.11	0.01	0.11	0.11	0.11	0.11	0.111	0.108	0.108
He II	4686	0.33	0.01	0.29	0.32	0.35	0.36	0.319	0.343	0.345
C II]	2326+	0.33	0.07	0.36	0.35	0.37	0.35	0.286	0.289	0.294
C III]	1909+	1.77	0.09	1.57	1.48	1.72	1.72	1.88	1.83	1.90
C IV	1549+	2.33	0.27	2.24	2.76	2.48	2.19	2.65	2.63	2.61
[N II]	6584	1.12	0.05	1.06	1.05	1.08	1.11	0.978	1.00	0.885
N III]	1749+	0.12	0.01	0.10	0.08	0.10	0.11	0.123	0.118	0.121
[NIII]	57 μ	0.13	0.00				0.12	0.128	0.125	0.127
N IV]	1487+	0.19	0.03	0.16	0.12	0.11	0.15	0.242	0.238	0.241
N V	1240+	0.17	0.05	0.14	0.09	0.06	0.09	0.175	0.18	0.18
[O I]	6300	0.105	0.007	0.11	0.11	0.12	0.12	0.116	0.118	0.119
[O II]	3727+	2.25	0.06	2.24	2.19	2.35	2.40	2.22	2.32	2.36
[O III]	5007	15.6	0.68	15.9	15.8	15.3	15.6	17.0	16.3	16.4
[O III]	4363	0.15	0.01	0.14	0.13	0.15	0.16	0.174	0.166	0.167
[O III]	52 μ	1.43	0.03	1.40	1.35	1.37	1.35	1.32	1.28	1.28
[O IV]	26 μ	3.67	0.30			3.42	3.65	3.42	3.59	3.64
O IV]	1403+	0.26	0.05	0.19	0.22	0.11	0.15	0.248	0.25	0.225
O V]	1218+	0.20	0.07	0.17	0.11	0.07	0.11	0.200	0.20	0.20
[Ne III]	15.5 μ	2.78	0.03	2.77	2.70	2.67	2.71	1.90	1.87	1.87
[Ne III]	3869	2.01	0.35	2.41	2.25	2.43	2.49	2.15	2.09	2.11
Ne IV]	2423+	0.78	0.09	0.62	0.51	0.51	0.63	0.823	0.827	0.829
[Ne V]	3426	0.50	0.06	0.48	0.40	0.40	0.48	0.589	0.614	0.621
[Ne V]	24.2 μ	0.87	0.11	0.24	0.25	1.01	1.04	1.03	1.09	1.098
Mg II	2798+	1.58	0.54	0.83	1.82	1.96	2.33	2.26	2.27	2.29
[Mg IV]	4.5 μ	0.12	0.00	0.12	0.13	0.14	0.13	0.118	0.12	0.123
[Si II]	34.8 μ	0.19	0.04	0.16	0.16	0.16	0.17	0.157	0.16	0.159
Si II]	2335+	0.16	0.01	0.15	0.14	0.16	0.18	0.150	0.15	0.173
Si III]	1892+	0.41	0.07	0.32	0.42	0.42	0.42	0.526	0.49	0.499
Si IV	1397+	0.18	0.03	0.17	0.24	0.22	0.15	0.235	0.223	0.229
[S II]	6720+	0.35	0.11	0.38	0.68	0.66	0.36	0.354	0.354	0.356
[S III]	18.7 μ	0.48	0.02	0.58	0.71	0.67	0.47	0.467	0.472	0.556
[S III]	9532	1.47	0.075	1.27	1.58	1.55	1.48	1.34	1.35	1.17
[S IV]	10.5 μ	1.98	0.17	1.64	1.32	1.53	1.78	2.20	2.04	2.04
L(total)	E35	137	1.60	117	124	128	121	135		
T(in)	E4	1.81	0.03				1.49	1.828	1.80	1.808
T(H+)	E4	1.25	0.07				1.28	1.22	1.22	1.23
<He+>/<H+>		0.72	0.02				0.72	0.74	0.72	0.72
R(out)	E17	4.06	0.02				3.90	4.03	3.99	4.00

9.6 Paris NLR Model

This compares current predictions of the code with those of other participants at the Meudon meeting on photoionization calculations, for a model similar to the NLR of active nuclei. Results for other codes are from the 1985 Meudon meeting. The input stream is the file *parisnlr.in* in the test suite.

Table 9 Paris Meeting NLR Model

Line		Netzer	Pequignot	Binette	Kraemer	Mean
H β	erg/s/cm ²	0.129	0.134	0.124	0.12	0.127 \pm 0.006
H β	4861	1.00	1.00	1.00	1.00	1.00
L α	1216	35.3	33.1	-	24.0	30.8 \pm 6.0
He I	5876	0.095	0.098	0.092	0.090	0.094 \pm 0.004
He II	4686	0.36	0.32	0.38	0.37	0.358 \pm 0.026
C II]	2326	0.96	0.77	1.70	1.06	1.12 \pm 0.40
C II	1335	0.14	0.14	0.20	0.08	0.14 \pm 0.05
C III]	1909	4.59	4.99	6.50	4.91	5.25 \pm 0.85
C IV	1549	7.03	7.20	5.30	7.20	6.68 \pm 0.93
[N I]	5200	0.31	0.33	0.82	0.37	0.46 \pm 0.24
[N II]	6548	2.68	1.52	1.77	1.63	1.90 \pm 0.53
N III]	1749	0.40	0.40	0.43	0.48	0.428 \pm 0.038
N IV]	1487	0.45	0.43	0.51	0.48	0.468 \pm 0.035
N V	1240	0.32	0.30	0.32	0.28	0.305 \pm 0.019
[O I]	63.2 μ m	-	0.62	0.14	0.10	0.29 \pm 0.29
[O I]	6300	1.32	0.90	1.62	1.04	1.22 \pm 0.32
[O II]	7325	0.11	0.094	0.16	0.10	0.116 \pm 0.03
[O II]	3727	3.4	2.62	4.41	2.73	3.29 \pm 0.82
[O III]	52 μ m	2.5	2.54	2.31	2.65	2.50 \pm 0.14
[O III]	5007	27.36	27.36	23.28	27.76	26.44 \pm 2.11
[O III]	4363	0.42	0.41	0.44	0.44	0.428 \pm 0.015
O III]	1663	0.97	0.95	0.92	1.01	0.963 \pm 0.038
[O IV]	25.9 μ m	5.69	5.19	5.49	-	5.46 \pm 0.25
O IV]	1403	0.53	0.44	0.51	0.66	0.534 \pm 0.092
O V]	1218	0.33	0.32	0.45	0.24	0.335 \pm 0.086
O VI	1035	0.17	0.17	0.22	0.10	0.165 \pm 0.049
[Ne II]	12.8 μ m	0.28	0.18	0.48	0.13	0.268 \pm 0.155
[Ne III]	15.5 μ m	2.8	2.62	1.83	1.25	2.13 \pm 0.72
[Ne III]	3869	2.70	2.59	2.27	1.67	2.31 \pm 0.46
Ne IV]	2423	0.82	0.79	1.03	1.12	0.94 \pm 0.16
[Ne V]	24.2 μ m	3.54	2.64	3.54	-	3.24 \pm 0.52
[Ne V]	3426	1.17	1.02	1.13	1.05	1.095 \pm 0.066
Mg II	2798	1.58	1.43	1.51	1.10	1.40 \pm 0.21
Si II	34.8 μ m	1.73	0.97	0.51	-	1.07 \pm 0.62
Si II	2335	0.21	0.17	0.09	-	0.16 \pm 0.06
Si III]	1892	0.15	0.19	0.69	0.14	0.29 \pm 0.26
Si IV	1397	0.21	0.14	0.02	0.13	0.13 \pm 0.08
S II	6720	1.00	0.62	1.29	0.37	0.82 \pm 0.41
S II	4070	0.07	0.04	0.078	0.03	0.055 \pm 0.023
S III	18.7 μ m	0.75	0.49	0.68	0.65	0.64 \pm 0.11
S III	9532	2.25	1.38	1.73	1.62	1.74 \pm 0.37
S IV	10.5 μ m	1.39	0.73	0.94	1.57	1.16 \pm 0.39

Table 10 Paris NLR Model vs CLOUDY

Line		Mean	84.15	90.05	94.00	95.00
H β	erg/s/cm ²	0.127 \pm 0.006	0.133	0.136	0.131	0.136
H β	4861	1.00	1.00	1.00	1.00	1.00
L α	1216	30.8 \pm 6.0	32.3	32.3	33.8	33.0
He I	5876	0.094 \pm 0.004	0.104	0.103	0.100	0.107
He II	4686	0.358 \pm 0.026	0.351	0.34	0.330	0.330
C II]	2326	1.12 \pm 0.40	0.766	0.652	0.693	0.754
C II	1335	0.14 \pm 0.05	0.126	0.141	0.142	0.145
C III]	1909	5.25 \pm 0.85	5.02	4.64	4.54	4.49
C IV	1549	6.68 \pm 0.93	8.42	7.47	7.15	6.84
[N I]	5200	0.46 \pm 0.24	0.14	0.144	0.151	0.154
[N II]	6548	1.90 \pm 0.53	2.32	2.36	2.50	2.52
N III]	1749	0.428 \pm 0.038	0.45	0.388	0.376	0.369
N IV]	1487	0.468 \pm 0.035	0.553	0.544	0.518	0.500
N V	1240	0.305 \pm 0.019	0.391	0.302	0.272	0.269
[O I]	63.2 μ m	0.29 \pm 0.29	0.36	0.439	0.464	0.476
[O I]	6300	1.22 \pm 0.32	1.02	1.02	1.08	1.19
[O II]	7325	0.116 \pm 0.03	0.111	0.106	0.117	0.125
[O II]	3727	3.29 \pm 0.82	2.99	2.69	2.97	3.32
[O III]	52 μ m	2.50 \pm 0.14	2.34	2.23	2.20	2.10
[O III]	5007	26.44 \pm 2.11	25.3	26.05	25.4	24.4
[O III]	4363	0.428 \pm 0.015	0.45	0.441	0.428	0.414
O III]	1663	0.963 \pm 0.038	1.05	1.03	1.00	0.973
[O IV]	25.9 μ m	5.46 \pm 0.25	5.76	5.91	5.79	5.59
O IV]	1403	0.534 \pm 0.092	0.54	0.517	0.490	0.419
O V]	1218	0.335 \pm 0.086	0.453	0.292	0.264	0.260
O VI	1035	0.165 \pm 0.049	0.222	0.142	0.120	0.123
[Ne II]	12.8 μ m	0.268 \pm 0.155	0.220	0.168	0.172	0.227
[Ne III]	15.5 μ m	2.13 \pm 0.72	3.13	2.12	2.12	2.13
[Ne III]	3869	2.31 \pm 0.46	3.99	3.22	3.19	3.12
Ne IV]	2423	0.94 \pm 0.16	1.19	1.16	1.12	1.09
[Ne V]	24.2 μ m	3.24 \pm 0.52	2.74	2.69	2.69	2.59
[Ne V]	3426	1.095 \pm 0.066	1.45	1.25	1.20	1.17
Mg II	2798	1.40 \pm 0.21	1.76	1.66	1.62	1.79
Si II	34.8 μ m	1.07 \pm 0.62	1.12	1.03	1.07	1.20
Si II	2335	0.16 \pm 0.06	0.218	0.191	0.201	0.239
Si III]	1892	0.29 \pm 0.26	0.401	0.501	0.474	0.464
Si IV	1397	0.13 \pm 0.08	0.140	0.215	0.188	0.185
[S II]	6720	0.82 \pm 0.41	1.50	0.988	0.958	1.21
[S II]	4070	0.055 \pm 0.023	0.1004	0.095		0.114
[S III]	18.7 μ m	0.64 \pm 0.11	0.625	0.685	0.715	0.837
[S III]	9532	1.74 \pm 0.37	1.92	1.67	1.74	1.57
[S IV]	10.5 μ m	1.16 \pm 0.39	1.73	1.44	1.32	1.27

9.7 Lexington NLR Model

This is the NLR model computed at the 1994 Lexington meeting, and is called `nlr.in` in the test suites.

Table 11 Lexington NLR Model

		Lexington	Binette	Ferland	Netzer	Pequignot	Viegas
I(H β)	E0	1.36	1.33	1.31	1.37	1.43	1.34
Ly α	1216	33.70	38.3	32.1	32.4	31.5	34.2
HeI	5876	0.12	0.11	0.13	0.12	0.13	0.13
HeII	4686	0.24	0.25	0.25	0.25	0.23	0.24
HeII	1640	1.62	1.60	1.74	1.53	1.56	1.67
CIII]	1909+	2.90	2.90	2.99	2.87	2.83	2.90
CIV	1549+	3.35	2.70	3.85	3.69	3.17	3.36
[NII]	6584+	2.55	1.40	3.20	3.10	2.67	2.40
NIII]	1749+	0.23	0.24	0.24	0.22	0.22	0.22
NIV]	1487+	0.21	0.20	0.23	0.22	0.21	0.21
[OI]	6300+	1.65	2.20	1.61	1.67	1.31	1.46
[OI]	63 μ	1.12	0.25	1.13		1.44	1.64
[OII]	3727+	1.42	1.60	1.44	1.58	1.30	1.20
OIII]	1663+	0.56	0.35	0.63	0.61	0.57	0.63
[OIII]	5007+	33.54	31.4	34.5	33.0	32.8	36.0
[OIII]	4363	0.32	0.30	0.34	0.31	0.30	0.33
OIV	1403+	0.36	0.49	0.30	0.36	0.42	0.25
[NeIII]	15.5 μ	1.89	1.50	2.01	1.94	2.05	1.95
[NeIII]	3869+	2.13	1.90	2.51	2.16	1.72	2.34
[Ne IV]	2423+	0.44	0.52	0.42	0.47	0.41	0.38
[NeV]	3426+	0.52	0.59	0.55	0.53	0.44	0.50
MgII	2798+	1.84	3.50	1.72	1.23	1.12	1.61
[SIII]	34.8 μ	0.90	1.00	0.96	1.07	0.96	0.52
[SII]	6720+	1.29	2.40	1.01	0.93	0.99	1.10
[SIII]	9532+	1.91	1.60	2.15	2.06	1.67	2.08
[SIII]	18.7 μ	0.49	0.36	0.61	0.57	0.52	0.37
[SIV]	10.5 μ	1.02	0.86	1.24	0.82	0.94	1.22
I(total)	E0	130	131	128	128	131	133
T(in)	E4	1.70	1.71	1.70	1.72	1.68	1.68
T(H+)	E4	1.18		1.24	1.06	1.20	1.23

Table 12 Lexington NLR vs Cloudy

		Mean	STD	90.05	94.01	95.00
I(H β)	E0	1.36	0.05	1.38	1.33	1.39
L α	1216	33.70	2.76	30.0	31.6	29.7
HeI	5876	0.12	0.01	0.128	0.12	0.129
HeII	4686	0.24	0.01	0.24	0.24	0.232
HeII	1640	1.62	0.09	1.76	1.68	1.63
CIII]	1909+	2.90	0.06	2.62	2.55	2.52
CIV	1549+	3.35	0.45	3.99	3.78	3.64
[NII]	6584	1.91	0.52	2.26	2.39	2.448
NIII]	1749+	0.23	0.01	0.20	0.20	0.192
NIV]	1487+	0.21	0.01	0.286	0.27	0.260
[OI]	6300+	1.65	0.34	1.63	1.73	1.96
[OI]	63 μ	1.12	0.61	0.77	0.79	0.784
[OII]	3727+	1.42	0.17	1.33	1.45	1.60
OIII]	1663+	0.56	0.12	0.62	0.60	0.584
[OIII]	5007+	33.54	1.76	34.4	33.4	32.1
[OIII]	4363	0.32	0.02	0.327	0.32	0.306
OIV	1403+	0.36	0.09	0.35	0.38	0.326
[NeIII]	15.5 μ	1.89	0.22	1.43	1.44	1.42
[NeIII]	3869+	2.13	0.32	1.96	1.98	1.92
[Ne IV]	2423+	0.44	0.06	0.47	0.45	0.437
[NeV]	3426+	0.52	0.06	0.60	0.60	0.544
MgII	2798+	1.84	0.96	1.48	1.48	1.63
[SiII]	34.8 μ	0.90	0.22	0.87	0.90	0.935
[SII]	6720+	1.29	0.63	0.91	0.91	1.18
[SIII]	9532+	1.91	0.26	2.02	2.04	1.85
[SIII]	18.7 μ	0.49	0.12	0.68	0.72	0.928
[SIV]	10.5 μ	1.02	0.20	1.14	1.06	1.02
I(total)	E0	130	1.86	125	125	124
T(in)	E4	1.70	0.02	1.7	1.69	1.70
T(H+)	E4	1.18	0.08	1.24	1.23	1.02

9.8 The DQ Her Shell

This is more or less the model of the DQ Her nebula proposed by Ferland et al. (1984). The input stream for this model is called `dqher.in` in the test suite. The big difference between C90 and previous versions is in the intensity of H β predicted. The code no longer assumes case B when the temperature is too low to do the matrix solution. The nebula is optically thin to many Lyman lines and these escape, robbing flux from H β .

Table 13 The DQ Her Shell

Line		80.09	84.15	90.05	94.00	95.00
H β (+30)	4861	1.65	1.62	0.879	1.07	1.07
Totl	4861			0.649	0.883	0.887
Case B	4861			1.00	1.00	1.00
L α	1216	40.9	20.7	20.6	27.9	28.7
He I	5876	0.786	0.315	0.246	0.274	0.273
He II	4686	0.166	0.085	0.0877	0.052	0.054
C II	158 μ	0.777	0.62	0.092	0.102	0.103
C II	1335		0.062	0.0195	0.021	0.677
[N I]	5200	0.144		0.078	0.087	0.087
[N II]	122 μ	7.85	3.43	3.22	3.58	3.603
[N III]	57 μ	5.30	3.78	9.20	10.2	10.4
[O II]	3727	1.16	0.304	0.294	0.325	0.331
[O III]	88 μ	12.1	6.36	4.95	5.45	5.40
[O III]	52 μ	12.4	6.60	5.51	6.06	6.00
[Si II]	35 μ	0.586	0.79	0.83	0.92	0.93
[S III]	34 μ	0.114	0.167	0.230	0.258	0.398
[Fe II]		0.121	0.187	0.137	0.151	0.154
<Te>	K		643	862	853	860

9.9 The Kwan and Krolik Standard Model

Table 14. The Kwan and Krolik Standard Model.

Line	KK81	80	84.15	90.05	94.00	95.00
L α 1216	5.512+ ⁷⁷	5.78	5.59	6.03	5.98	6.03
L α 1216	100	100	100	100	100	100
H β 4861	10.3	6.37	6.05	5.85	5.55	5.97
H α 6563	42.8	19.9	19.56	24.4	17.2	18.2
BaC	47.0	38.8	39.5	59.6	58.5	64.6
PaC	30.7	20.2	20.0	37.4	34.1	39.8
2s-1s	3.3	2.35	3.19	2.70	2.95	2.78
free-free	29.0	25.3	23.3	40.5	39.2	43.9
He I 10830	4.4	2.86	2.84	3.30	2.96	3.30
He I 5876	0.9	0.697	0.845	0.68	0.64	0.690
He II 4686	-	0.365	0.316	0.29	0.27	0.272
He II 1640	2.5	3.17	2.77	2.49	2.29	2.29
C II] 2326	3.3	2.70	0.686	3.17	3.70	4.83
C II 1335	-	0.661	0.611	0.68	0.71	0.587
C III 977	5.1	7.53	6.52	5.58	5.47	5.76
C III] 1909	13.0	20.7	15.9	13.5	13.2	13.2
C IV 1549	67.0	95.7	76.3	70.2	68.1	68.1
N III] 1750	1.5	4.26	4.95	4.23	4.17	3.67
N III 990	-	0.324	1.35	1.40	1.35	1.35
N IV] 1486	5.8	4.88	5.24	5.27	5.13	5.13
N V 1240	8.4	3.61	4.95	3.78	3.72	3.72
O I 1304	6.9	4.01	3.93	5.13	8.21	6.44
O III] 1663	9.5	18.4	18.3	18.6	18.3	18.6
O IV] 1402	5.2	4.81	5.84	6.20	6.04	5.15
O V] 1218	6.7	3.07	4.63	3.23	3.15	3.25
O VI 1034	15.0	2.85	4.71	2.92	2.85	2.98
Mg II 2798	18.0	17.6	9.71	30.0	33.8	34.6
Si III] 1892	-	14.6	15.7	11.6	12.2	12.6
Si IV 1397	5.5	14.2	9.77	8.54	7.60	7.75
Fe II	23.9	9.58	11.3	13.6	15.5	13.2
Fe K α		2.11	1.74	1.78	1.81	1.80

Table 14 gives the spectrum of the Kwan and Krolik (1981) standard model, called `kk.in` in the test suite.

⁷⁷Line intensity in $\text{erg s}^{-1} \text{cm}^{-2}$. The entries which follow are relative to a scale where $\text{Ly}\alpha=100$.

9.10 Rees, Netzer, and Ferland, low density

This is the lower density model listed in Table 1 of Rees et al. (1989). It is `rnfa.in` in the sample input streams.

Table 15 Rees, Netzer, and Ferland, low density case.

	1989	90.05	94.00	95.00	Ion
log n = 10					
L α	1.000	100	100	100	1.000
H β	0.030	2.76	3.18	3.44	0.026
H α	0.197	13.0	12.1	12.5	0.180
P α	0.022	1.45	1.43	1.31	0.020
Ba C	0.127	13.3	12.1	12.9	0.125
Pa C	0.052	5.01	4.54	4.90	0.025
ff	0.082	8.38	8.24	8.61	0.089
HeI 5876	0.007	0.656	0.584	0.623	0.008
HeII 4686	0.005	0.39	0.392	0.397	0.004
HeII 1640	0.039	3.44	3.38	3.42	0.034
CII 1335	0.009	1.04	1.00	0.880	0.008
CII] 2326	0.014	1.05	0.96	1.59	0.011
CIII 977	0.030	2.59	2.52	2.58	0.035
CIII] 1909	0.106	9.91	9.66	9.96	0.103
CIV 1549	0.424	44.8	42.9	42.3	0.453
NIII] 1750	0.012	1.31	1.26	1.06	0.014
NIV] 1486	0.009	1.19	1.14	1.13	0.009
NV 1240	0.002	0.235	0.22	0.225	0.003
OI 1304	0.033	0.91	0.915	1.27	0.013
OI 8446	0.005	0.0579	0.064	0.066	0.005
OIII] 1663	0.055	6.17	5.92	6.00	0.060
OIV] 1402	0.011	1.41	1.35	1.18	0.017
Mg II 2798	0.076	19.2	18.4	21.5	0.160
SIII 1207	0.012	1.04	1.05	1.07	0.013
SIII] 1892	0.085	8.56	8.51	9.13	0.090
SiIV 1397	0.048	4.47	4.10	4.16	0.044
FeII	0	1.79	2.77	3.46	0.052
K α	0	0.098	0.102	0.103	0.001
sum	2.492	254.2169	248.787	255.764	2.602
L α	4.85E+07	4.54+07	4.42+7	4.39	4.55+07
I(total)	6.03E+07	1.15+8	1.09+8		6.02+07

9.11 Rees, Netzer, and Ferland, high density

Table 16 Rees, Netzer, and Ferland, high density case.

	1989	90.05	94.00	95.00	Ion
log n = 12					
L α	1.000	100	100	100	1.000
H β	0.063	1.99	1.66	1.62	0.054
H α	0.175	5.87	3.00	2.94	0.134
P α	0.009	0.45	0.22	0.212	0.008
Ba C	2.938	230	234	233	2.680
Pa C	1.313	104	104	107	0.590
ff	1.563	114	120	119	1.300
HeI 5876	0.038	3.13	3.42	3.27	0.032
HeII 4686	0.030	2.78	2.76	2.69	0.017
HeII 1640	0.266	17.5	18.3	17.8	0.140
CII 1335	0.082	4.67	4.84	4.34	0.099
CII] 2326	0.003	0.14	0.18	0.312	0.002
CIII 977	0.225	9.95	11.5	11.2	0.150
CIII] 1909	0.023	0.85	0.92	0.895	0.014
CIV 1549	0.938	55.4	62.7	59.9	0.860
NIII] 1750	0.006	0.26	0.28	0.182	0.004
NIV] 1486	0.006	0.35	0.36	0.341	0.004
NV 1240	0.031	2.15	2.34	2.20	0.021
OI 1304	0.033	2.01	1.98	2.05	0.028
OI 8446	0.006	0.26	0.27	0.276	0.005
OIII] 1663	0.039	1.77	1.90	1.82	0.024
OIII 835	0.041	1.90	0.043	2.03	0.002
OIV] 1402	0.023	1.62	1.69	1.60	0.020
Mg II 2798	0.088	12.6	13.6	13.0	0.170
SIII 1207	0.088	4.46	4.70	4.37	0.068
SIII] 1892	0.113	3.37	3.72	6.79	0.061
SIV 1397	0.300	17.5	18.9	18.3	0.230
FeII		0.77	0.99	0.832	0.015
K α		0.38	0.42	0.405	0.003
sum	8.188	700.13	718.693	718.375	7.521
L α	8.54+08	1.17+09	1.08+09	1.11+09	1.08+09
I(total)	7.00+09	8.19+09	7.74+09		7.04+09

This is the higher density model listed in Table 1 of Rees et al. (1989). It is **rnfb.in** in the sample input streams.

10 SAMPLE INPUT

In versions 90 and before this section listed all test cases and discussed the motivations for each. Now each file is self-descriptive, and should be consulted to use as examples and understand how the code is run. Column 1 of Table 17 lists the single-model test cases and the second column gives a brief description of the test.

These tests are included in the distribution, and are the input streams that are used to exercise Cloudy under both standard and extreme conditions.

Perl script to Run all models

Concept of assert commands

Perl script to check results

Table 17 The single-model standard test cases

albedo.in	measure rayleigh scattering of La
amoeba.in	find density and temperature from HST FOS spectra of [NII] in orion
atlas.in	Model of a Compact HII Region
high.in	high case B with largest possible H atom
blister.in	conditions similar to Orion nebula blister
blr89.in	final F+P 1989 BLR model table 3
blr92.in	standard blr cloud in Ferland et al. 1992
brems.in	generate continuum due to hot ism in high Z,z starburst
casea.in	case A
caseblowt.in	log density case B, T=5000, log n=2
casebn2.in	log density case B, T=5000, log n=2
casebn8.in	casebn8 high density case B
casec.in	case C
comphi.in	compton exchange in high temper limit
complo.in	compton exchange near low temperature limit
compton.in	Compton limit, test continuum partition
conserv.in	test that energy is conserved
coolhii.in	cool HII region model from Lexington Meeting
coronal.in	model of active region of solar corona
cosmicray.in	background cosmic ray ionization by suprathermal electrons only
costar.in	test costar continuum
dqher.in	(roughly) Ferland et al. DQ Her model
eden.in	Martin Gaskell's funny model
fe2.in	test feii
feii.in	thermal equilibrium of FeII in STE limit
feiihin.in	test feii
feiihirad.in	feii in case of high radiation density limit
fileopac1.in	part of test of stored file opacities
fileopac2.in	part of test using on-the-fly file opacities
fluc.in	Paris meeting Planetary nebula with density fluctuations
grainlte.in	check that grains equilibrate at correct temp in ste limit
grains.in	test all grain species temperature
hcexc.in	test collisional excitation only, very high density to force H to LTE
hcion.in	test collisional ionization only, no excitation, should be in lte
hemis.in	test hydrogen atom continuous emissivity, used for plot in hazy
hheonlyotssp.in	plane parallel conservation and hydrogenic emission for pure hydrogen
hheonlyotssp.in	pherical conservation and hydrogenic emission for hydrogen and helium
hheonlyoutpp.in	plane parallel conservation and hydrogenic emission for pure hydrogen
hheonlyoutppff.in	plane parallel filling factor for pure hydrogen
hheonlyoutsp.in	spherical conservation and hydrogenic emission for hydrogen and helium

highn.in	high electron density approach to lte
honlyotsopen.in	test ots, inward fractions for pure hydrogen, open geo, filling factor
honlyotspp.in	plane parallel conservation and hydrogenic emission for pure hydrogen
honlyotssp.in	spherical conservation and hydrogenic emission for pure hydrogen
honlyoutopen.in	test inward fractions, open geo, filling factor, 2-photon emission
honlyoutpp.in	plane parallel conservation and hydrogenic emission for pure hydrogen
honlyoutsp.in	spherical conservation and hydrogenic emission for pure hydrogen
induc.in	constant temper black body limit from Ferland and Rees 1988
induchhe.in	half H, He gas with induced BF processes dominate
ism.in	interstellar cloud irradiated by ism background
kk.in	Kwan+Krolik Ap.J. 250, 478 BLR model
lalpha.in	Lyman alpha forest cloud
laser1.in	test of H ionization in optically thin limit
laser2.in	test of H and HeI ionization in optically thin limit
laser3.in	test of H, HeI, and HeII ionization in optically thin limit
level2.in	test dominant level2 lines
liner.in	liner model
lowd0.in	test low density limit, this and lowdm6 should get same results
lowden.in	test optically thin model that extends to very low densities
lowdm6.in	low density test
lte.in	thermal equil black body limit from Ferland and Rees 1988
ltemetl.in	thermodynamic equilibrium with metals
map.in	map of heating vs cooling
newpdr.in	Tielens and Hollenbach pdr model, Table 2, paper b
nlr.in	Hagai's nlr model for Lexington Meeting
nova.in	dense nova shell
oldblr.in	Netzer and Ferland PASP 96, 593, Table 1
optimal.in	find density and temperature from HST FOS spectra of [NII] in orion
parishii.in	"New" Paris meeting HII region
parisnlr.in	paris meeting NLR model
parispn.in	parispn.in Meudon Planetary nebula
pdr.in	Tielens and Hollenbach pdr model, Table 2, paper b
pgrains.in	test all grain species temperature
pnots.in	Paris meeting Planetary nebula with ots
powell.in	find density and temperature from HST FOS spectra of [NII] in orion
primal.in	primordial abundances
qheat.in	grain quantum heating test
rauch.in	very hot PN model
rnfa.in	table 1 of Rees et al. ApJ 347, 648
rnfb.in	table 1 of Rees et al. ApJ 347, 648
sqrden.in	test with density falling as R^{-2} , and filling factor
strom.in	check pure hydrogen Stromgren sphere
subplex.in	find density and temperature from HST FOS spectra of [NII] in orion
test.in	this runs the standard, one command, test, which contains many asserts
transpunch.in	first of transpunch/transread pair, punch continuum
transread.in	first of transpunch/transread pair, punch continuum
vbhum.in	test against Van Blerkom and Hummer, fig 4
werner.in	test run with Werner stellar atmosphere
wind.in	test of equations of motion in a wind
z3.in	redshift 1000 recombination epoch

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